

# Parity Violation Parallel Sessions

Intense Electron Beam Workshop  
Ithaca, NY

June 17-19, 2015

Convenors: M. Perelstein, K. Paschke

<b>Krishna Kumar</b>	<b>Parity-Violating Electron Scattering: Status and Prospects</b>
<b>Charles Horowitz</b>	<b>Nuclear Weak Form-Factors</b>
<b>Hooman Davoudiasl</b>	<b>Dark Z and Parity Violation</b>
<b>Carl Carlson</b>	<b>Calculations of Gamma-Z Box Diagrams</b>
<b>Oscar Moreno</b>	<b>Nuclear and nucleon structure effects in low-energy parity-violating electron scattering</b>
<b>Roger Carlini</b>	<b>Qweak + Torroidal Spectrometer Options</b>
<b>Dominik Becker</b>	<b>Monte Carlo simulations of a solenoid spectrometer for Project P2</b>
<b>Paul Souder</b>	<b>Carbon with Solenoidal Spectrometer</b>

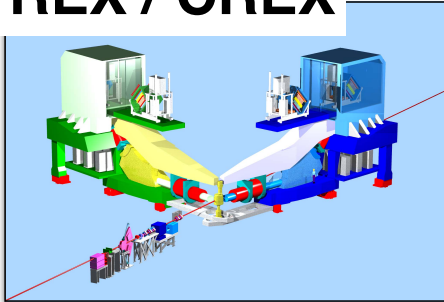
# 3 Topics

- Neutron distribution in heavy nuclei
- Strange form factors
- Standard Model tests and beyond Standard Model reach

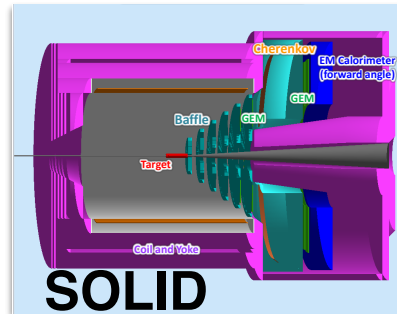
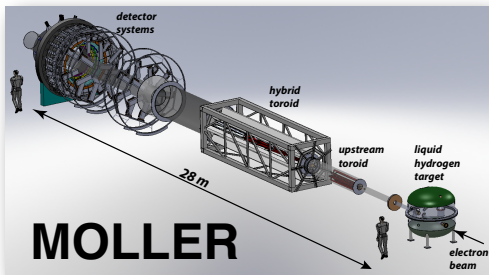
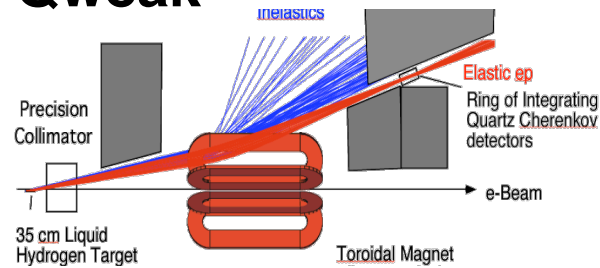
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## PREX / CREX



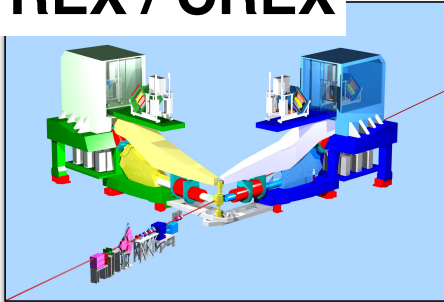
## Qweak



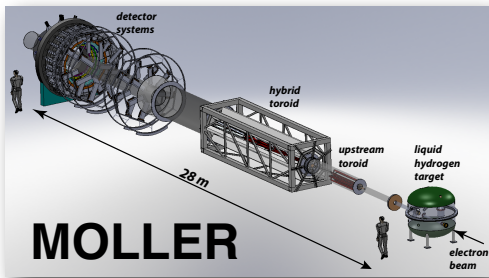
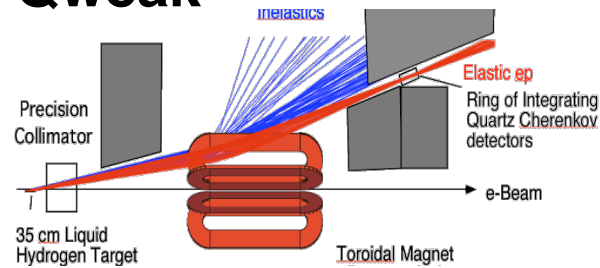
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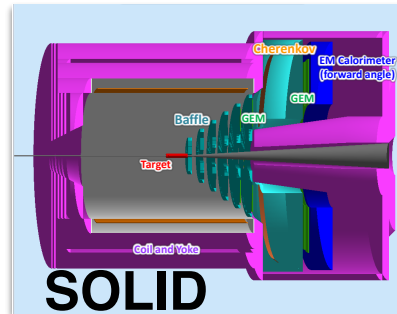
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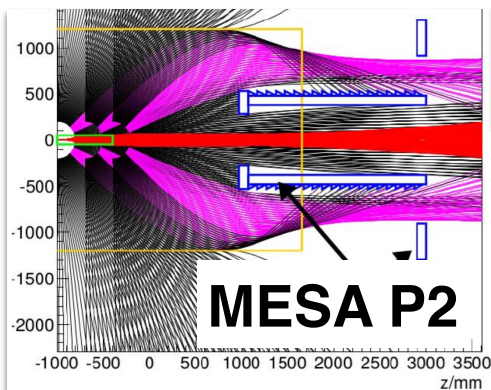
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## MOLLER



## SOLID

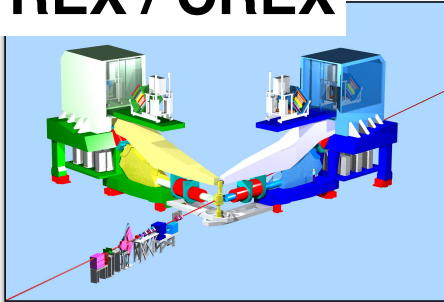


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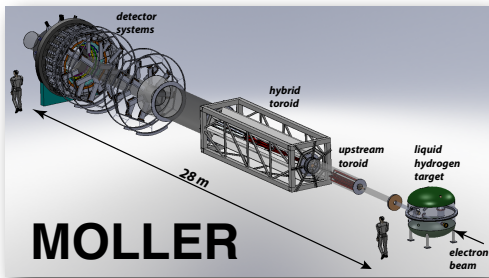
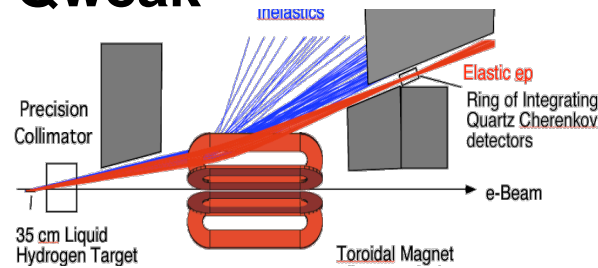
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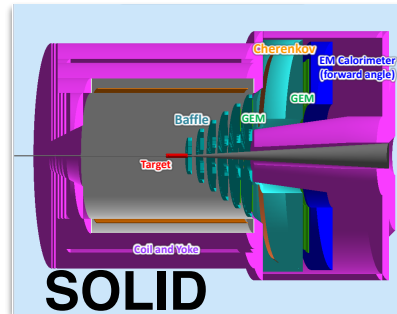
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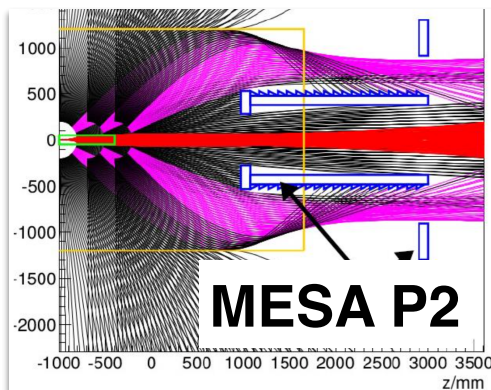
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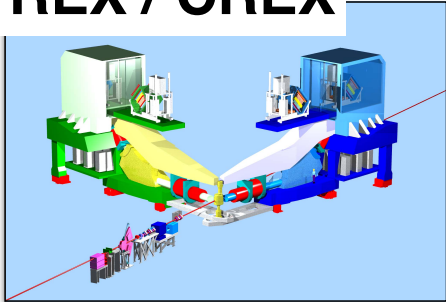
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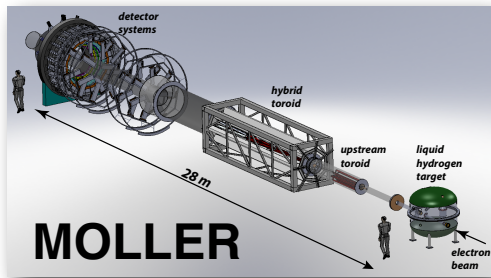
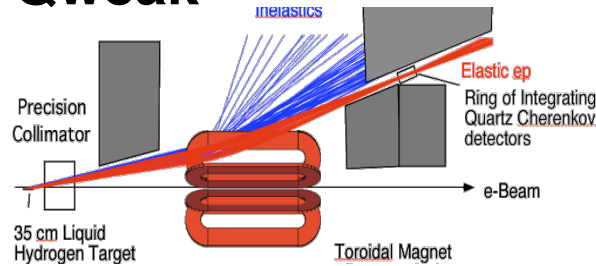
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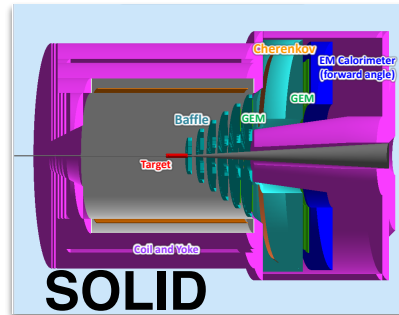
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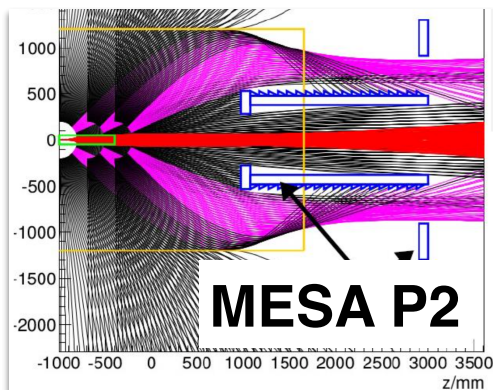
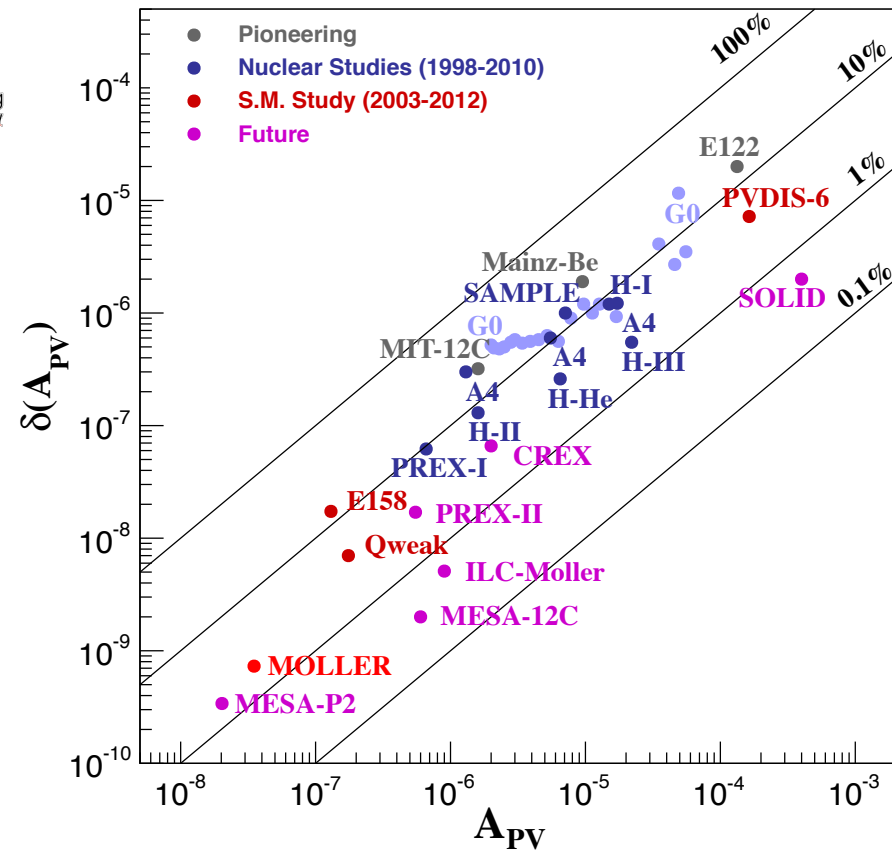
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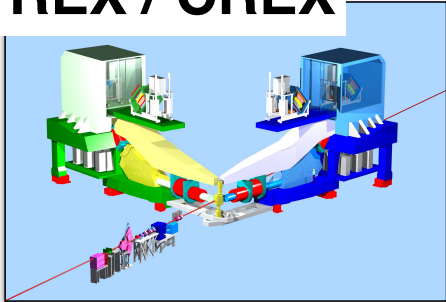
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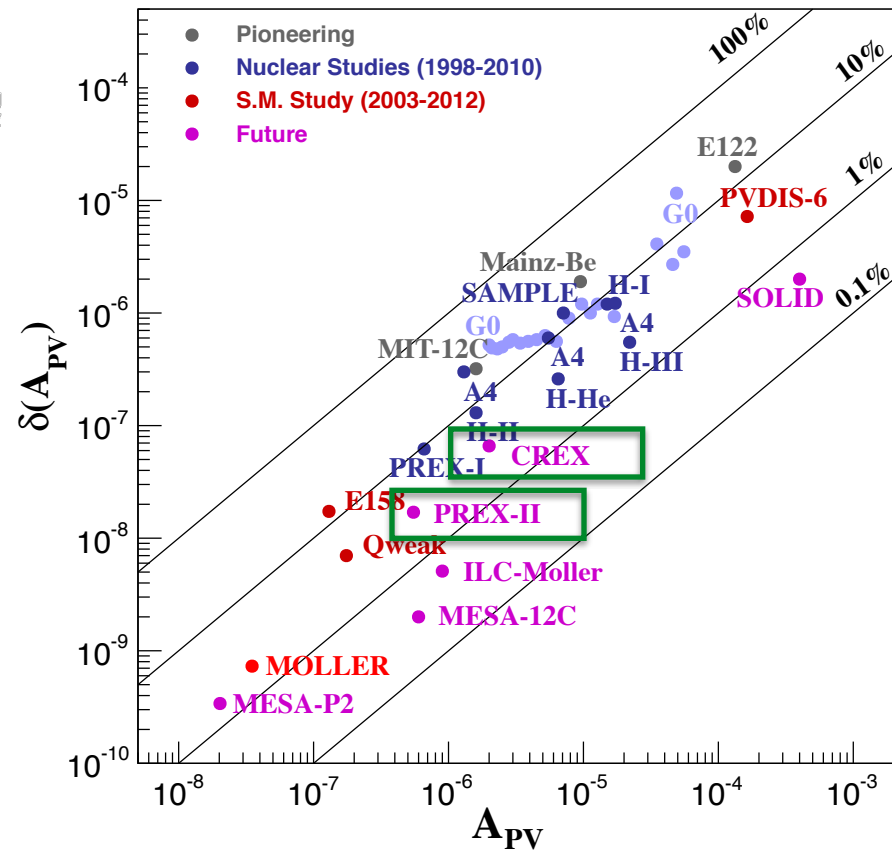
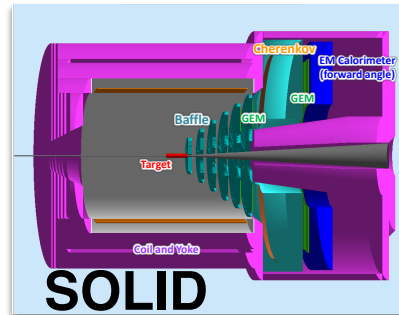
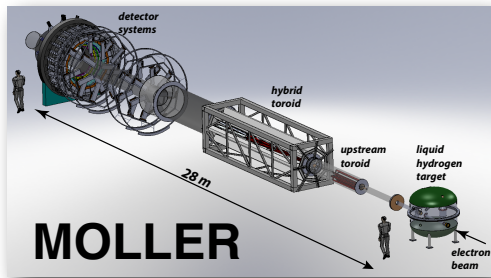
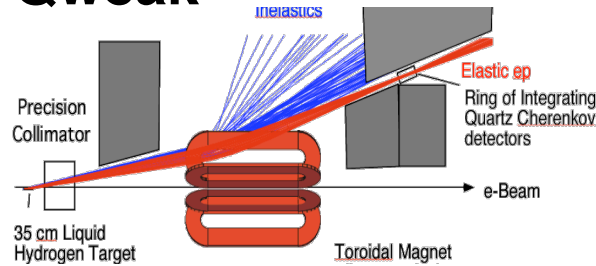
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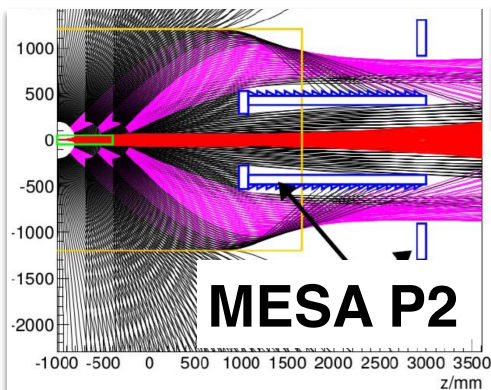
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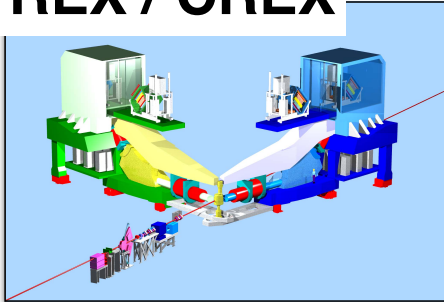
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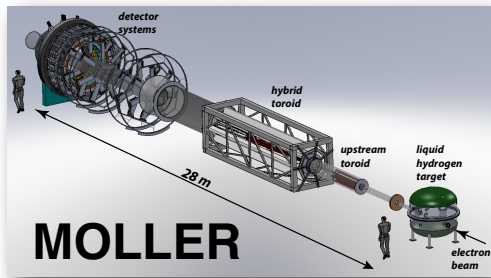
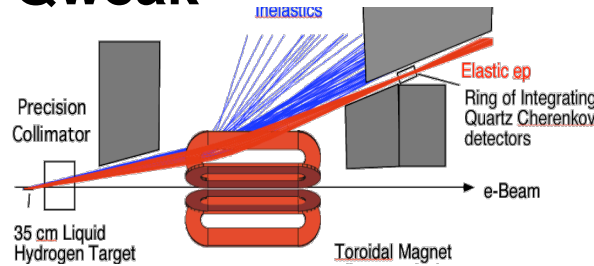
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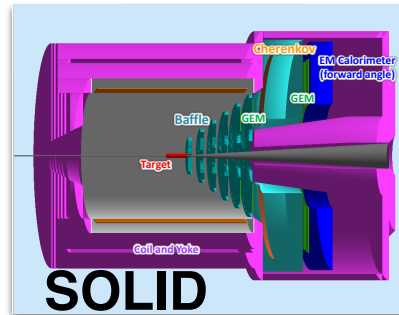
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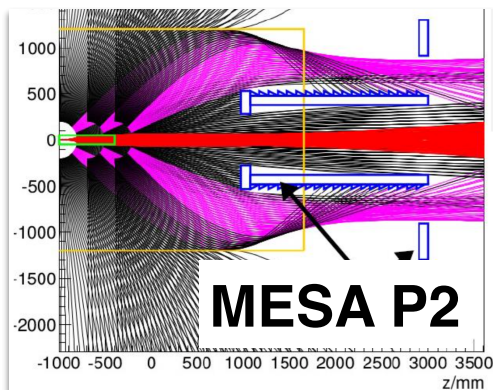
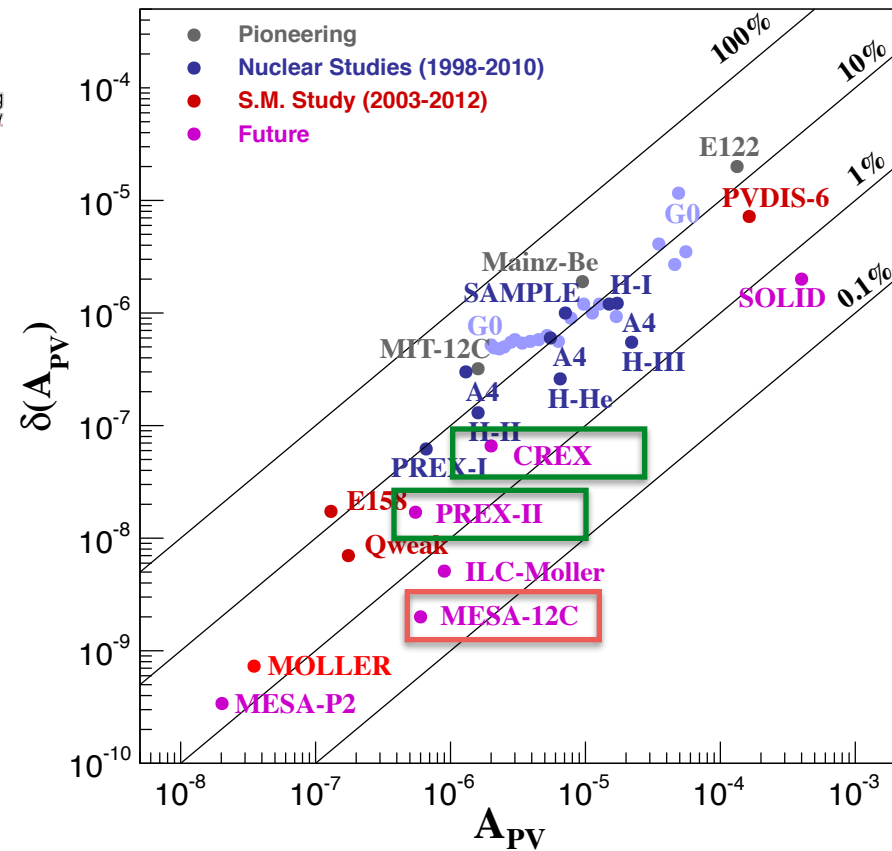
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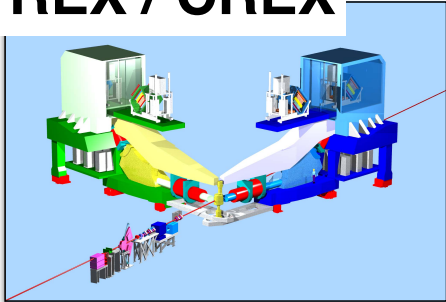
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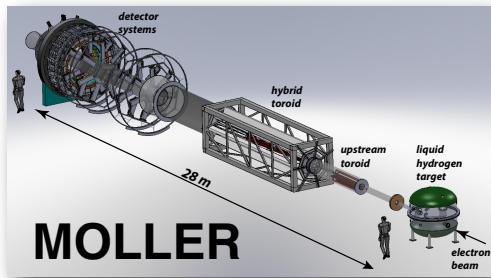
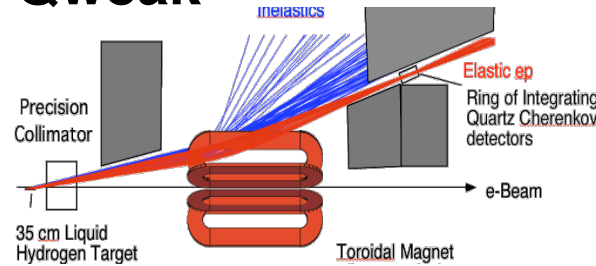
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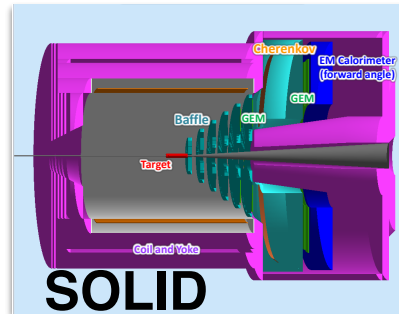
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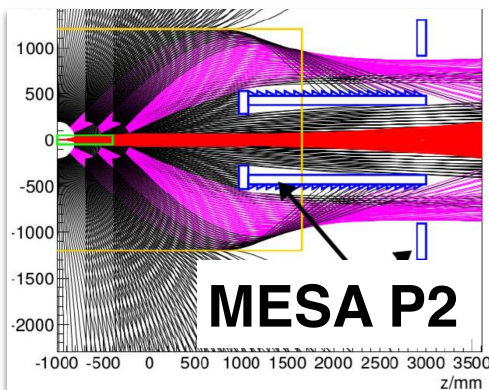
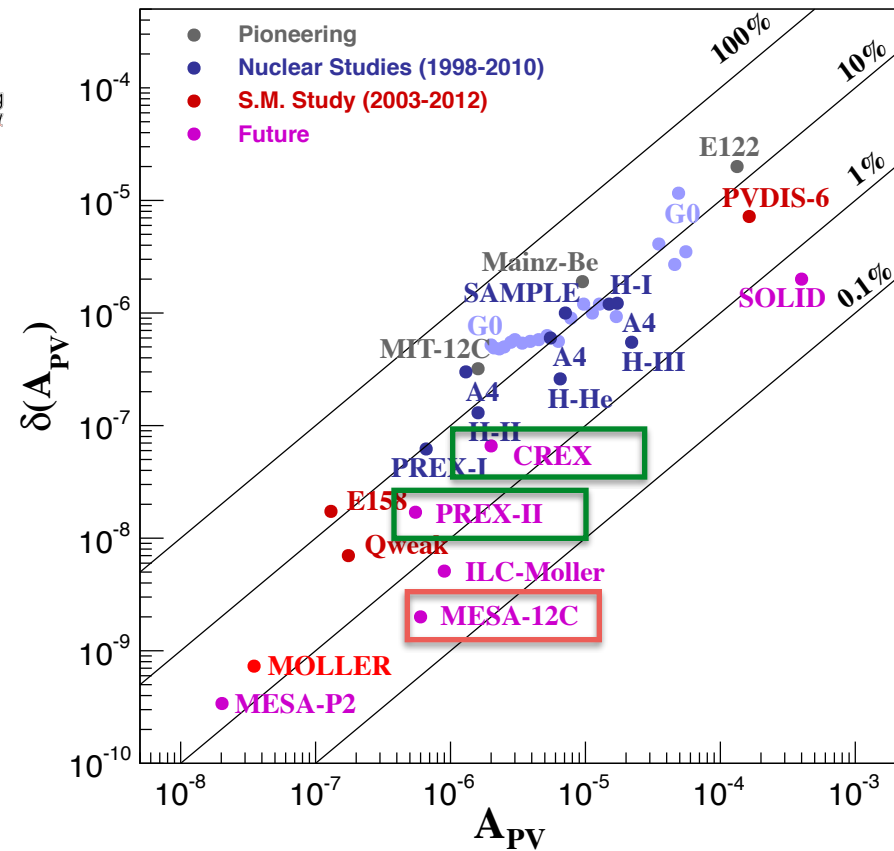
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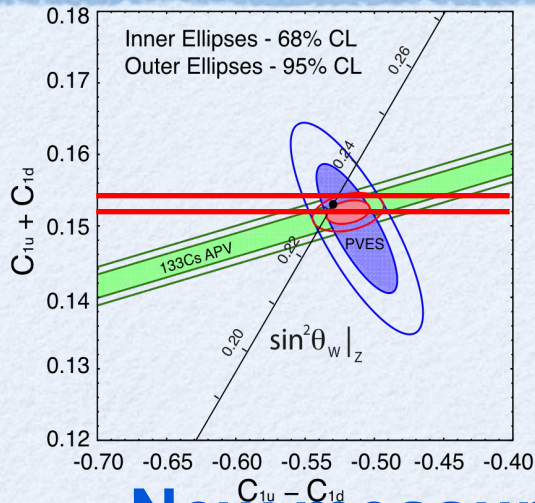
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$$\text{MESA: } E_b = 150 \text{ MeV}$$

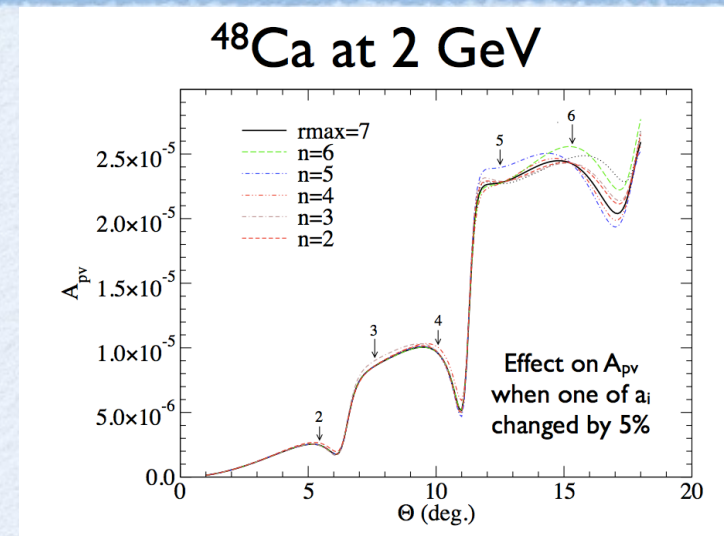
$$I = 150 \mu\text{A}$$

# Potential of New Machine



$$A_{PV} = \frac{G_F Q^2}{\sqrt{2\pi\alpha}} \sin^2 \theta_W$$

for spinless, isoscalar nucleus



## • New measurements on Carbon-12

- *A Standard Model test extremely interesting if 0.3% can be reached*
- *Must be coupled with higher  $Q^2$  measurements to constrain strange quark radius (strange quark contribution to charge radius)*

## • New measurements on Calcium-48

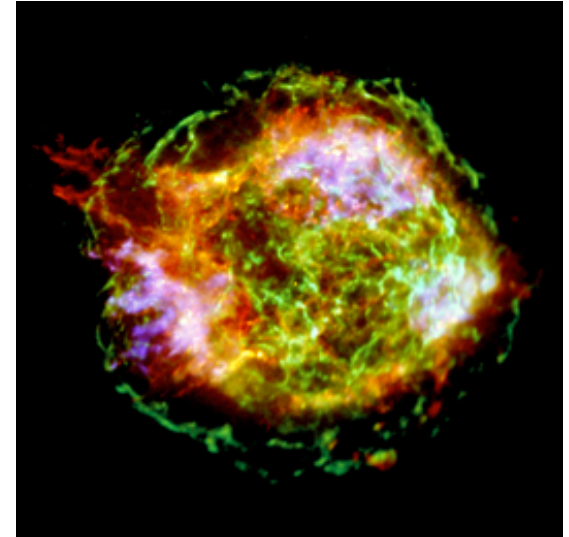
- *CREX will make a very precise low  $Q^2$  measurement*
- *Higher  $Q^2$  measurements will provide a complete and model-independent distribution of neutrons in the ground state*

## • Ideal requirements:

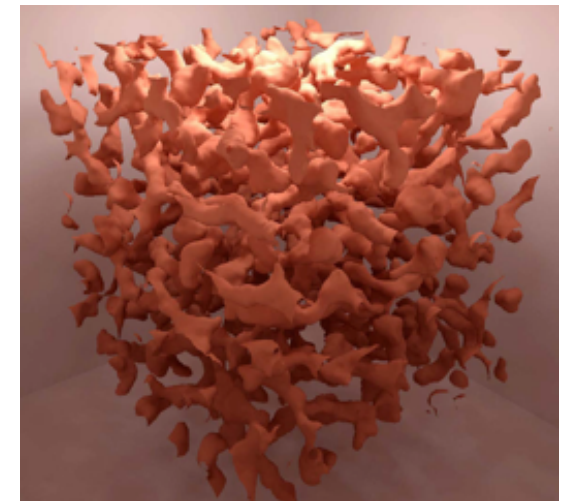
- *several hundred microamps (polarized) with up to 500 MeV*
  - *could do quite a bit with 286 MeV and 100 microamps*

# Neutron Rich Matter

- Compress almost anything to  $10^{11}+$  g/cm<sup>3</sup> and electrons react with protons to make neutron rich matter. This material is at the heart of many fundamental questions in nuclear physics and astrophysics.
  - What are the high density phases of QCD?
  - Where did chemical elements come from?
  - What is the structure of many compact and energetic objects in the heavens, and what determines their electromagnetic, neutrino, and gravitational-wave radiations?
- Interested in neutron rich matter over a tremendous range of density and temperature were it can be a *gas, liquid, solid, plasma, liquid crystal (nuclear pasta), superconductor ( $T_c=10^{10}$  K!), superfluid, color superconductor...*

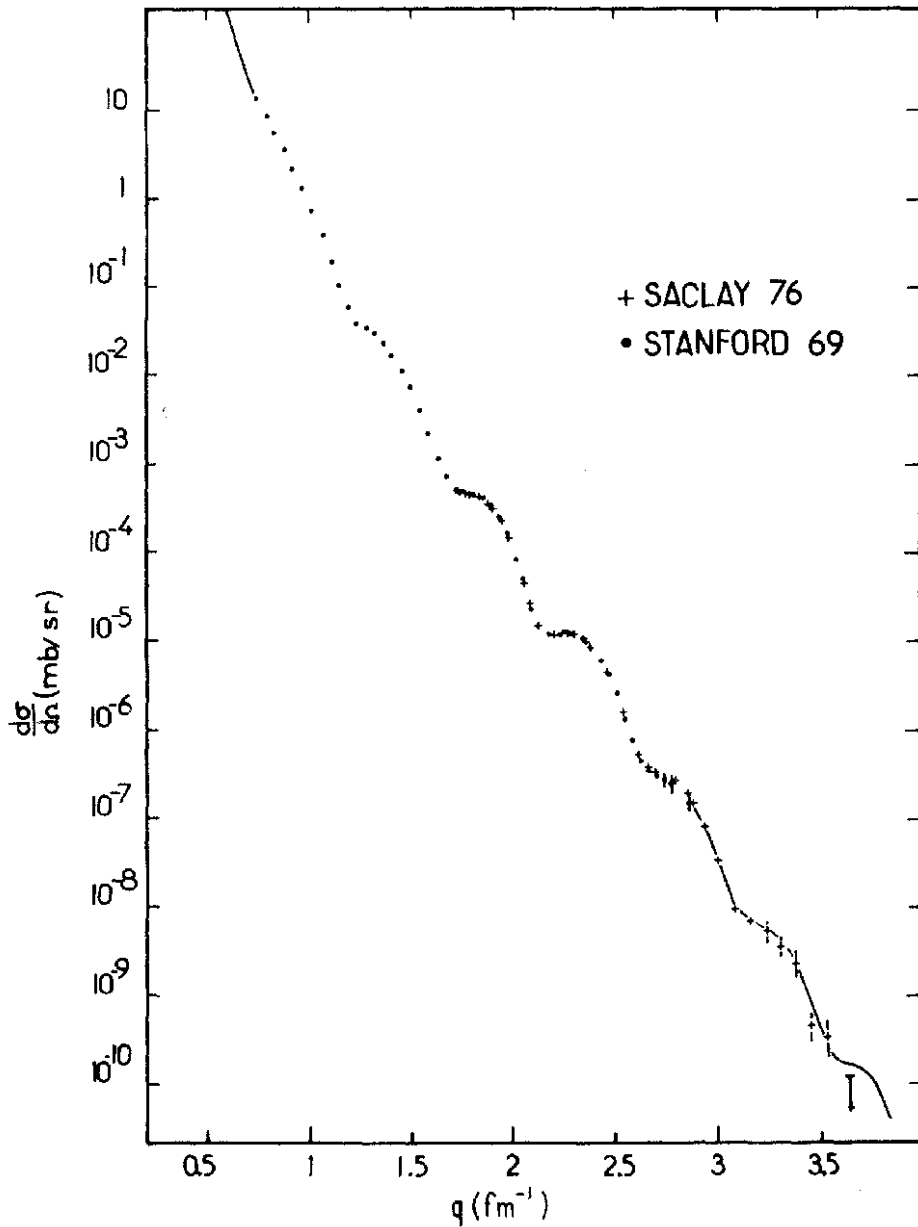


Supernova remnant  
Cassiopea A in X-rays

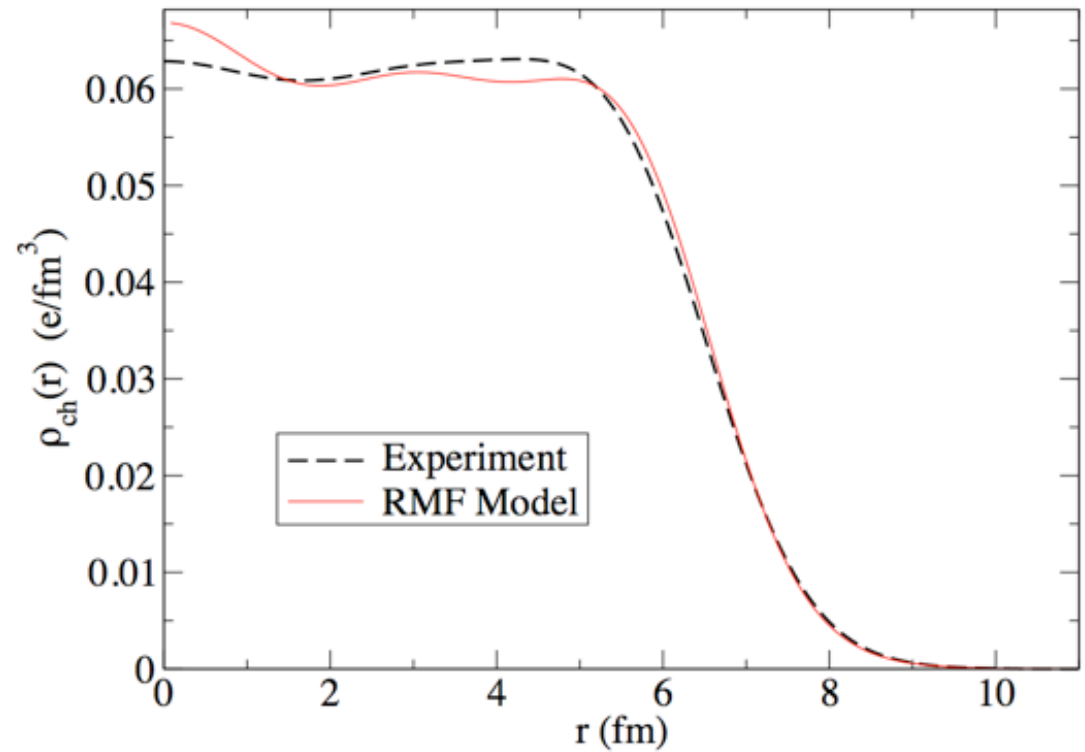


MD simulation of Nuclear  
Pasta with 100,000 nucleons

Charge Density of  $^{208}\text{Pb}$ , accurately measured in elastic electron scattering.

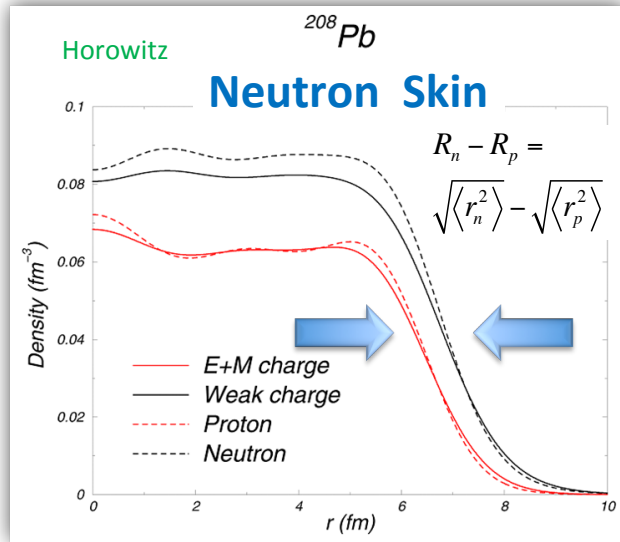


Cross section measured over 12 orders of magnitude.



These elastic charge densities **are** our picture of the atomic nucleus!

# Neutron Skin of Heavy Nuclei

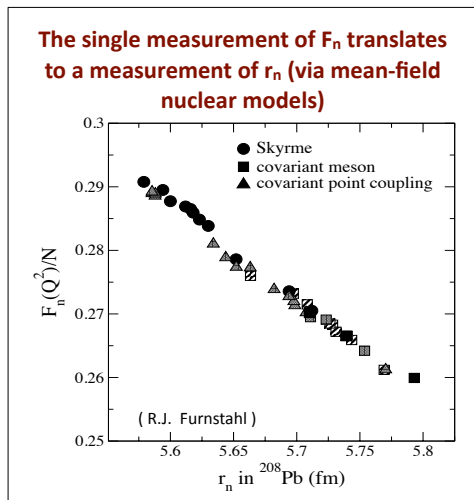


Nuclear theory predicts a neutron “skin” in heavy nuclei

Neutron distribution is not sensitive to the charge-sensitive photon

→ access through weak charge distribution

	proton	neutron
Electric charge	1	0
Weak charge	~0.08	1



For spin 0 nucleus:

$$A_{PV} = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \left[ \frac{F_n(Q^2)}{F_p(Q^2)} \right] F_{n,p}(Q^2) = \frac{1}{4\pi} \int d^3r j_0(qr) \rho_{n,p}(r)$$

- Measurement of  $R_n$  in <sup>208</sup>Pb **calibrates** the equation of state in neutron rich nuclear matter (determines density dependence of symmetry energy)
- Applications to neutron stars, heavy ion physics, atomic parity violation

# PREX/CREX

$$Q^2 \sim 0.01 \text{ GeV}^2$$

5° scattering angle

$$A_{PV} \sim 0.6 \text{ ppm}$$

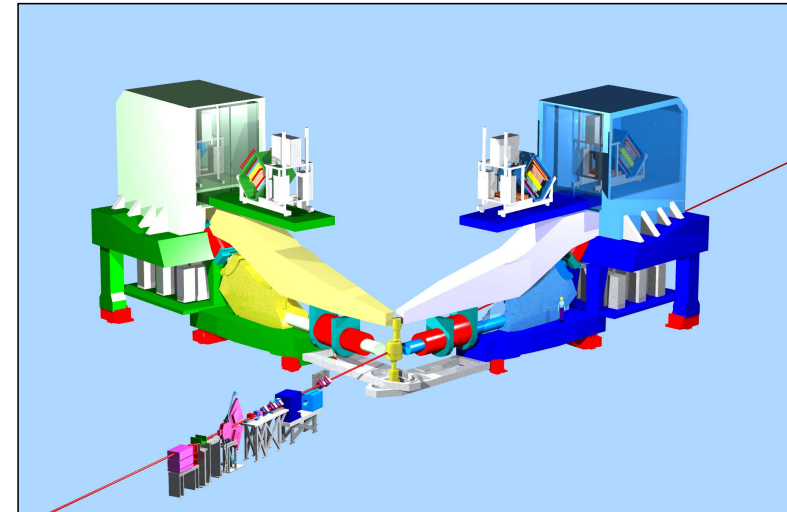
Rate  $\sim 1.5 \text{ GHz}$

**PREX-I (2012):**

$$A_{PV} = 0.657 \pm 0.060(\text{stat}) \pm 0.014(\text{sys}) \text{ ppm}$$

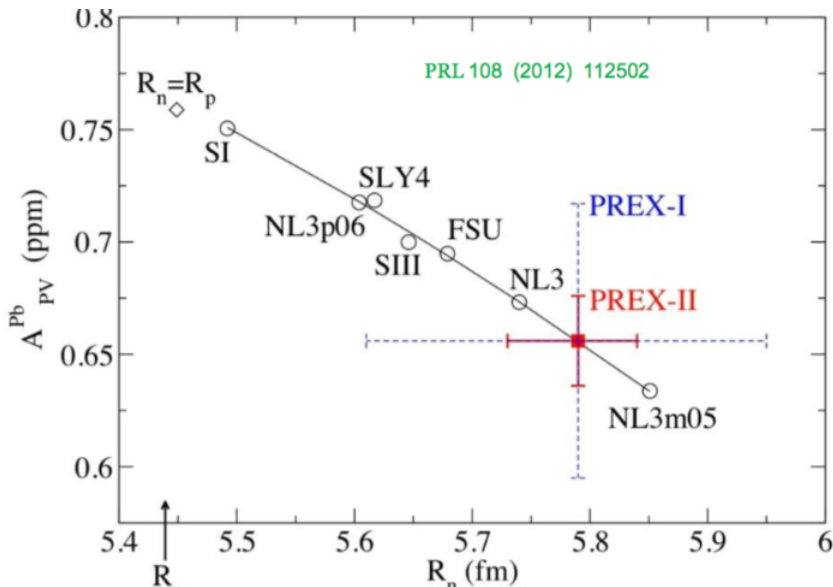
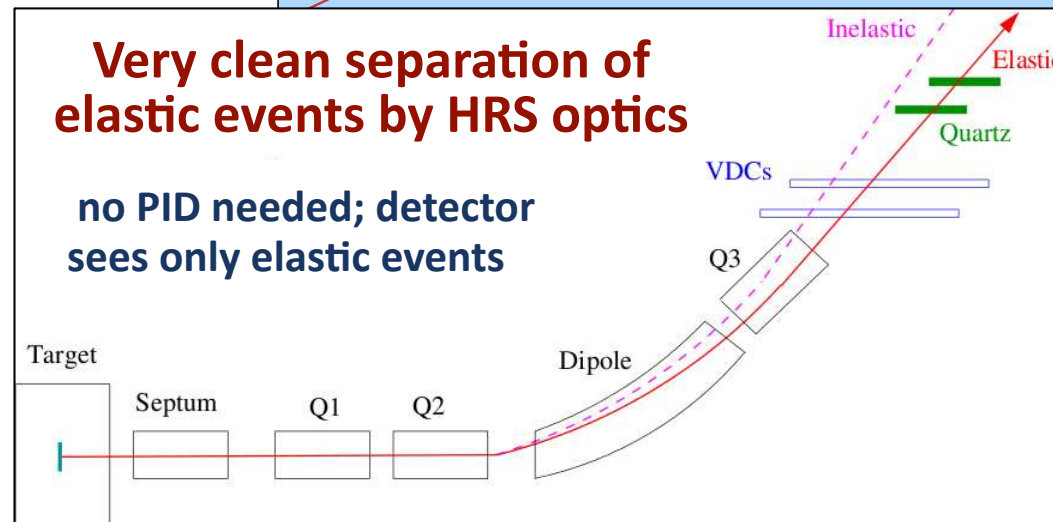
$$r_n - r_p = 0.33^{+0.16}_{-0.18} \text{ fm}$$

- 0.5 mm  $^{208}\text{Pb}$  foil, 70  $\mu\text{A}$
- 5° scattering
- $P_b \sim 90\% \pm 1\%$



**Very clean separation of elastic events by HRS optics**

**no PID needed; detector sees only elastic events**

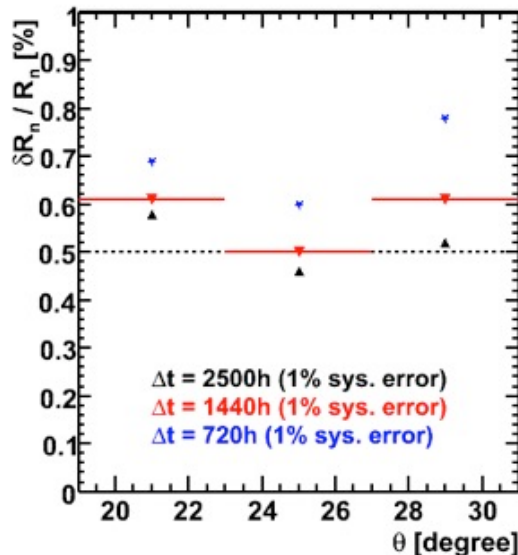
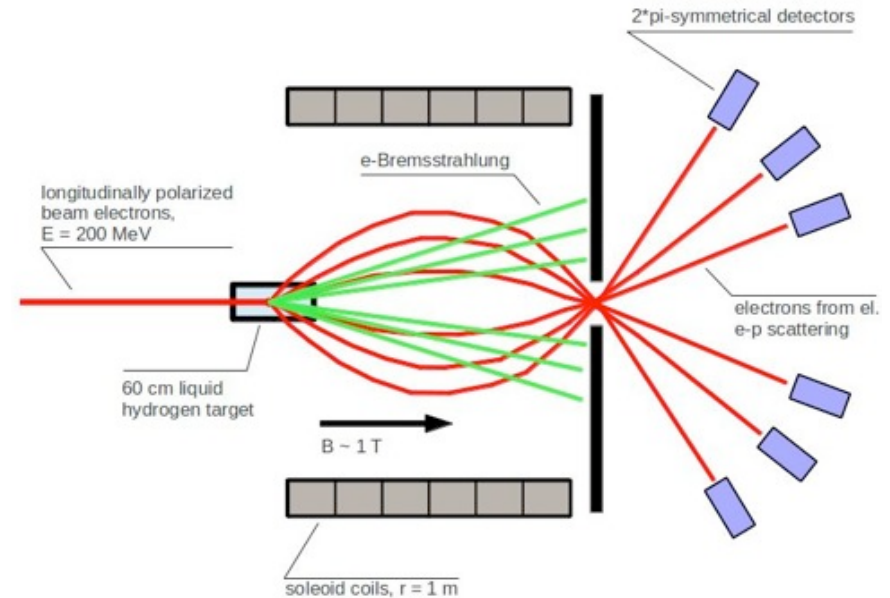
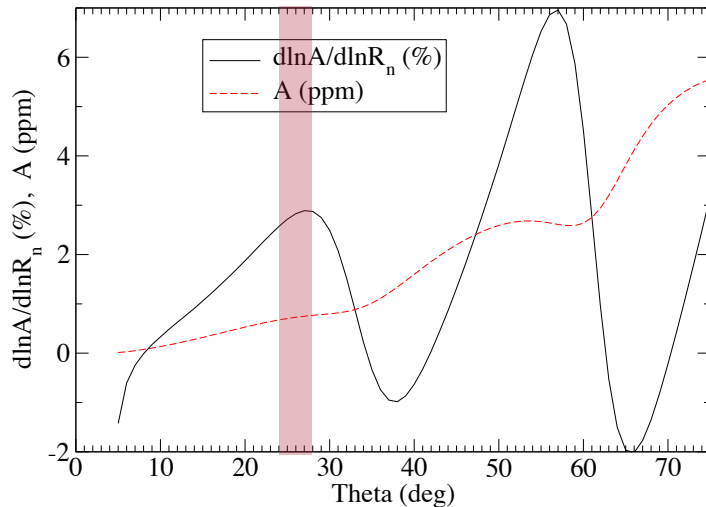


**Summer 2017: PREX (3%  $A_{PV}$ ,  $r_n$  to 0.06 fm), CREX (2.5%  $A_{PV}$ ,  $r_n$  to 0.02 fm)**

# Opportunities: “Super PREX”

## $N_{\text{skin}}$ measurement@MESA

### Sensitivity (C. Horowitz)



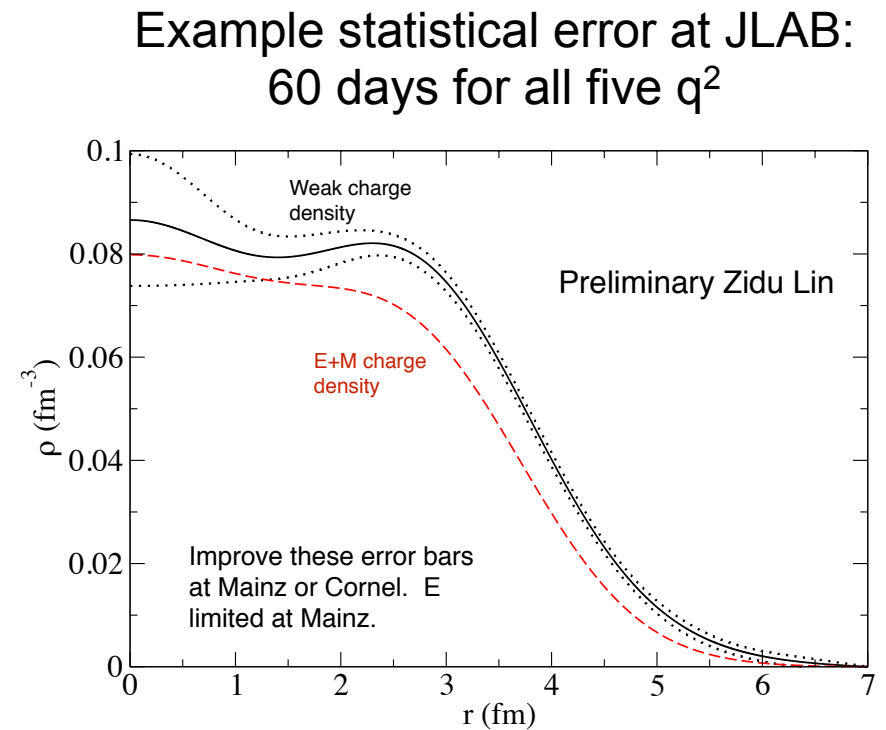
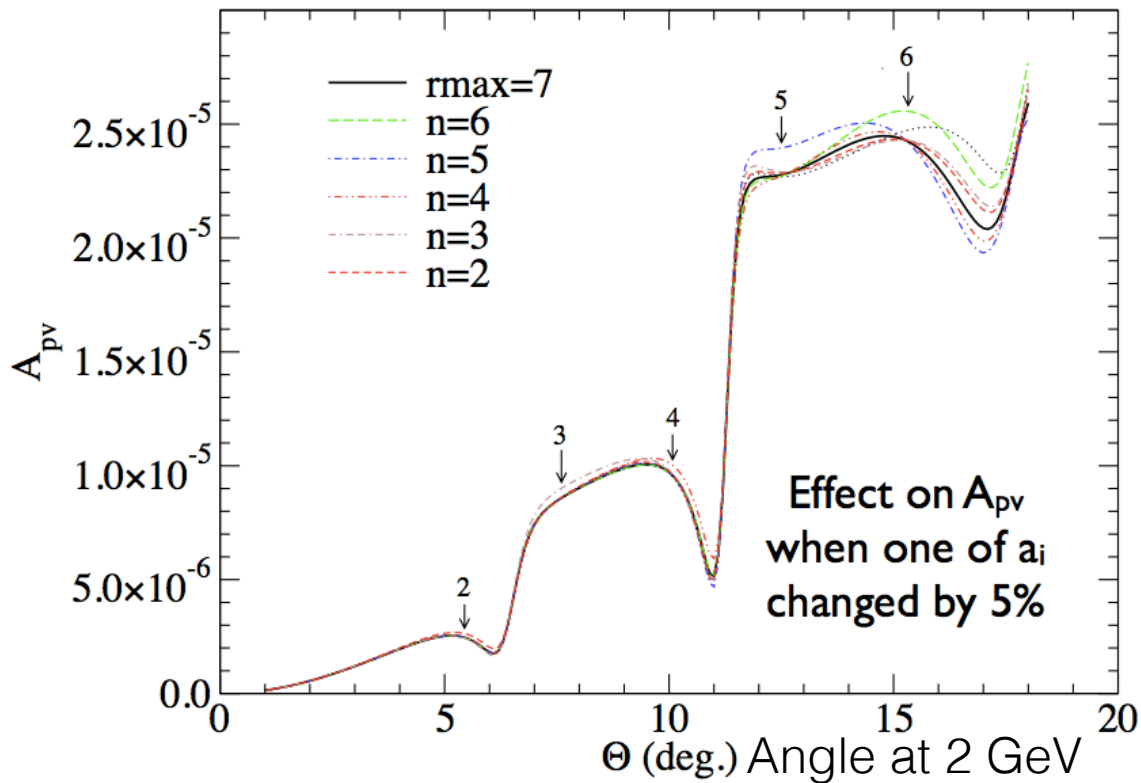
Same PREX Luminosity ( $0.25\text{mm } ^{208}\text{Pb}$ )  
 $\Delta\theta = 4^\circ$  : Rate = 9.75 GHz,  $A_{\text{PV}} = 0.68 \times 10^{-6}$

$1440\text{h} \rightarrow \delta A_{\text{PV}} / A_{\text{PV}} = 6.52 \times 10^{-3}$

$\rightarrow \delta R_n / R_n = 5.04 \times 10^{-3}$   
 (stat + syst 1%)

# Opportunities

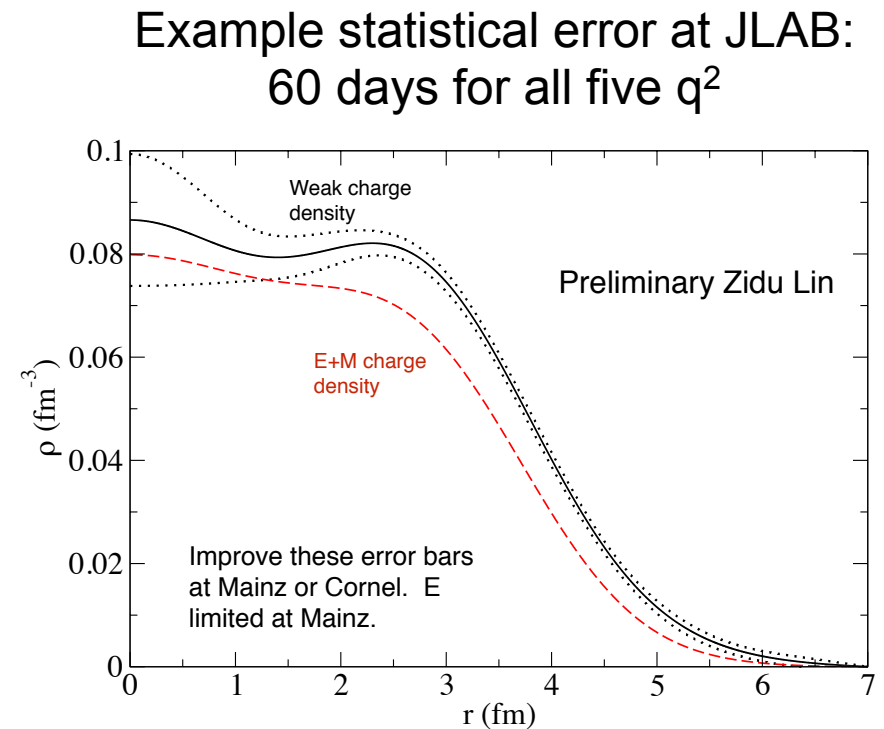
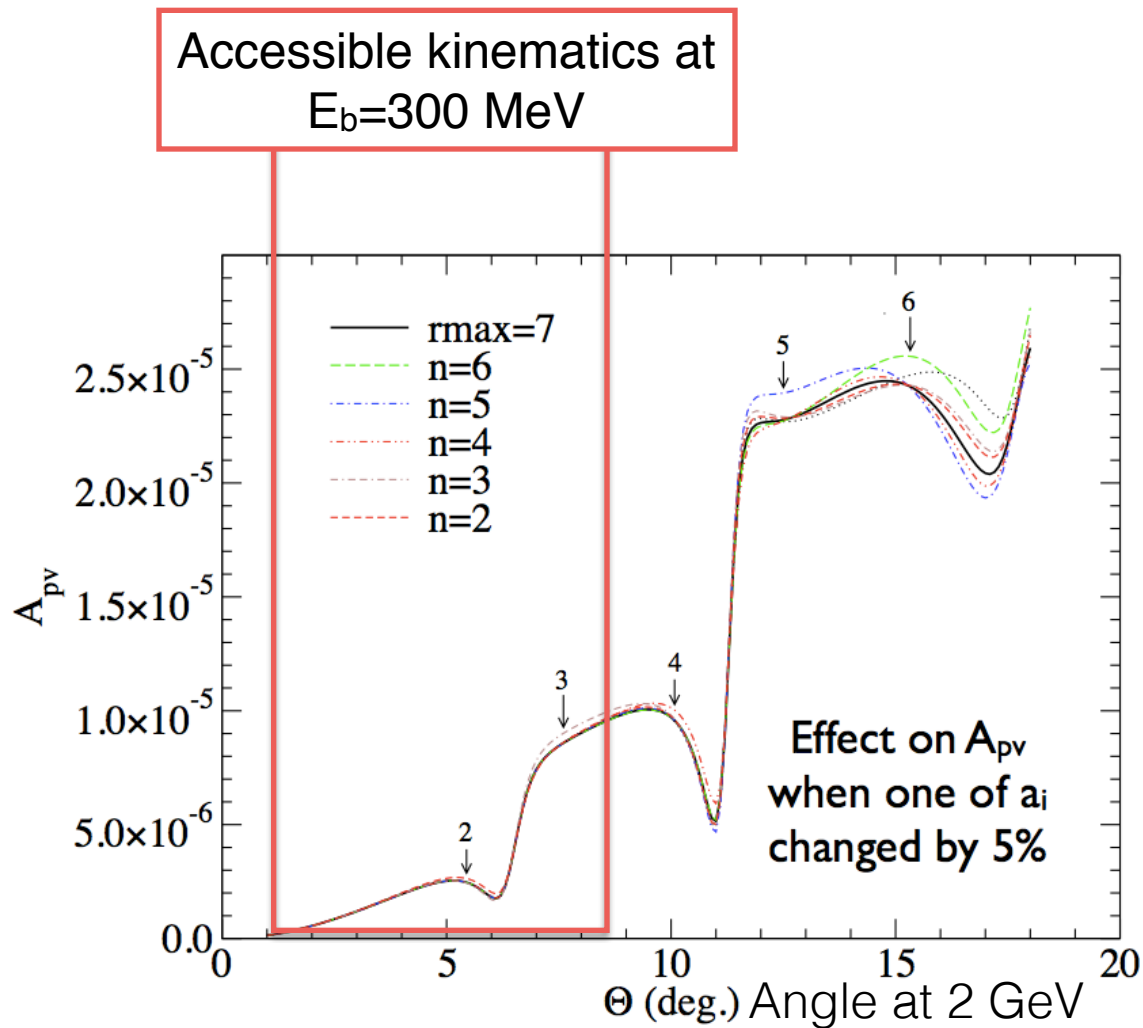
## Map neutron distribution of $^{48}\text{Ca}$





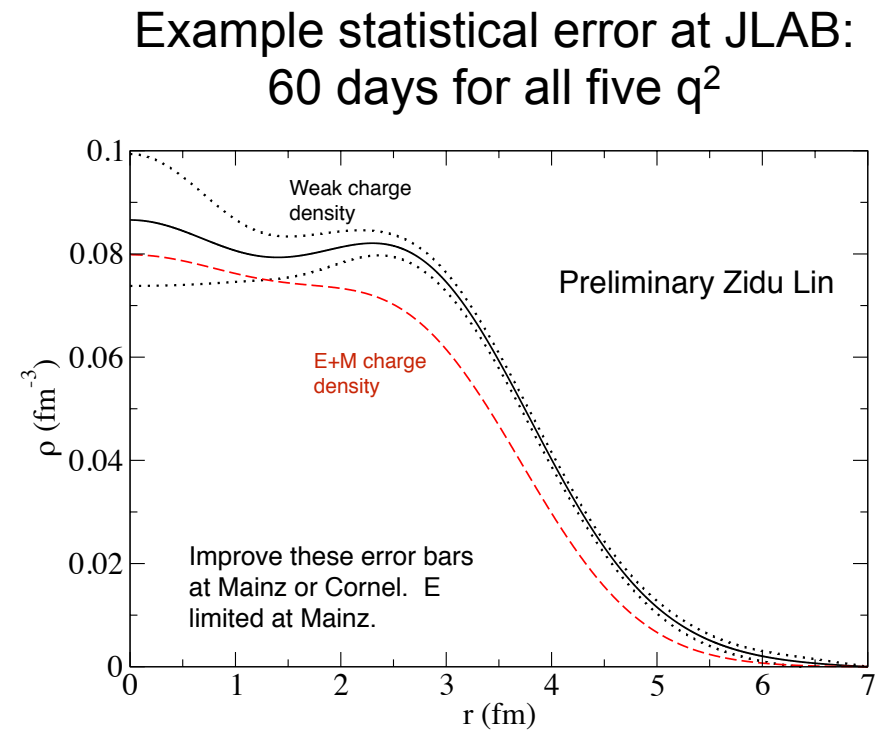
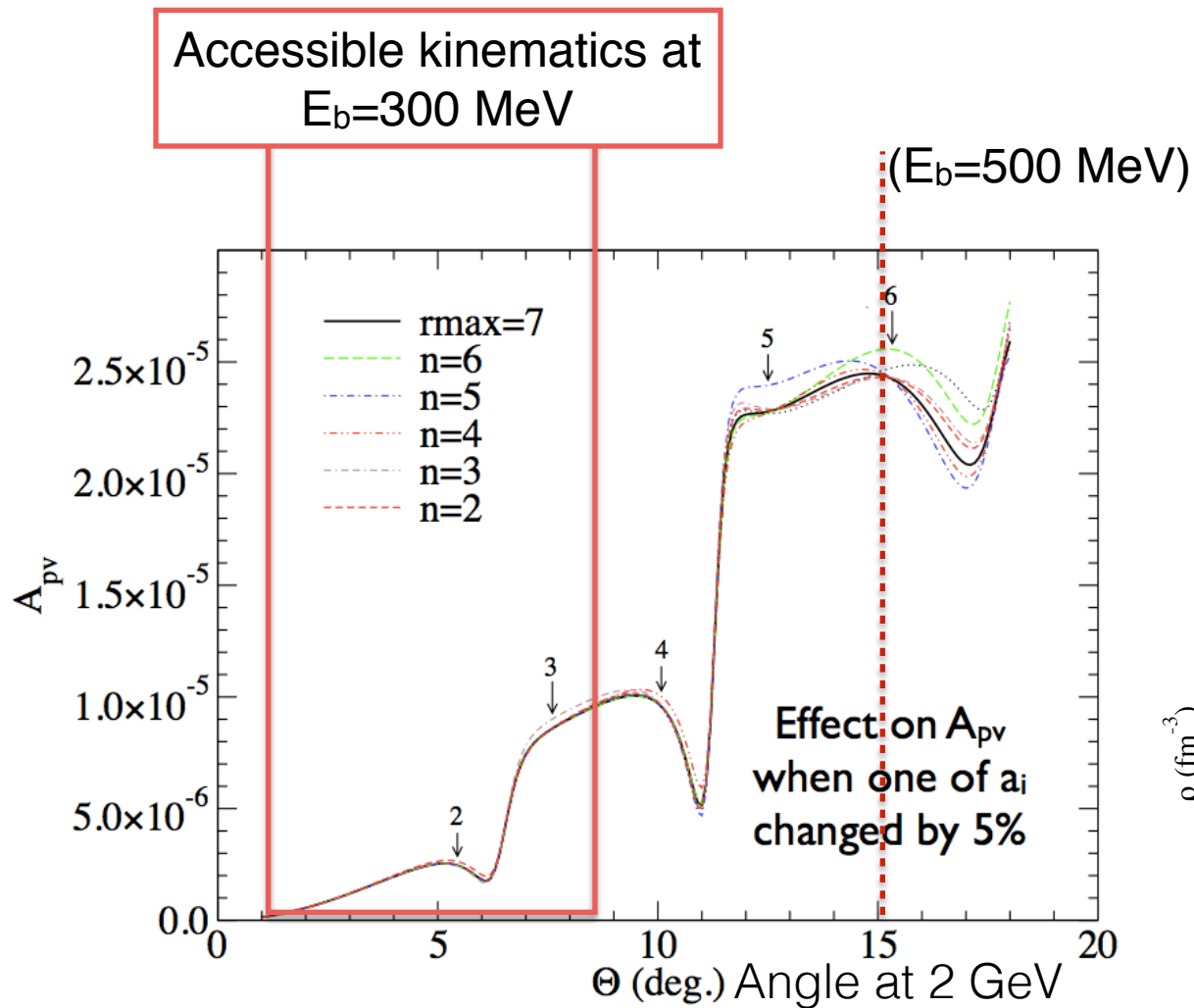
# Opportunities

## Map neutron distribution of $^{48}\text{Ca}$



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# Full $^{48}\text{Ca}$ weak charge density

- Would provide **text book picture of where neutrons and protons are in a nucleus.**
- Learn about shell oscillations of neutrons, saturation density of nuclear matter, neutron skin thickness, surface thickness of the neutrons...
- We expect central baryon density in  $^{208}\text{Pb}$  to be approximately constant but we only know what the proton density is.
- Compare to new microscopic calculations of the neutron density in  $^{48}\text{Ca}$  based on chiral effective field theory two and three nucleon interactions.

## Summary: Neutron distributions

- Crucial calibration on nuclear structure models
- “Super PREX” (also  $^{48}\text{Ca}$ ,  $^{128}\text{Sn}$ ? ~1000 hr each)
- Optimize program of neutron distribution measurements (this is sometime MESA cannot do)

# Precision Measurements To Date

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## ◆ Atomic Parity Violation

- ◆ future measurements and theory challenging

## ◆ Neutrino Deep Inelastic Scattering

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## ◆ PV Møller Scattering

- ◆ E158 at SLAC (total uncertainty 17 ppb)

- ◆ *statistics limited, theory robust*

***Technology developed to improve uncertainty by factor ~ 25***

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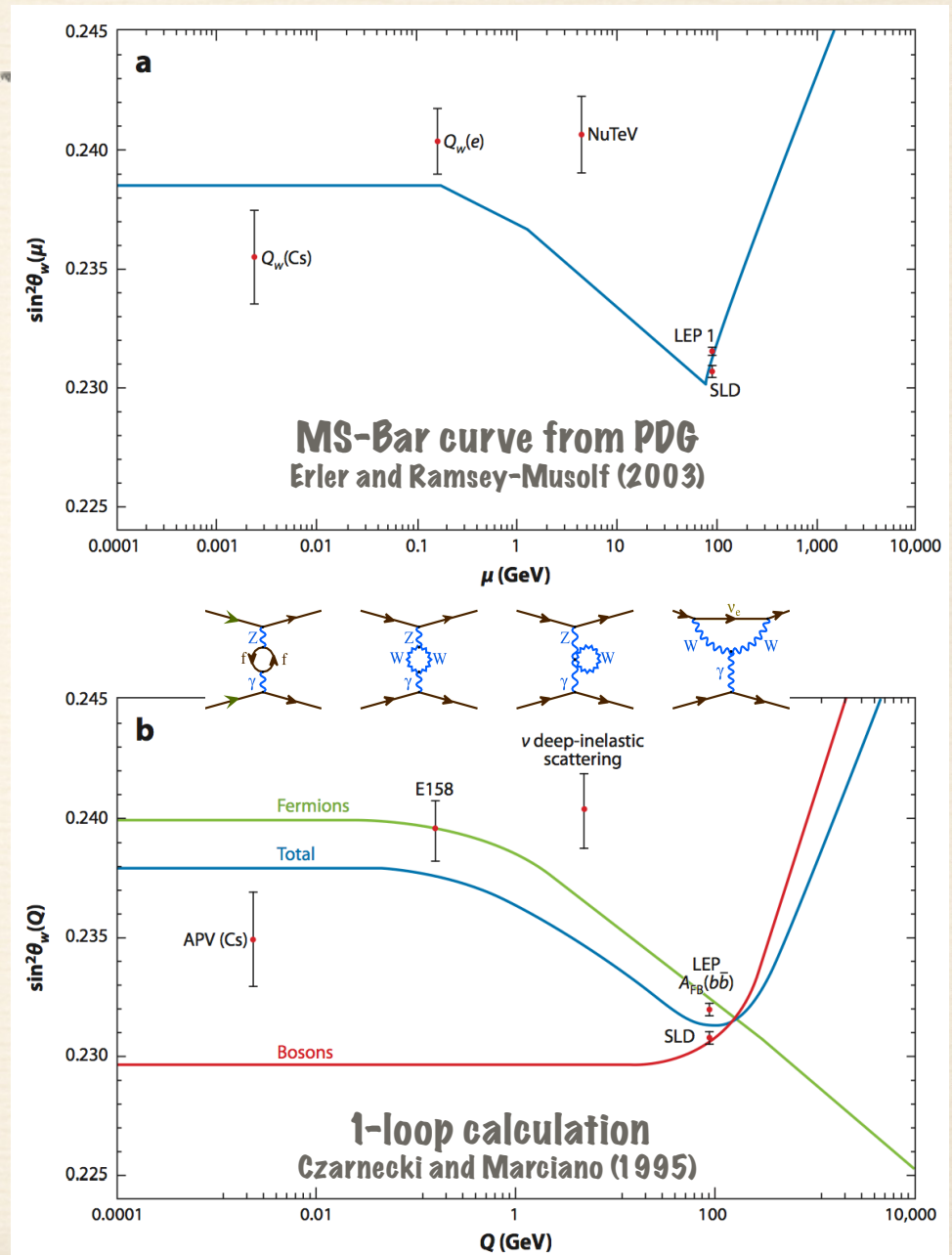
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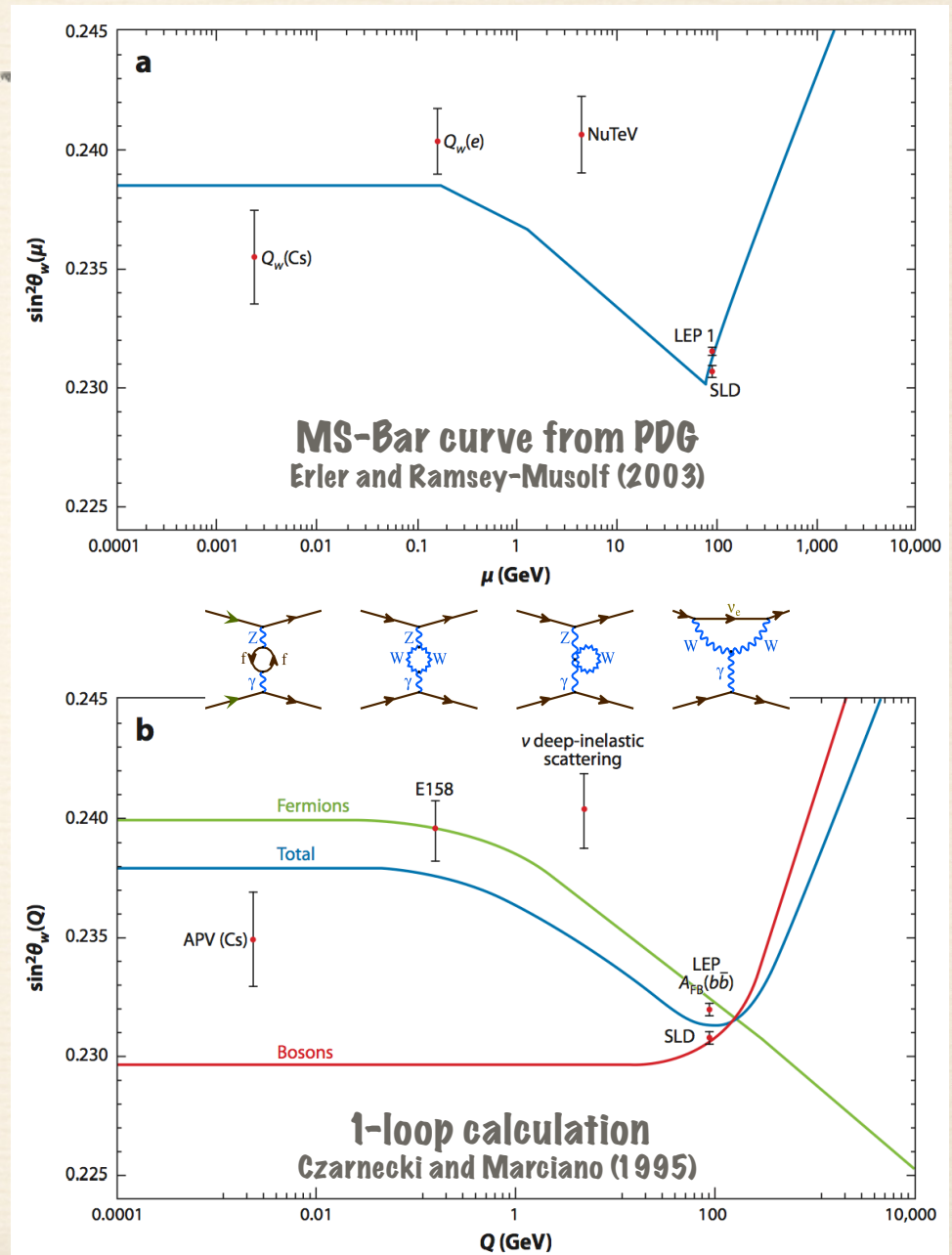
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### Recent Progress

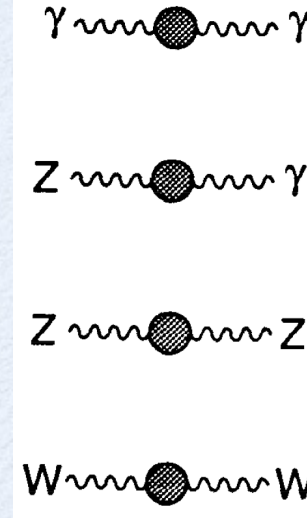
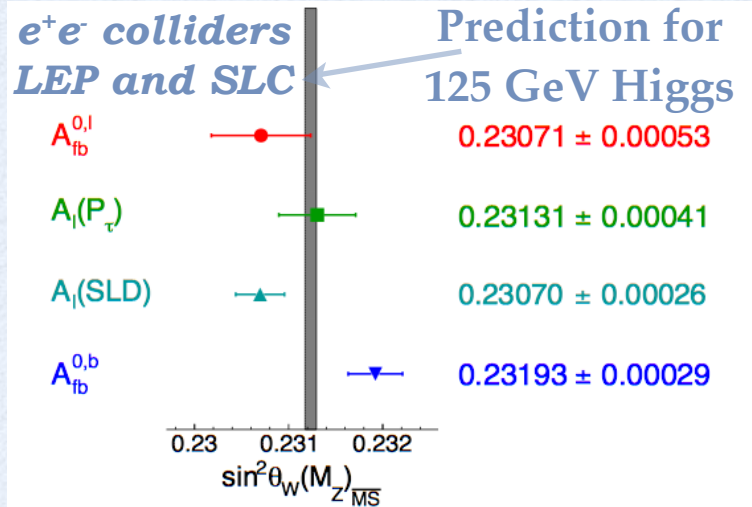
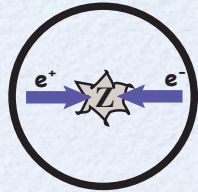
6 GeV PVDIS at JLab: first non-zero determination of axial-vector quark couplings

Qweak at JLab: should produce precision measurement soon



# Measurements of $\sin^2\theta_w$

*The most precise measurements are from LEP/SLC*



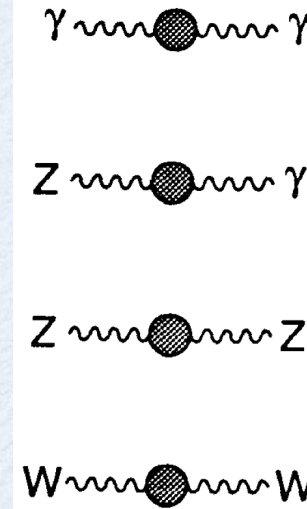
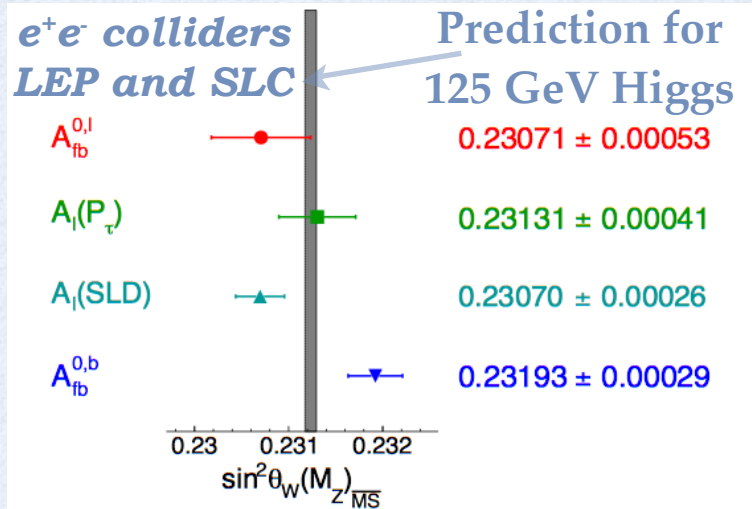
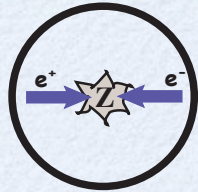
S, T, U  
parameters

Stringent constraints  
on large classes of  
new physics models



# Measurements of $\sin^2\theta_w$

The most precise measurements are from LEP/SLC



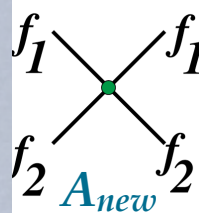
S, T, U  
parameters

Stringent constraints  
on large classes of  
new physics models

## Flavor Diagonal Contact Interactions

Consider  $f_1 \bar{f}_1 \rightarrow f_2 \bar{f}_2$  or  $f_1 f_2 \rightarrow f_1 f_2$

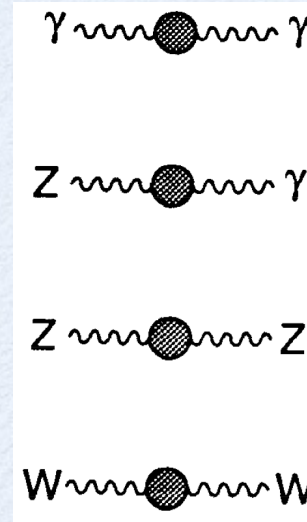
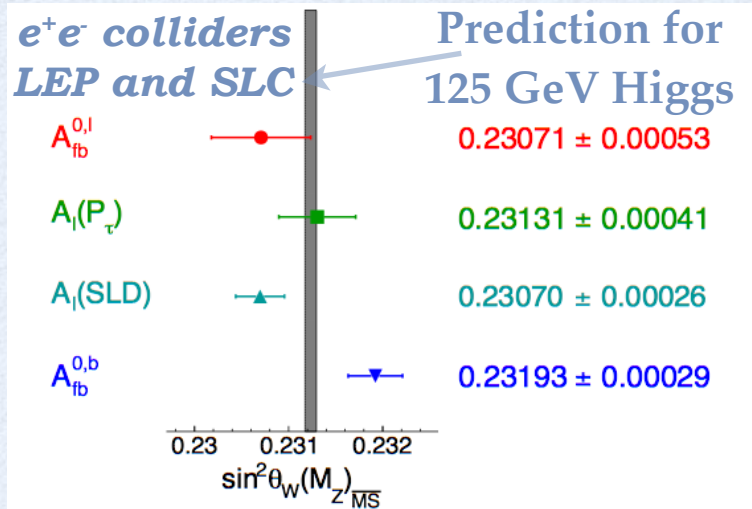
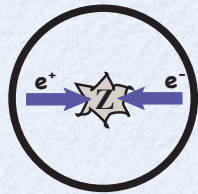
$$L_{f_1 f_2} = \sum_{i,j=L,R} \frac{4\pi}{\Lambda_{ij}^2} \eta_{ij} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma^\mu f_{2j}$$



**New heavy physics that does not  
couple directly to SM gauge bosons**

# Measurements of $\sin^2\theta_w$

The most precise measurements are from LEP/SLC



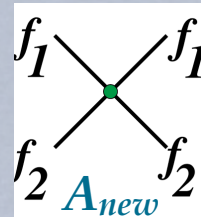
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## Flavor Diagonal Contact Interactions

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New heavy physics that does not  
couple directly to SM gauge bosons

on resonance:  $A_Z$  is imaginary

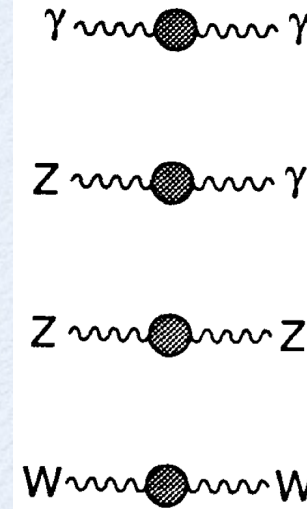
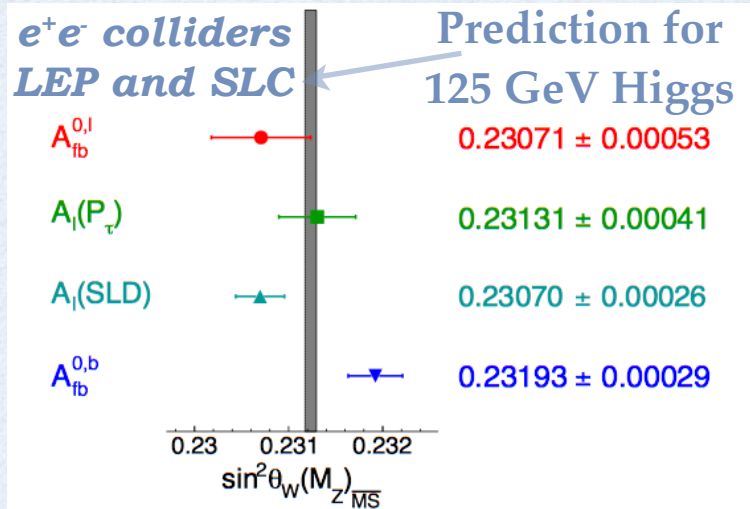
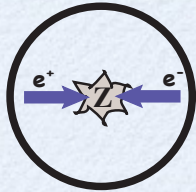
$$|A_Z + A_{\text{new}}|^2 \rightarrow A_Z^2 \left[ 1 + \left( \frac{A_{\text{new}}}{A_Z} \right)^2 \right]$$

no interference!

# Unique role for Low Energy Weak Neutral Current Measurements

## Measurements of $\sin^2\theta_w$

The most precise measurements are from LEP/SLC



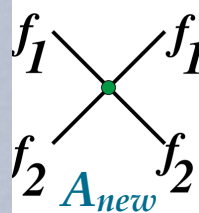
S, T, U  
parameters

Stringent constraints  
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new physics models

### Flavor Diagonal Contact Interactions

Consider  $f_1 \bar{f}_1 \rightarrow f_2 \bar{f}_2$  or  $f_1 f_2 \rightarrow f_1 f_2$

$$L_{f_1 f_2} = \sum_{i,j=L,R} \frac{4\pi}{\Lambda_{ij}^2} \eta_{ij} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma^\mu f_{2j}$$



New heavy physics that does not  
couple directly to SM gauge bosons

on resonance:  $A_Z$  is imaginary

$$|A_Z + A_{\text{new}}|^2 \rightarrow A_Z^2 \left[ 1 + \left( \frac{A_{\text{new}}}{A_Z} \right)^2 \right]$$

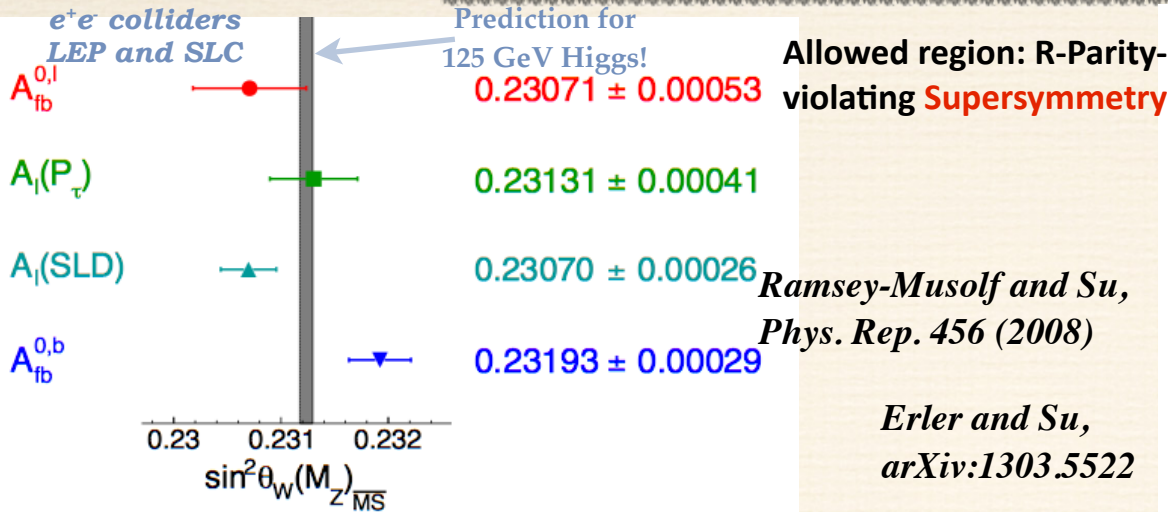
no interference!

New flavor diagonal interactions mediated by  
a new light boson such as the “dark Z”

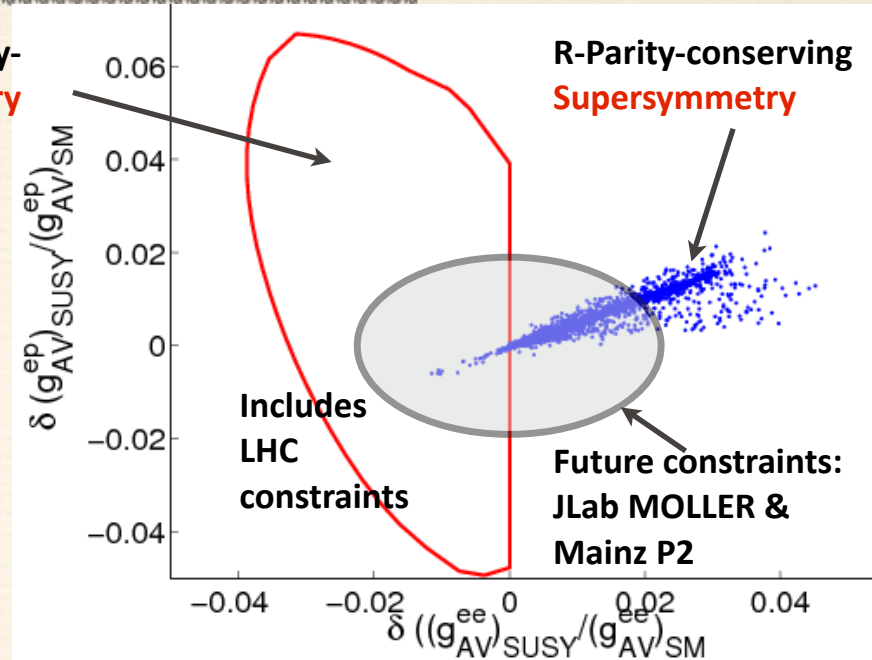
$$Q^2 \ll M_Z^2$$

# New Physics Complementarity

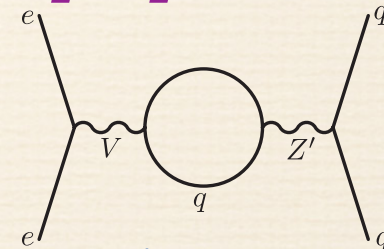
**Z resonance measurements: little sensitivity to new contact interactions**



MOLLER	—	proposed	$\pm 0.00029$
Qweak (Mainz)	—	proposed	$\pm 0.00037$
SOLID (JLab)	—	ongoing	$\pm 0.00060$
Qweak (JLab)	—	ongoing	$\pm 0.00072$
$A_{PV}^{Cs}$	●	published	$\pm 0.0014$
E158	●	published	$\pm 0.0014$



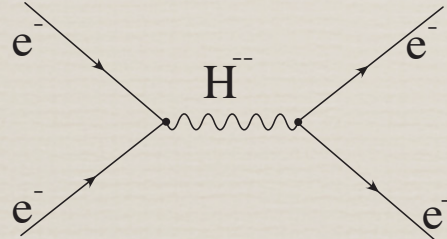
## Leptophobic Z'



**SOLID can improve sensitivity:  
100-200 GeV range**

## Lepton Number Violation

$\Lambda > 5 \text{ TeV}$   
**Doubly-  
Charged  
Scalars**



**Significant reach beyond LEP-200**

## Dark $Z$ and Parity Violation

- Low  $Q^2$  ( $< m_{Z_d}^2$ ) parity violation from  $Z - Z_d$  mixing
- $Z_d$  effects can be parameterized by HD, Lee, Marciano, 2012

$$G_F \rightarrow \rho_d G_F \quad \text{and} \quad \sin^2 \theta_W \rightarrow \kappa_d \sin^2 \theta_W$$

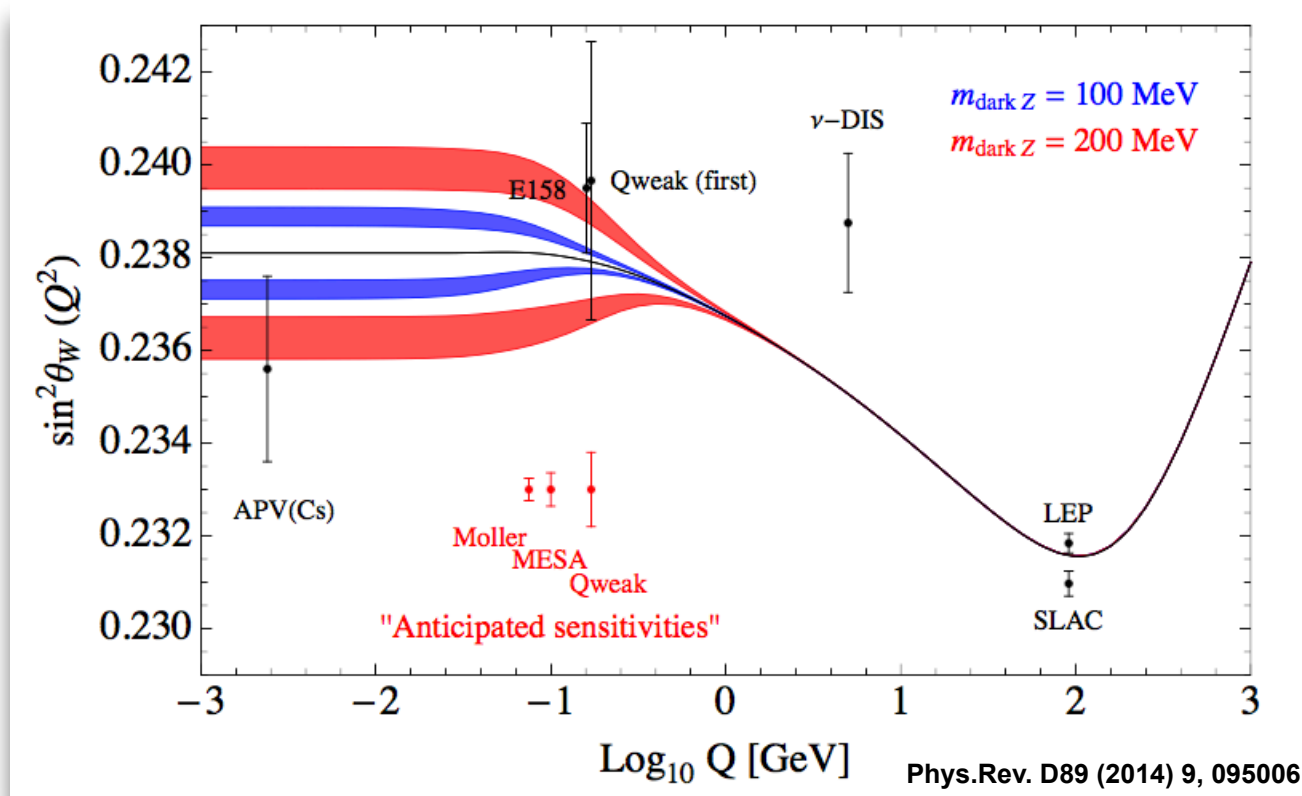
$$\text{with } \rho_d = 1 + \delta^2 \frac{m_{Z_d}^2}{Q^2 + m_{Z_d}^2} \quad \text{and} \quad \kappa_d = 1 - \varepsilon \frac{m_Z}{m_{Z_d}} \delta \frac{\cos \theta_W}{\sin \theta_W} \frac{m_{Z_d}^2}{Q^2 + m_{Z_d}^2}$$

- Leads to variation of  $\sin^2 \theta_W$  with  $Q^2$ :

$$\Delta \sin^2 \theta_W(Q^2) = -\varepsilon \delta \frac{m_Z}{m_{Z_d}} \sin \theta_W \cos \theta_W f(Q^2/m_{Z_d}^2)$$

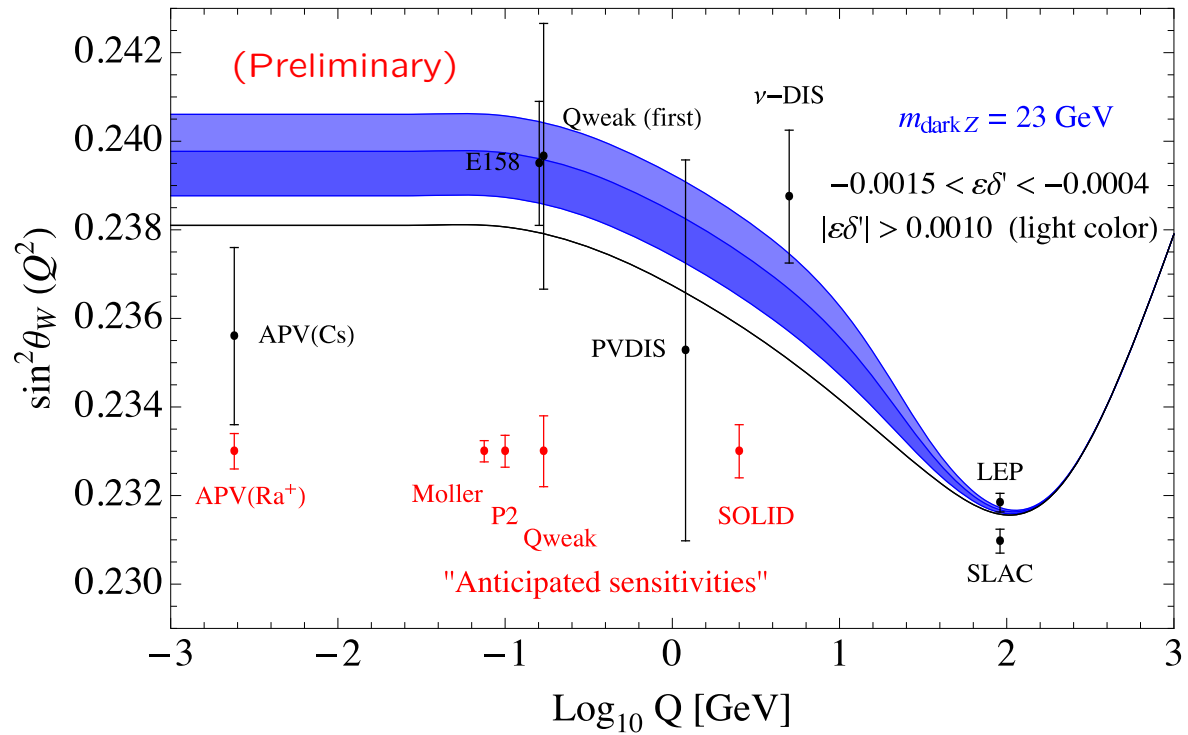
$$f(Q^2/m_{Z_d}^2) = 1/(1 + Q^2/m_{Z_d}^2)$$

Complementary “dark” U(1) search, not dependent on decay or production modes



light  $Z_d$   $Q^2$  dependent shift

HD, Lee, Marciano, work in progress



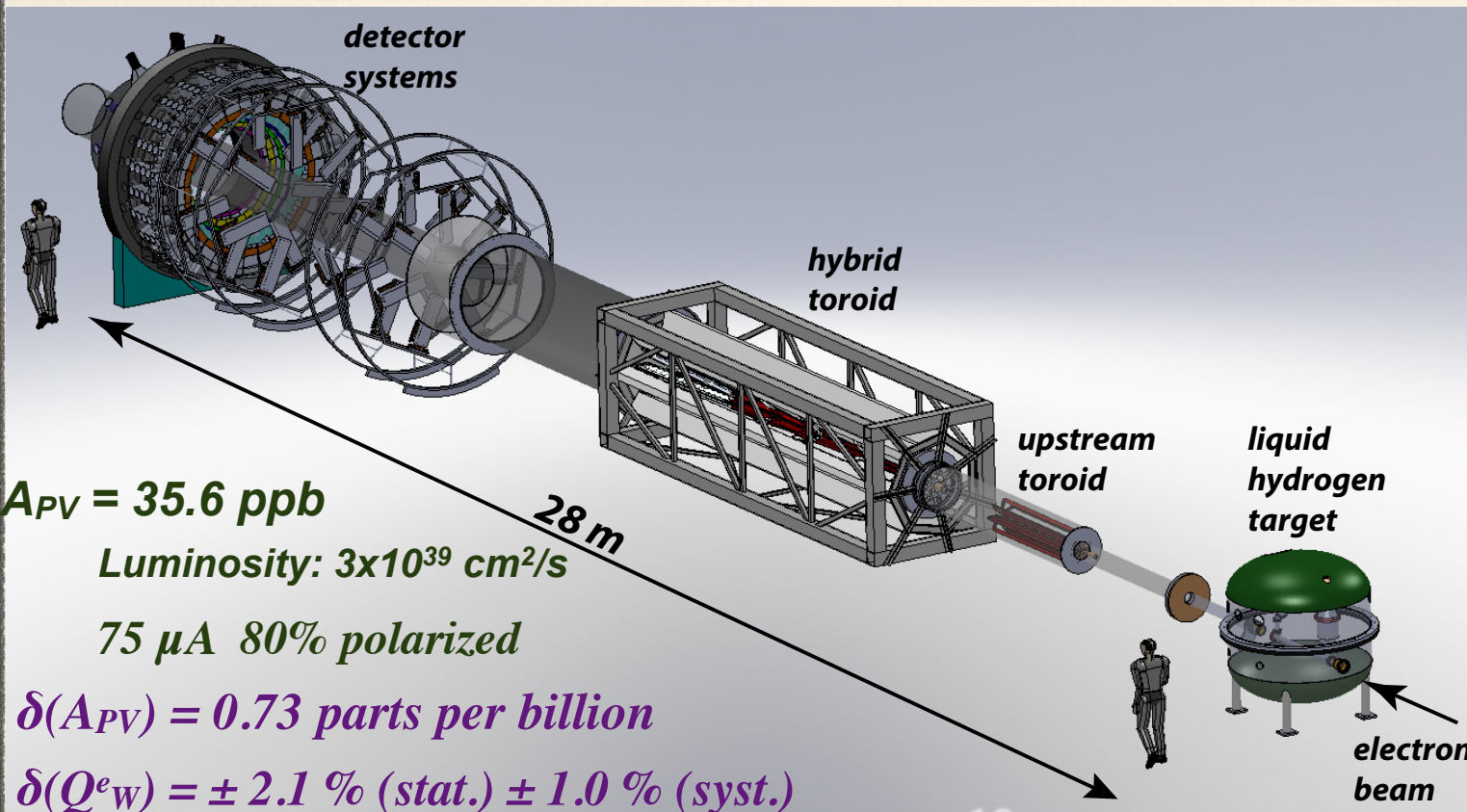
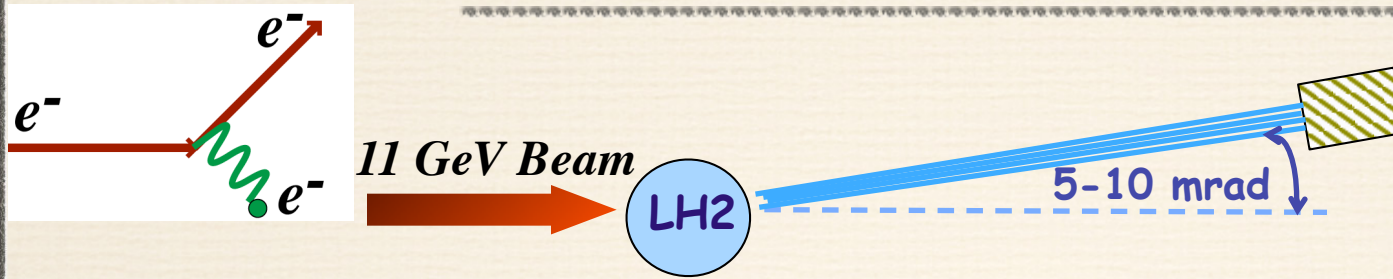
- $\varepsilon\delta' < 0$  range corresponds to  $1 \sigma$  band for  $\sin^2\theta_W$  deviation
- The upper region of the band: tension with constraints
- Interesting implications for planned experiments at different  $Q^2$
- Near future:  $Q_{\text{weak}}$  results can shed further light on this scenario

# An ultra-precise measurement of the weak mixing angle using Møller scattering

11 GeV Møller scattering

## MOLLER at JLab

### Measurement Of Lepton Lepton Electroweak Reaction



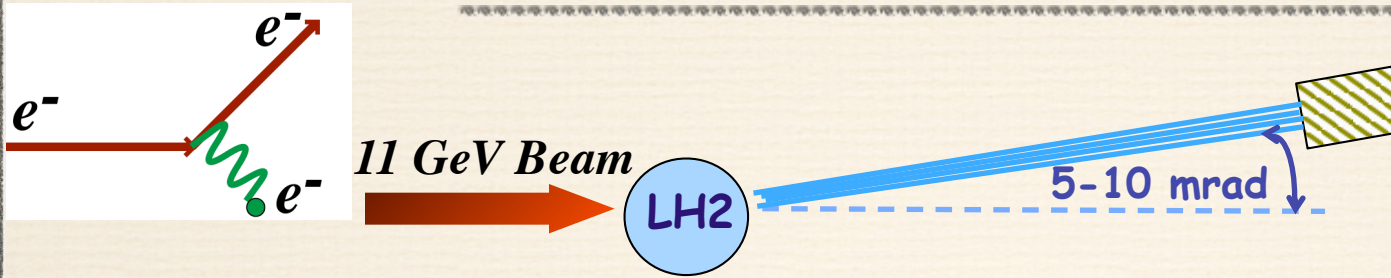


# An ultra-precise measurement of the weak mixing angle using Møller scattering

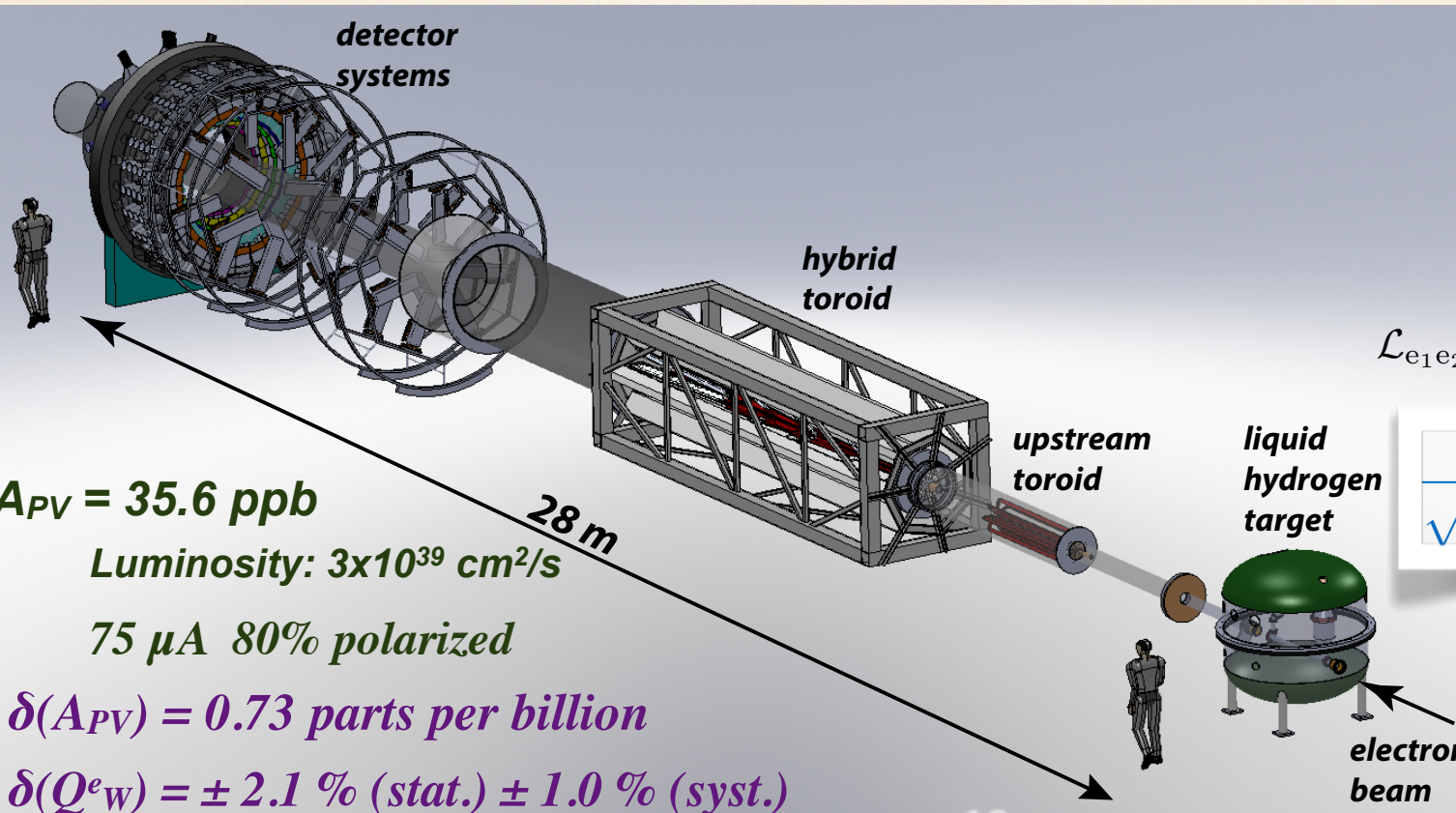
11 GeV Møller scattering

## MOLLER at JLab

Measurement Of Lepton Lepton Electroweak Reaction



$$Q_W = 1 - 4 \sin^2 \theta_W$$



$$+ \frac{1}{\Lambda^2} \mathcal{L}_6$$

$$\mathcal{L}_{e_1 e_2} = \sum_{i, j=L, R} \frac{g_{ij}^2}{2\Lambda^2} \bar{e}_i \gamma_\mu e_i \bar{e}_j \gamma^\mu e_j$$

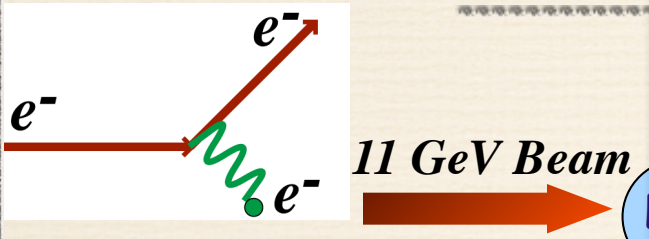
$$\frac{\Lambda}{\sqrt{|g_{RR}^2 - g_{LL}^2|}} = 7.5 \text{ TeV}$$

# An ultra-precise measurement of the weak mixing angle using Møller scattering

11 GeV Møller scattering

## MOLLER at JLab

Measurement Of Lepton Lepton Electroweak Reaction

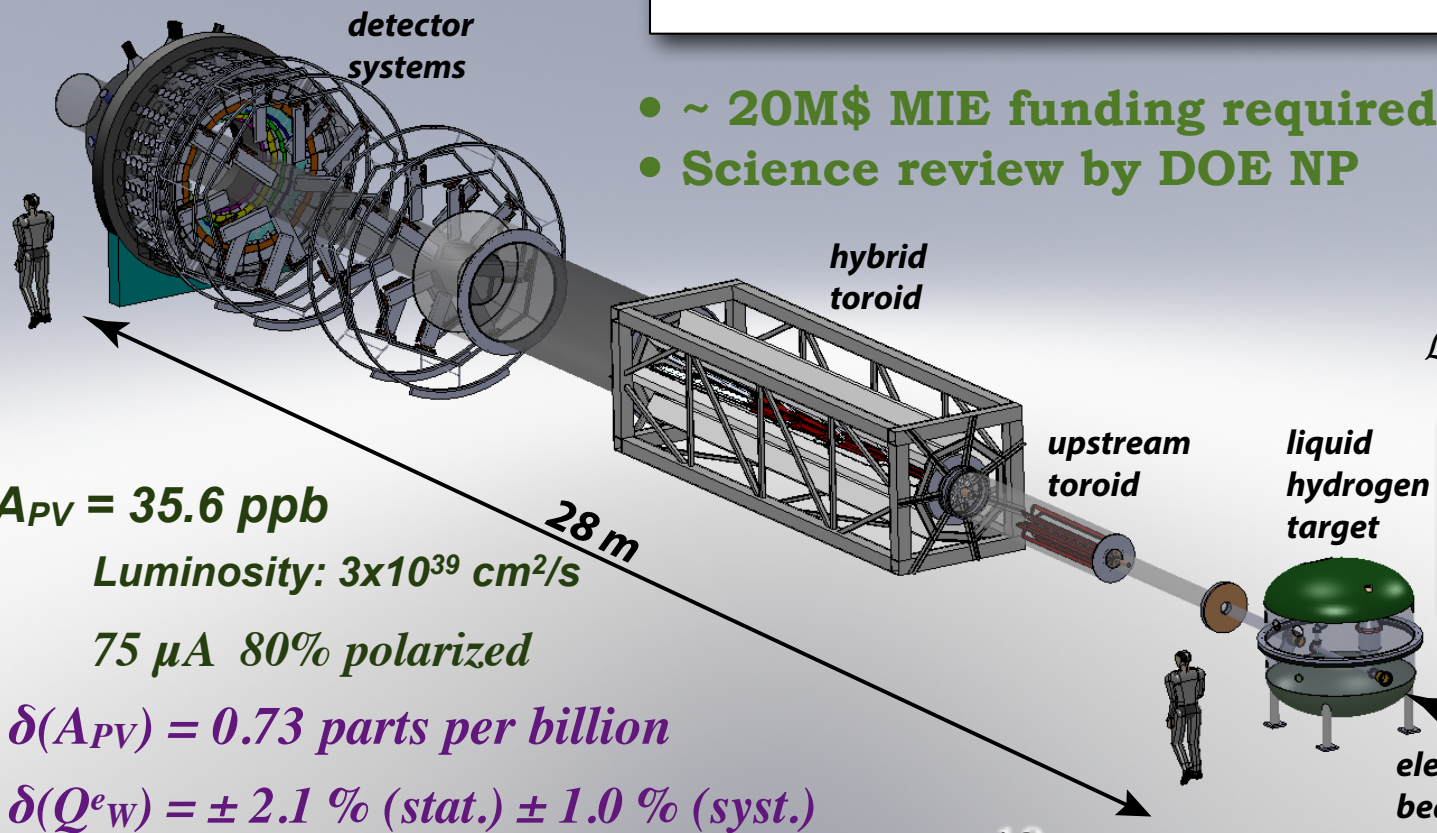


$$\delta(\sin^2\theta_W) = \pm 0.00026 \text{ (stat.)} \pm 0.00012 \text{ (syst.)} \rightarrow \sim 0.1\%$$

Matches best collider (Z-pole) measurements!

best contact interaction reach for leptons at low OR high energy

To do better for a 4-lepton contact interaction would require:  
Giga-Z factory, linear collider, neutrino factory or muon collider



- ~ 20M\$ MIE funding required
- Science review by DOE NP

$$+ \frac{1}{\Lambda^2} \mathcal{L}_6$$

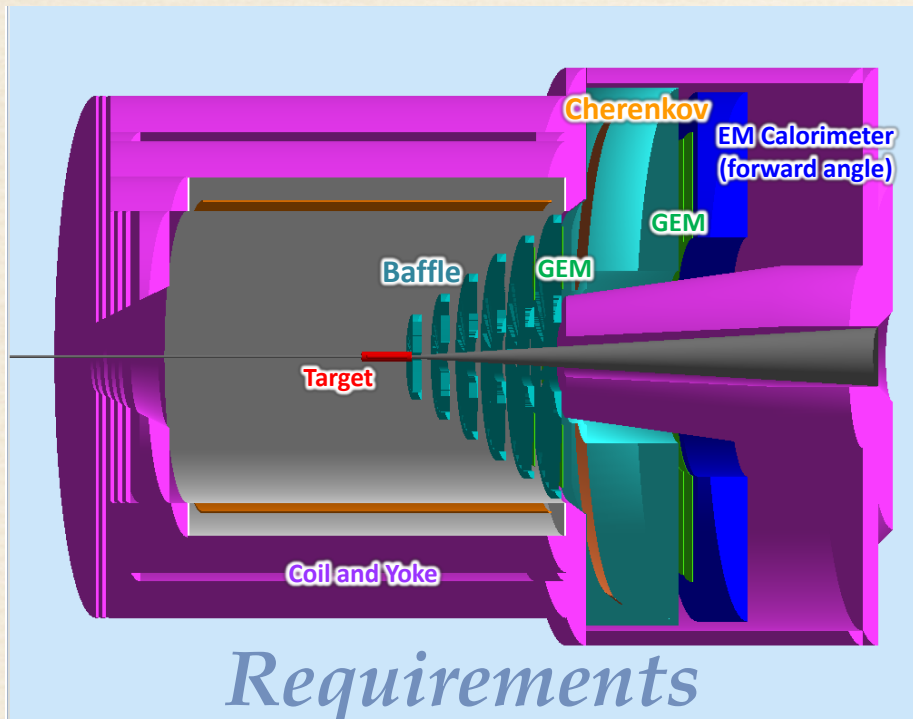
$$\mathcal{L}_{e_1 e_2} = \sum_{i,j=L,R} \frac{g_{ij}^2}{2\Lambda^2} \bar{e}_i \gamma_\mu e_i \bar{e}_j \gamma^\mu e_j$$

$$\frac{\Lambda}{\sqrt{|g_{RR}^2 - g_{LL}^2|}} = 7.5 \text{ TeV}$$

$A_{PV} = 35.6 \text{ ppb}$   
Luminosity:  $3 \times 10^{39} \text{ cm}^2/\text{s}$   
75  $\mu\text{A}$  80% polarized

$\delta(A_{PV}) = 0.73 \text{ parts per billion}$   
 $\delta(Q_e^W) = \pm 2.1 \% \text{ (stat.)} \pm 1.0 \% \text{ (syst.)}$

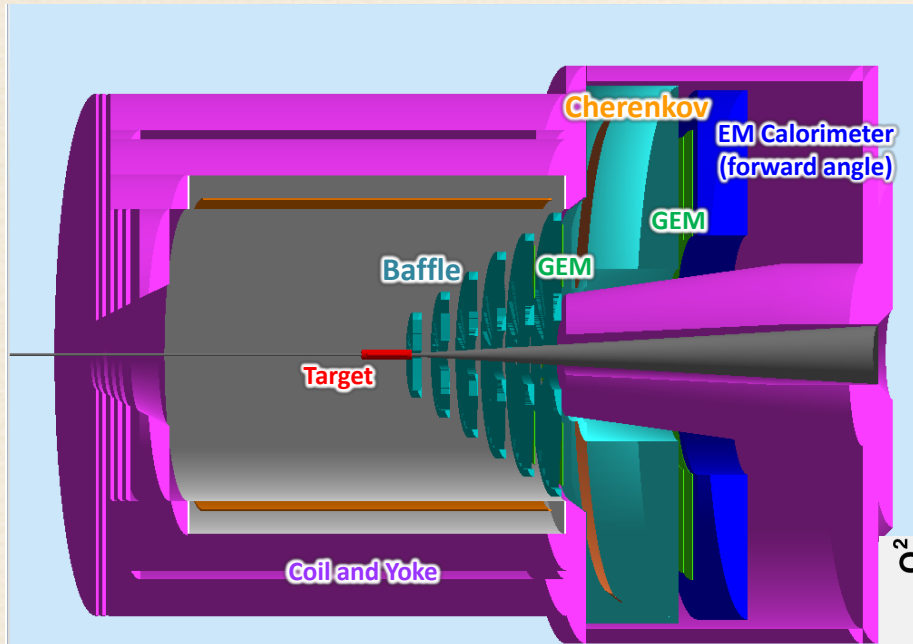
# SOLID with the 12 GeV Upgrade



Strategy: sub-1% precision over broad kinematic range: sensitive Standard Model test *and* detailed study of hadronic structure contributions

- *High Luminosity with  $E > 10$  GeV*
- *Large scattering angles (for high  $x$  &  $y$ )*
- *Better than 1% errors for small bins*
- *$x$ -range 0.25-0.75*
- *$W^2 > 4$  GeV<sup>2</sup>*
- *$Q^2$  range a factor of 2 for each  $x$* 
  - (Except at very high  $x$ )
- *Moderate running times*

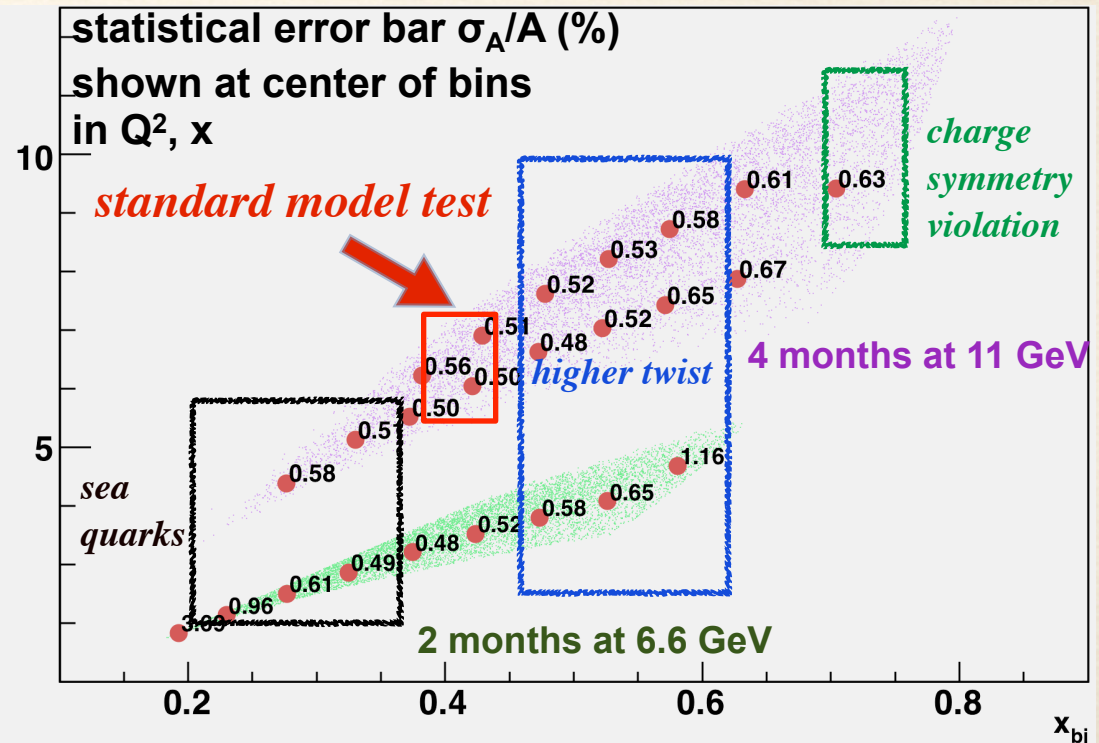
# SOLID with the 12 GeV Upgrade



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**Strategy:** sub-1% precision over broad kinematic range: sensitive Standard Model test and detailed study of hadronic structure contributions



# JLab Qweak: First Measurement of Proton's Weak Charge

$$Q_W^p \equiv -2[2C_{1u} + C_{1d}] = (1 - 4\sin^2 \theta_W)$$

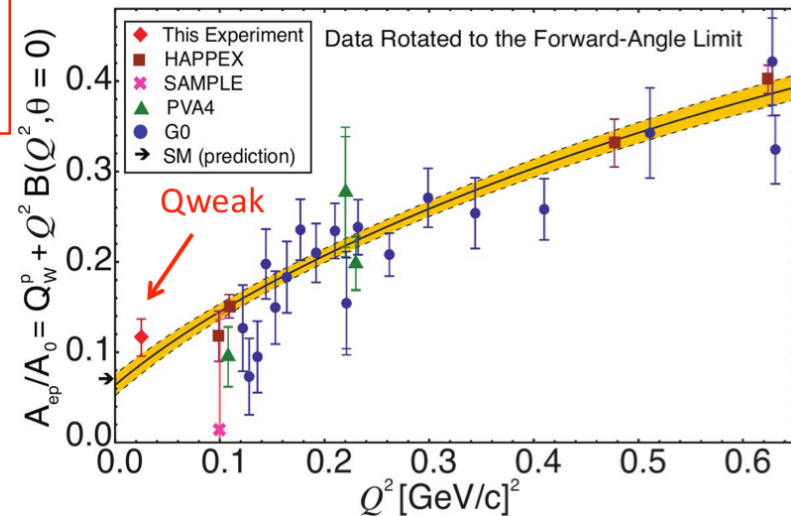
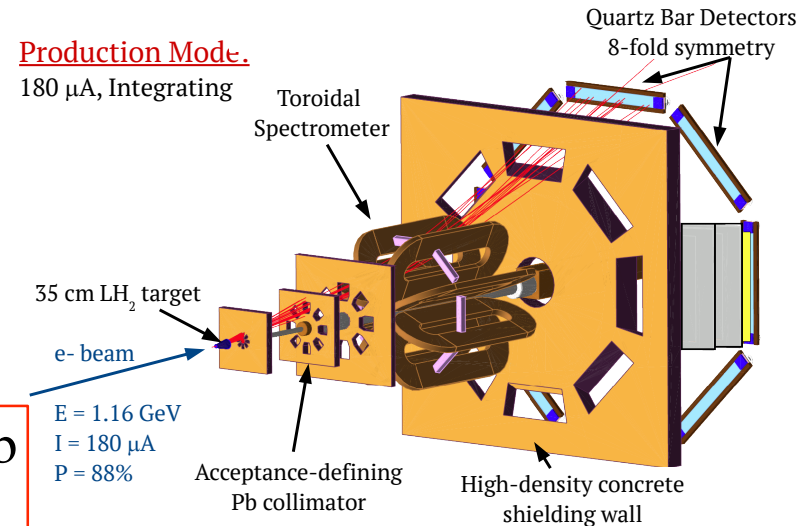
$$A \sim \left[ \frac{-G_F}{4\pi\alpha\sqrt{2}} \right] \left[ Q^2 Q_{weak}^p + Q^4 B(Q^2) \right]$$

(1/25<sup>th</sup> of total dataset) – published in PRL **111**, 141803 (2013)

$$A_{ep} = -279 \pm 35(\text{stat}) \pm 31(\text{syst}) \text{ ppb}$$

$$\langle Q^2 \rangle = 0.0250 \text{ (GeV/c)}^2$$

$$Q_W^p \text{ (PVES)} = 0.064 \pm 0.012$$

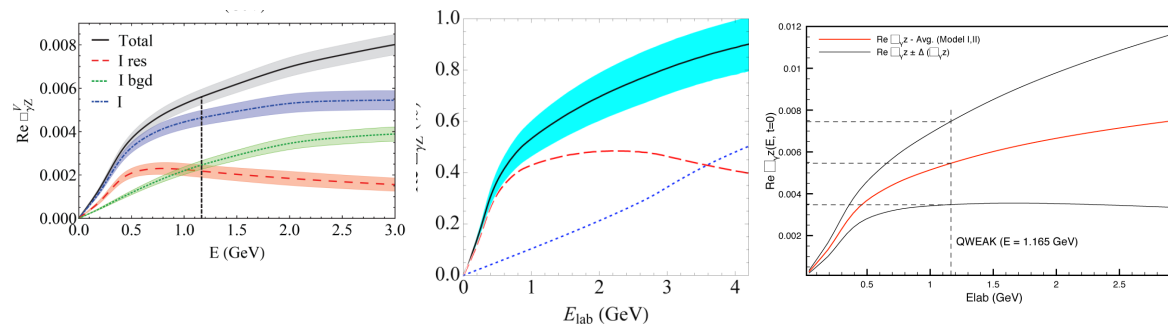


# Vector box plots today

Hall *et al.*  
PRD 88, 013011 (2013)

Carlson and Rislow  
PRD 83, 113007 (2011)

Gorchtein *et al.*  
PRC 84, 015502 (2011)



$\text{Re} T_{\gamma Z}^V(E = 1.165 \text{ GeV})$		
$(5.6 \pm 0.36) \times 10^{-3}$	$(5.7 \pm 0.9) \times 10^{-3}$	$(5.4 \pm 2.0) \times 10^{-3}$

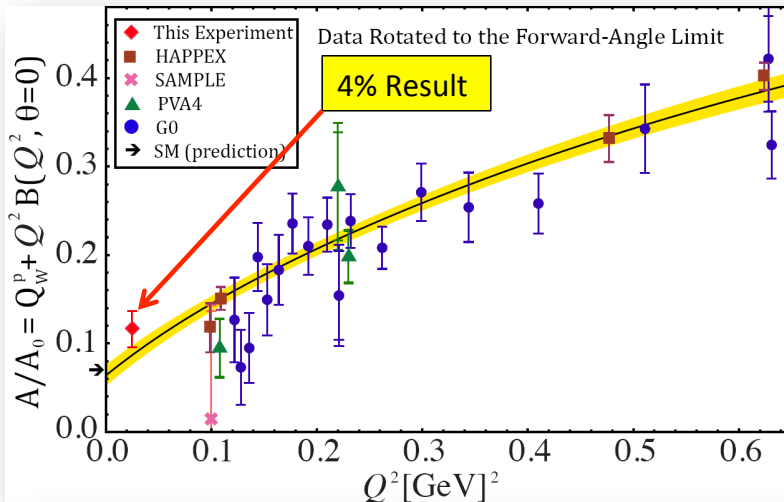
- Central values close
- Differences come from the treatment of the structure functions
- BTW, we combined errors directly, Hall et al. in quadrature. Could repeat:

$\text{Re} T_{\gamma Z}^V(E = 1.165 \text{ GeV})$		
$(5.6 \pm 0.36) \times 10^{-3}$	$(5.7 \pm 0.52) \times 10^{-3}$	$(5.4 \pm 2.0) \times 10^{-3}$

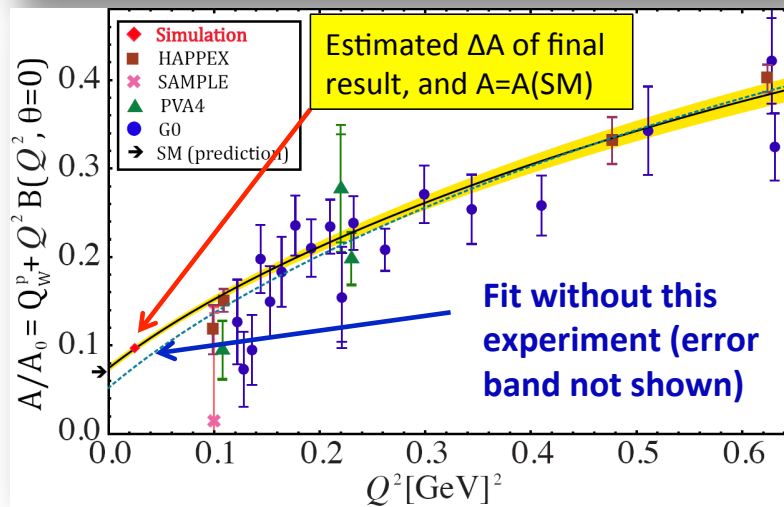
# Summary

- The world is saved—maybe—regarding the  $\gamma Z$  corr. to  $Q_{\text{Weak}}$ .
- I.e.,  $\square_{\gamma Z}^V$  now calculated.
- About  $(8.1 \pm 1.4)\%$  of  $Q_W^P$  at  $E_{\text{elec}} = 1.165$  GeV. Proportional to  $E_{\text{elec}}$ .
- Not discussed here:  $\square_{\gamma Z}^A$  also now calculated w/o guesswork certain log terms
- About  $(6.3 \pm 0.6)\%$  of  $Q_W^P$  at  $E_{\text{elec}}$  threshold. Small dependence on  $E_{\text{elec}}$ . Might still like to improve.
- For goal of 1% or better measurement of  $Q_{\text{Weak}}$  (Mesa), energy is about 1/6 of JLab experiment, and corrections and error in  $\square_{\gamma Z}^V$  scale with energy.
- PVDIS can help shrink uncertainty limits.

# Global fit of $Q^2 < 0.63$ (GeV/c)<sup>2</sup> PVES Data



**$A = -279 \pm 35 \pm 31$  ppb**  
 **$Q_W(p) = 0.064 \pm 0.012$**   
**(only 1/25<sup>th</sup> of all data collected)**  
**SM value = 0.0710(7)**



**Estimated Fit Uncertainties with Final Result (Assuming SM Value)**



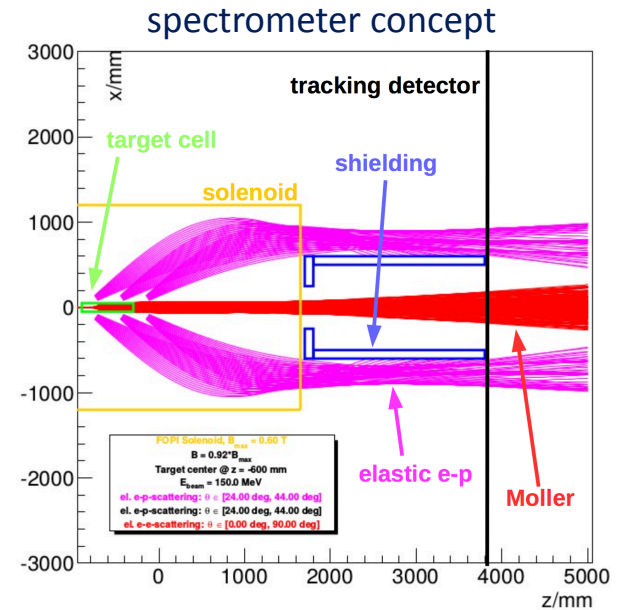
# P2 at Mainz MESA – Proton Weak Charge

$$\bar{e} + p \rightarrow e' + p \quad Q_W^p \equiv -2[2C_{1u} + C_{1d}] = (1 - 4\sin^2 \theta_W)$$

$$A \sim \left[ \frac{-G_F}{4\pi\alpha\sqrt{2}} \right] \left[ Q^2 Q_{weak}^p + Q^4 B(Q^2) \right]$$

Run at low energy; reduce hadronic contributions and gamma-Z box radiative

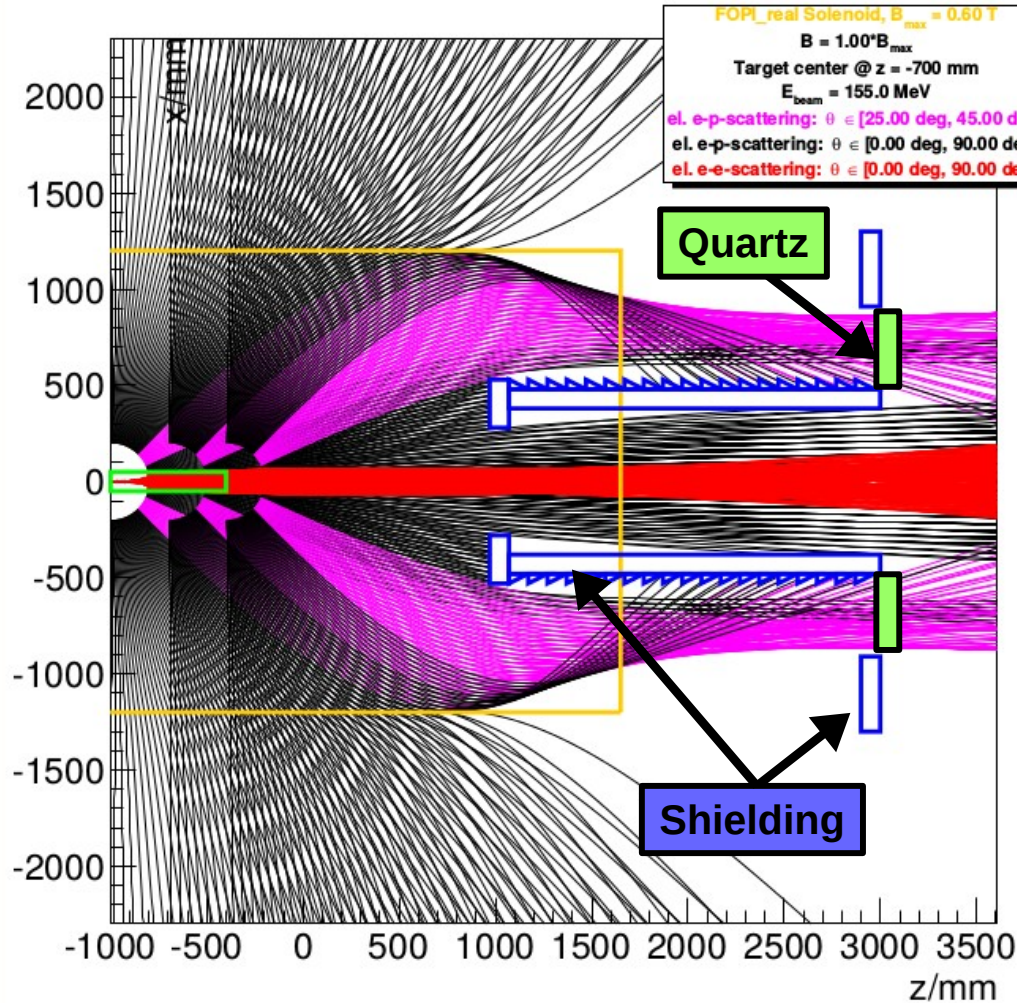
- $E_{beam} = 155 \text{ MeV}, 25\text{-}45^\circ$
- $Q^2 = 0.0049 \text{ GeV}^2$
- 60 cm  $\text{LH}_2$  target, 150  $\mu\text{A}$ , 10,000 hours
- Total rate  $\sim 0.5 \text{ THz}$
- $A = -28 \text{ ppb to } 1.5\%$
- **Improve Jlab  $Q_{weak}$ 's determination of proton weak charge by factor of 2.5**
- **0.13% precision on  $\sin^2\theta_W$**



Status:

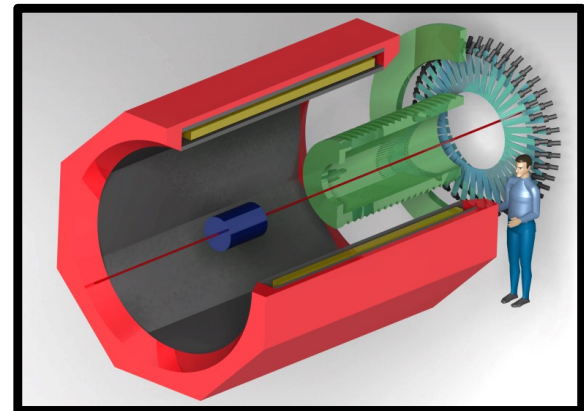
- Collaborators from Germany and US
- Funding approval by DFG
- R&D in progress
- Aim to run from 2017-2020

# Raytrace simulations in the magnetic field

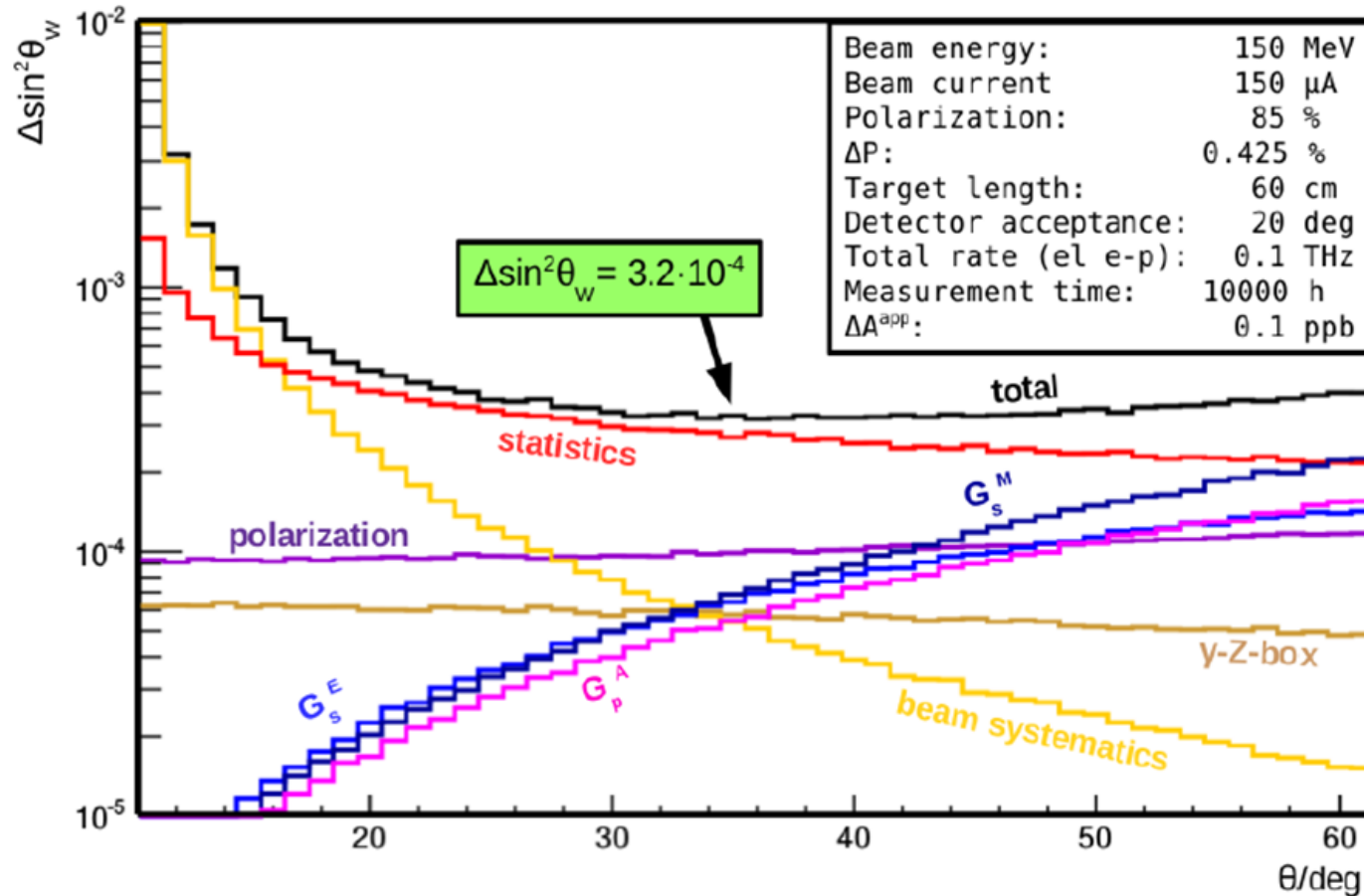


**Magnetic field:**  
**0.6 T**

Beam energy = 155 MeV  
Moller,  $\theta \in [0^\circ, 90^\circ]$   
Elastic e-p,  $\theta \in [25^\circ, 45^\circ]$   
Elastic e-p,  $\theta \in [0^\circ, 90^\circ]$



## Choice of kinematics for the P2 experiment



## Facts and Figures

The following results are based on error propagation calculations **including** the results of the Geant4 simulation of the experimental setup:

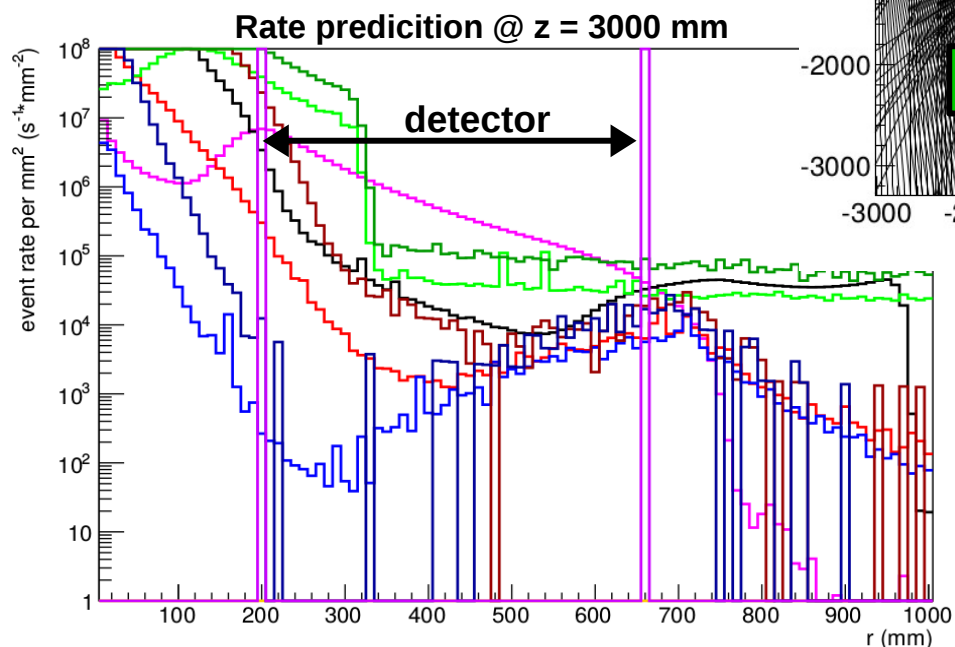
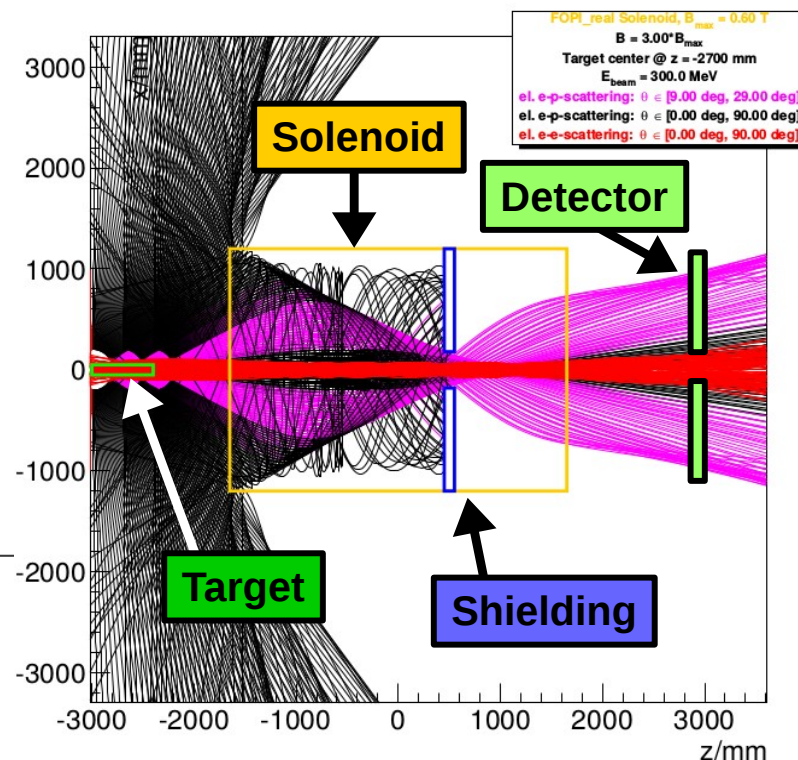
<b>Beam energy</b>	155 MeV	
<b>Beam current</b>	150 $\mu$ A	
<b>Polarization</b>	85 %	$\pm 0.425$ %
<b>Target</b>	60 cm	liquid hydrogen
<b>Detector acceptance</b>	$2\pi \cdot 20^\circ$	$\theta \in [25^\circ, 45^\circ]$
<b>Detector rate</b>	0.5 THz	
<b>Measurement time</b>	1e4 h	
<b><math>\langle Q^2 \rangle</math></b>	4.49e-3 GeV <sup>2</sup> /c <sup>2</sup>	
<b><math>A^{\text{exp}}</math></b>	-28.35 ppb	

	Total	Statistics	Polarization	Apparative	Form factors	Re( $\square_{yZA}$ )
$\Delta \sin^2(\theta_w)$	<b>3.1e-4</b> <b>(0.13 %)</b>	2.6e-4 (0.11 %)	9.7e-5 (0.04 %)	7.0e-5 (0.03 %)	1.4e-4 (0.04 %)	6e-5 (0.03 %)
$\Delta A^{\text{exp}}/\text{ppb}$	<b>0.44</b> <b>(1.5 %)</b>	0.38 (1.34 %)	0.14 (0.49 %)	0.10 (0.35 %)	0.11 (0.38 %)	0.09 (0.32 %)

# A very first idea for 300 MeV

**Beam energy: 300 MeV**  
**Beam current : 150  $\mu$ A**  
**Central magnetic field: 1.8 Tesla**

**Moller,  $\theta \in [0^\circ, 90^\circ]$**   
**Elastic e-p,  $\theta \in [9^\circ, 29^\circ]$**   
**Elastic e-p,  $\theta \in [0^\circ, 90^\circ]$**



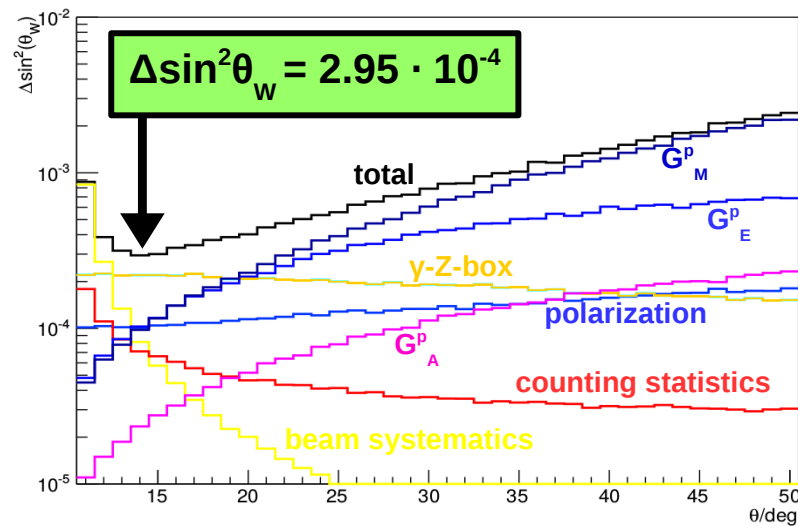
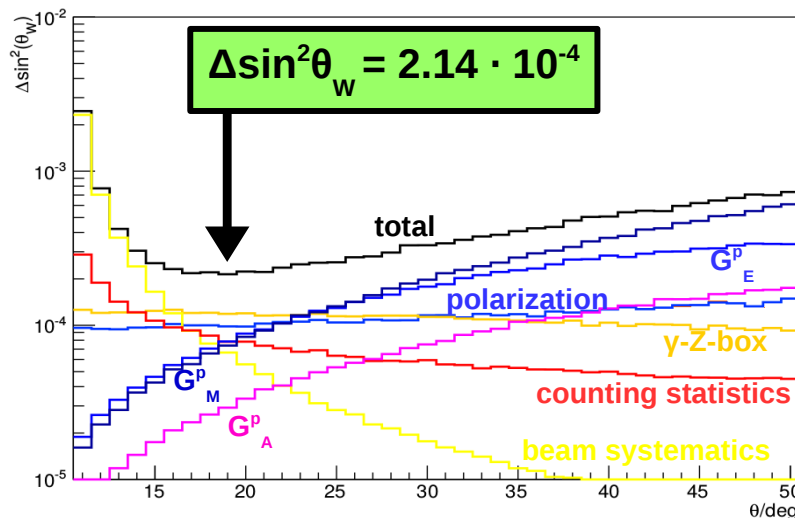
**Elastic e-p,  $\theta$  in  $[9^\circ, 29^\circ]$**   
**Elastic e-p,  $\theta$  not in  $[9^\circ, 29^\circ]$**   
**Moller, e-p**  
**Moller, background**  
**Positrons, e-p**  
**Positrons, background**  
**Photons, e-p**  
**Photons, background**

## Achievable precision @ higher energies/beam current

**Beam current:** 1 mA  
**Polarization:** 85 %  $\pm$  0.425 %  
**Target material:** liquid hydrogen  
**Target:** 60 cm  
**Measurement time:** 10000 h  
**Detector acceptance:**  $2\pi \cdot 20^\circ$   
 **$\Delta A^{\text{app}}$ :** 0.1 ppb

**Beam energy:** 300 MeV  
**Central scattering angle:**  $19^\circ$   
 $A^{\text{PV}} = (-30.8 \pm 0.34)$  ppb  
 $\langle Q^2 \rangle = 4.84e-3 \text{ GeV}^2/c^2$   
**Rate elastic e-p:** 1.8 THz

**Beam energy:** 500 MeV  
**Central scattering angle:**  $14^\circ$   
 $A^{\text{PV}} = (-24.8 \pm 0.36)$  ppb  
 $\langle Q^2 \rangle = 3.82e-3 \text{ GeV}^2/c^2$   
**Rate elastic e-p:** 3.6 THz



For 0.1mA, statistics dominated:  $\Delta(\sin^2 \theta_W) \sim 0.00032$   
 Matches P2, requires 10k hours

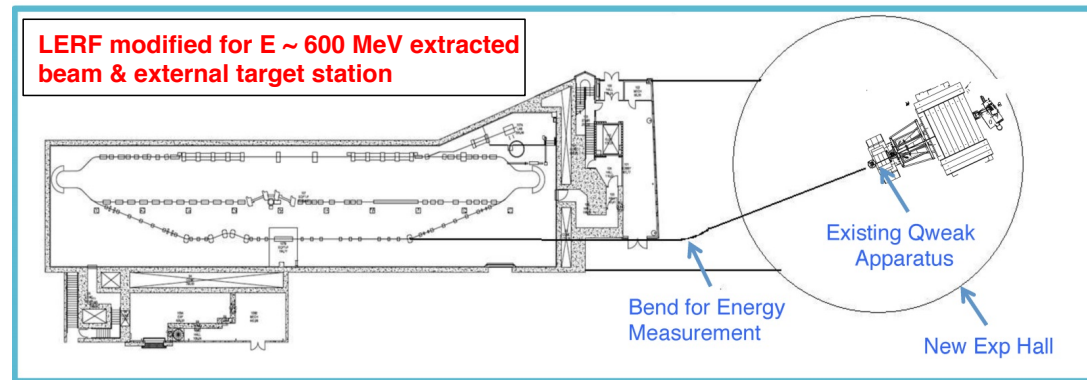
# Qweak Apparatus Reused at Lower Energy

What might be achievable by re-using the Qweak apparatus at lower beam energy for a much lower  $Q^2$  measurement of the proton's weak charge?

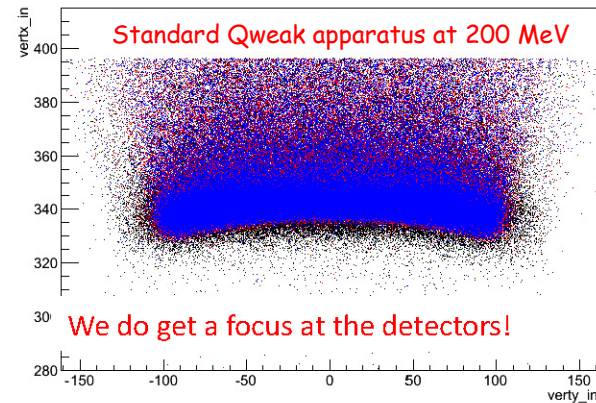
How Qweak Exp might look at an upgraded LERF

OR

other equivalent facility



Monte Carlo studies by **Juliette Mammei** and **Kurtis Bartlett** (using Qweak apparatus with same relative target/collimators/spectrometer positions, etc.) indicates there is a focus at lower energies (200 MeV to 600 MeV).



## Projections for Using Qweak Apparatus at 600 MeV

Projected rates/asymmetries for standard Qweak apparatus at 600 MeV:  
Case A: standard 2.5 kW LH<sub>2</sub> target; Case B: 3.8 kW LH<sub>2</sub> target

Parameter	MESA P2 <sup>*</sup>	Q-weak 600, case A	Q-weak 600, case B
E <sub>beam</sub>	200 MeV	600 MeV	600 MeV
Time	10000 hours	10000 hours	10000 hours
Current	150 μA	200 μA	300 μA
LH <sub>2</sub> Target Length	60 cm	35 cm	35 cm
Polarization	85%	85%	85%
Central θ	20°	8°	8°
<Q <sup>2</sup> >	.0029 GeV <sup>2</sup>	.0065 GeV <sup>2</sup>	.0065 GeV <sup>2</sup>
Total rate	440 GHz	30 GHz	44 GHz
Asym. Width @240 Hz	23 ppm	89 ppm	74 ppm
A <sub>phys</sub> (ppb)	-20 ppb	-46 ppb	-46 ppb
Hadronic “B” term	9%	10%	10%
ΔA (stat)	0.25 ppb (1.2%)	0.96 ppb (2.1%)	0.79 ppb (1.7%)
ΔA (syst)	0.19 ppb (0.9%)	0.41 ppb (0.9%)	0.41 ppb (0.9%)
ΔA (tot)	0.34 ppb (1.7%)	1.20 ppb (2.6%)	1.01 ppb (2.2%)
ΔQ <sub>W</sub> <sup>P</sup>	0.0014 (2.0%)	0.0021(3.0%)	0.0019 (2.6%)
Δsin <sup>2</sup> θ <sub>W</sub>	3.6x10 <sup>-4</sup> (0.15%)	5.4x10 <sup>-4</sup> (0.23%)	4.7x10 <sup>-4</sup> (0.20%)

\* MESA P2 parameters come from F. Maas talk at “Dark Forces at Accelerators” Frascati, Oct. 2012

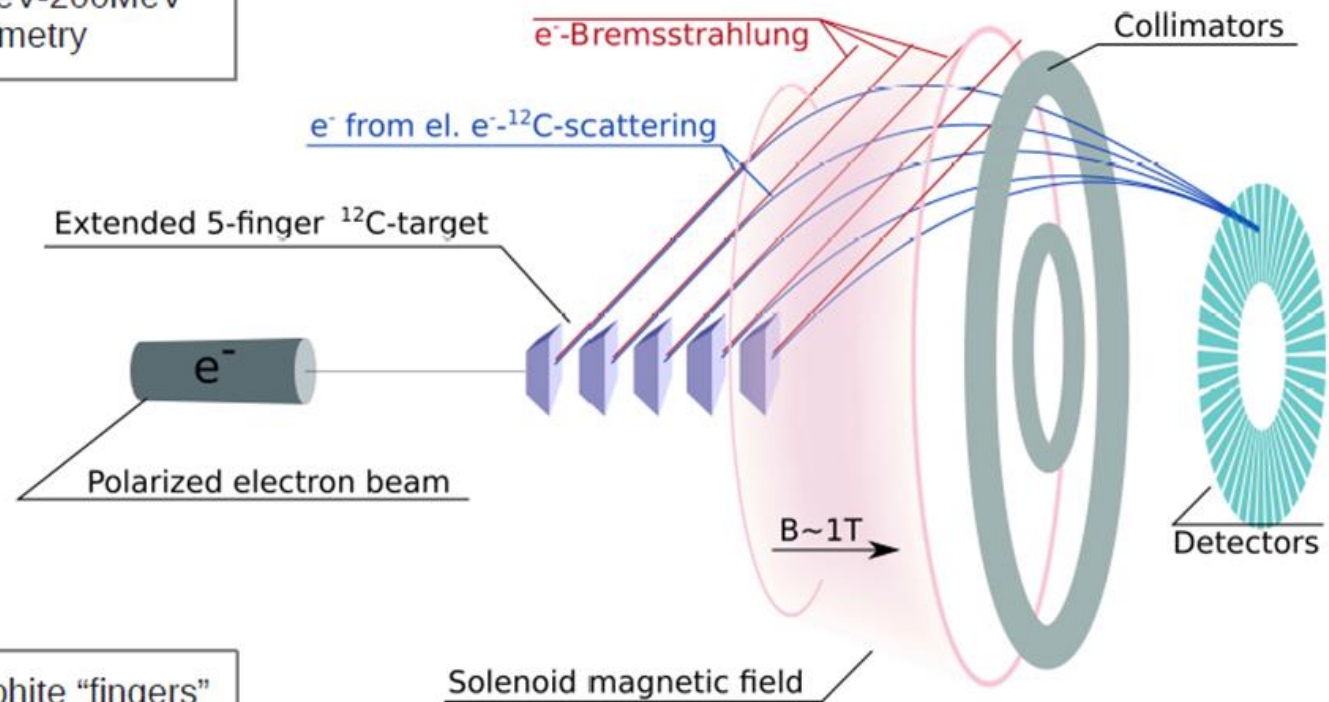
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# Measurements with other targets at P2

- MESA:
- 150 $\mu$ A
  - 150MeV-200MeV
  - Polarimetry



- Target:
- 5 graphite "fingers"
  - 5 g/cm $^2$  total
  - 36mm spacing

# C12 @ P2 MESA – Weak Charge of the $^{12}\text{C}$ Nucleus

PV Elastic Electron Scattering  
on J=0, T= 0 Nucleus:

$$A_{PV} = \frac{G_F Q^2}{\sqrt{2\pi\alpha}} \sin^2 \theta_W$$

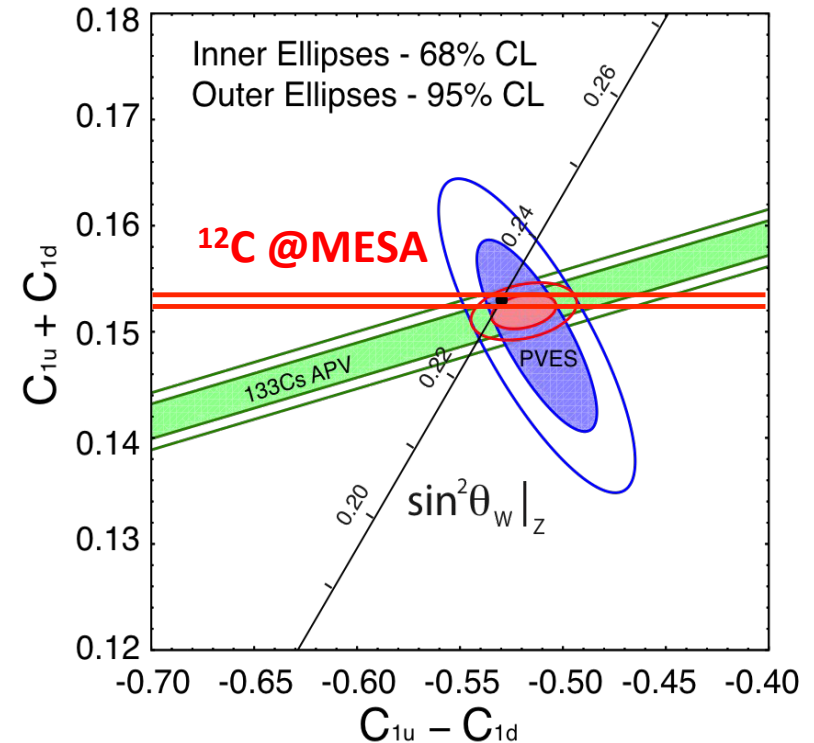
## $^{12}\text{C}$ measurement at P2

@ Beam energy  $E = 150\text{MeV}$       Measuring time  $t = 2500\text{h}$   
 Scattering angle  $\theta = 40^\circ \pm 9^\circ$       Beam current  $I = 150\mu\text{A}$   
 Target density  $d = 5\text{g/cm}^2$

We can achieve  $\frac{\delta \sin^2 \theta_W}{\sin^2 \theta_W} = 0.3\%$

$$A_{PV} = \frac{G_F \cdot Q^2}{\sqrt{2\pi\alpha}} \sin^2 \theta_W \longrightarrow \frac{\delta A_{PV}}{A_{PV}} = \frac{\delta Q_W^C}{Q_W^C} = \frac{\delta \sin^2 \theta_W}{\sin^2 \theta_W} = 0.3\%$$

$$Q_W^C = -24 \sin^2 \theta_W \longrightarrow$$



**S. Baunack**

3x better than APV, nearly same coupling combination  
 Moderate runtime

# Inelastic Levels: Experiment

$Q = 240 \text{ MeV}$   
 $= 1.2 \text{ fm}^{-1}$

566

1961 R. HOFSTADTER

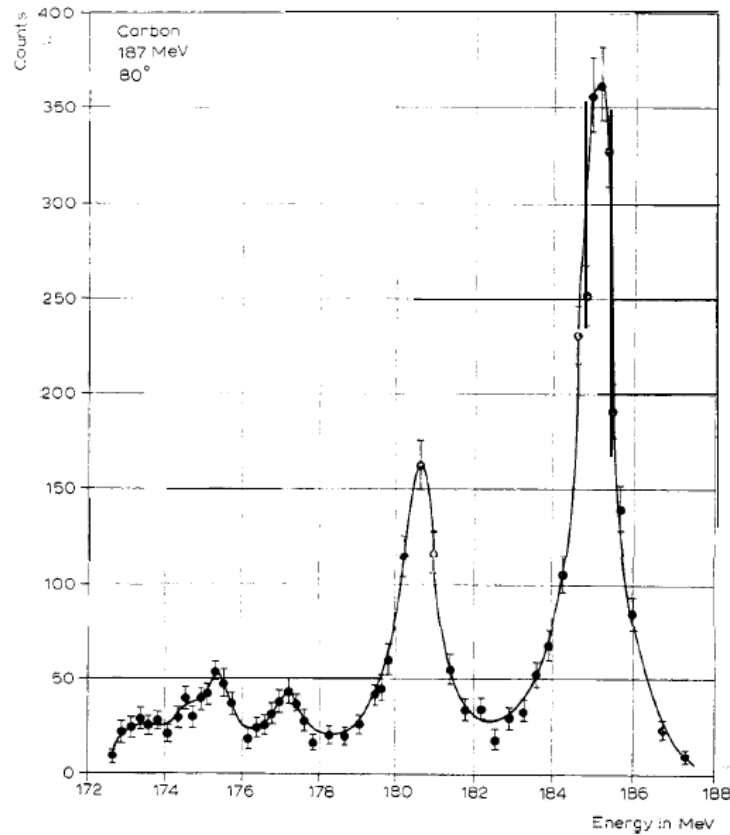
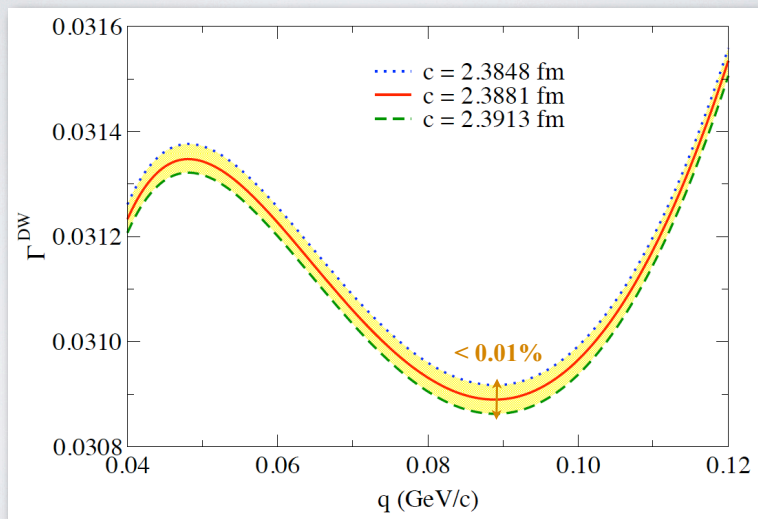


Fig. 4. This figure shows the elastic-scattering peak from carbon at an abscissa near 185 MeV, and the inelastic-scattering peaks from the excited states of  $^{12}\text{C}$ . The peak near 180.7 MeV is associated with the 4.43-MeV level.

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## COULOMB DISTORTION EFFECT

Theoretical uncertainties in the PV asymmetry due to a reasonable variability in the nuclear charge distributions

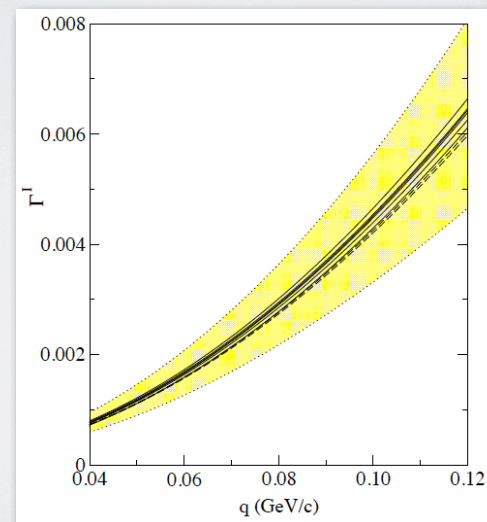


PRC 89 (2014) 015501

## NUCLEAR ISOSPIN MIXING EFFECT

Theoretical uncertainties in the PV asymmetry from a reasonable variability in the neutron vs. the proton distributions

DIFFERENT NUCLEON-NUCLEON FORCES: HF MEAN FIELD



PRC 89 (2014) 015501

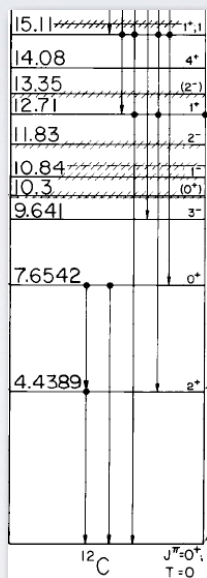
## INELASTIC TRANSITIONS EFFECT

Theoretical uncertainties in the PV asymmetry due to possible transitions to excited states in the target

Total asymmetry:

$$A = \sum_{i=0}^n f_i A_i$$

$$f_i = \frac{\sigma_i^{EM}}{\sum_{k=0}^n \sigma_k^{EM}}$$



T = 1 levels:  
Different asymmetry  
(isovector contribution)

T = 0 levels:  
Same asymmetry  
as for ground state  
( $\Gamma^{incl.} = 0$ )  
(in absence of strangeness  
and at tree-level)

PRC 89 (2014) 015501

## CONCLUSIONS

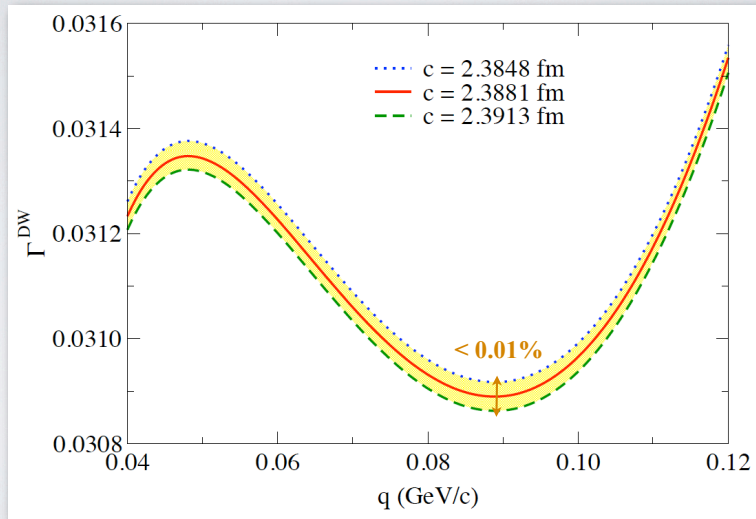
### SUMMARY OF SIZES AND UNCERTAINTIES

150 MeV incident energy, 25°-45° scattering angular range

Contribution to PV asymmetry	Relative size	Relative uncertainty
Coulomb distortion of projectile wave function	3%	0.01%
Nuclear isospin mixing (electromagnetic origin)	0.4%	0.05%
Nucleon strangeness content (mainly electric)	0 - 1 %	1%
Meson exchange currents	< 0.1 %	< 0.1 %
Inelastic contributions	< 0.1 %	--

## COULOMB DISTORTION EFFECT

Theoretical uncertainties in the PV asymmetry due to a reasonable variability in the nuclear charge distributions

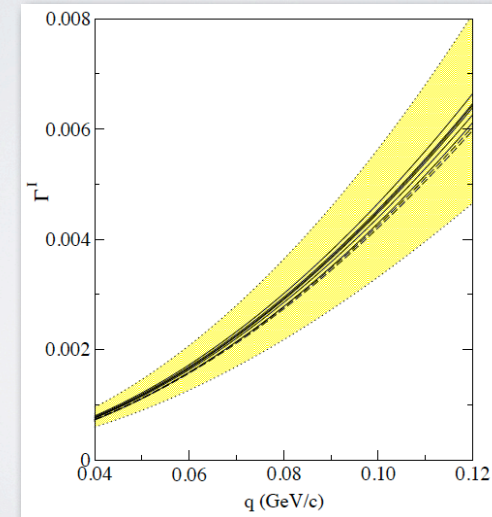


PRC 89 (2014) 015501

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Theoretical uncertainties in the PV asymmetry from a reasonable variability in the neutron vs. the proton distributions

DIFFERENT NUCLEON-NUCLEON FORCES: HF MEAN FIELD



PRC 89 (2014) 015501

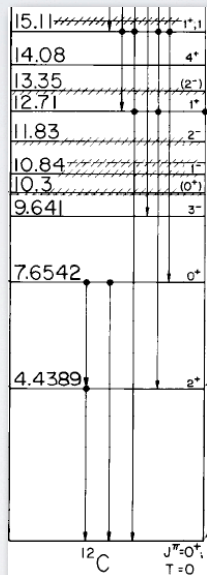
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T = 1 levels:  
Different asymmetry (isovector contribution)

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Same asymmetry as for ground state ( $\Gamma^{incl.} = 0$ )  
(in absence of strangeness and at tree-level)

PRC 89 (2014) 015501

## CONCLUSIONS

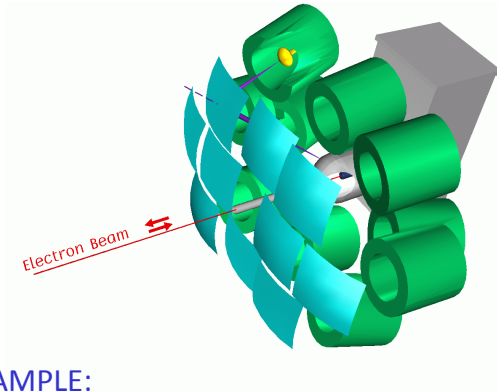
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Meson exchange currents	< 0.1 %	< 0.1 %
Inelastic contributions	< 0.1 %	--

# Strange Form Factors – Worldwide Program

1992 – 2011: Worldwide program on strange form factors measured with PVES



**SAMPLE:**

Location: MIT-Bates

Targets: p,d

Kinematics: backward angle,  $Q^2 = .038, .10 \text{ GeV}^2$



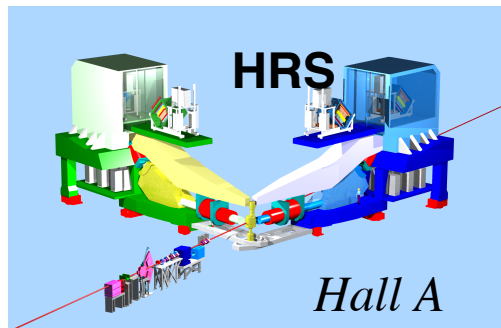
**Mainz PV-A4:**

Location: Mainz MAMI microtron

Targets: p,d

Kinematics: forward & backward angles

$Q^2 = .11, .23, .62 \text{ GeV}^2$

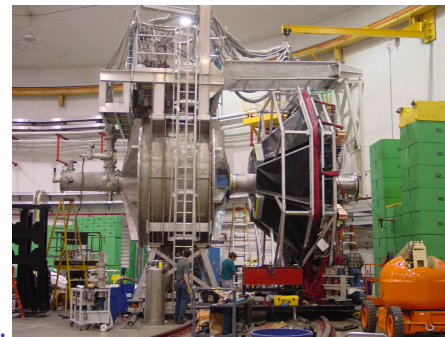


**HAPPEX I, II, III:**

Location: Jefferson Lab Hall A

Targets: p,  $^4\text{He}$

Kinematics: forward angle,  $Q^2 = .10, .48, .62 \text{ GeV}^2$



**G<sup>0</sup>:**

Location: Jefferson Lab Hall C

Targets: p, d

Kinematics: forward & backward angles

$Q^2 = .1 - 1 \text{ GeV}^2$

6/17/2015

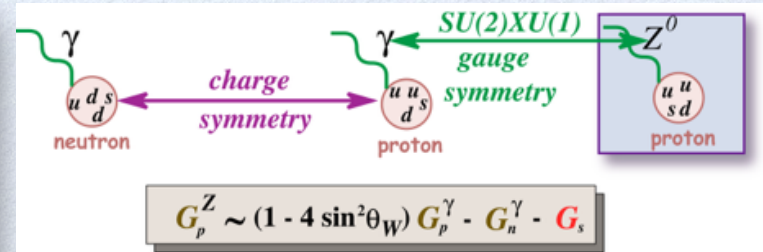
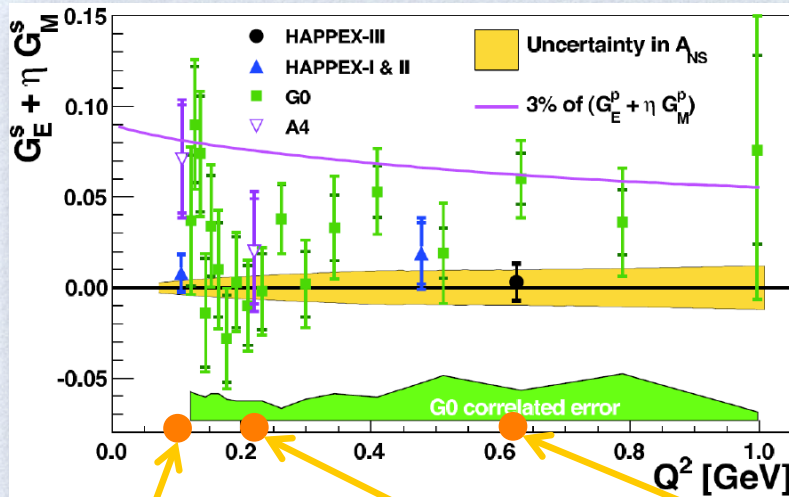
Intense Electron Beams Workshop

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2011: Completion of a 2-decade program

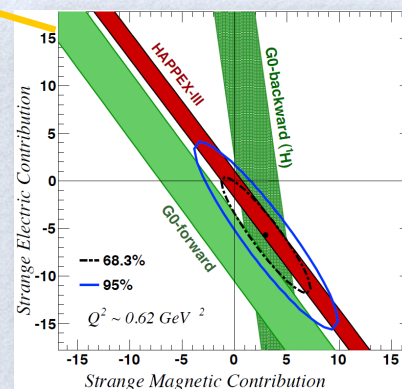
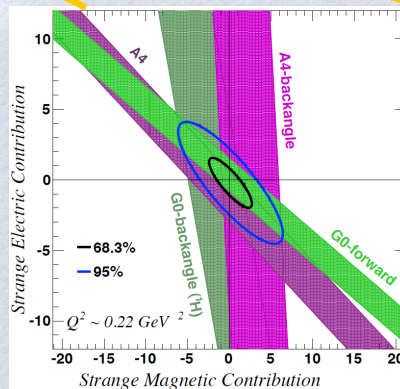
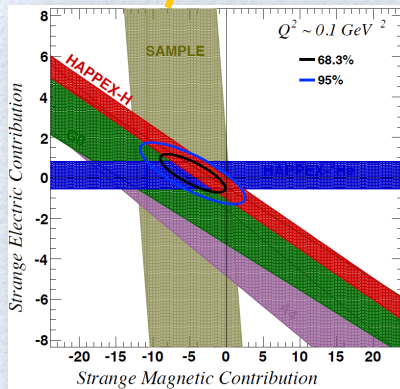
# Strange Form Factor Summary

Strange quarks carry nucleon momentum: Other external properties affected?



$\Rightarrow G_E^s(Q^2), G_M^s(Q^2)$

- Sensitive Flavor separation at 3  $Q^2$  values
- No more than few % of EM structure
- Recent lattice results in agreement





# Strange Form Factors – Measurements at Low Energy?

Are further strange form factor measurements warranted?

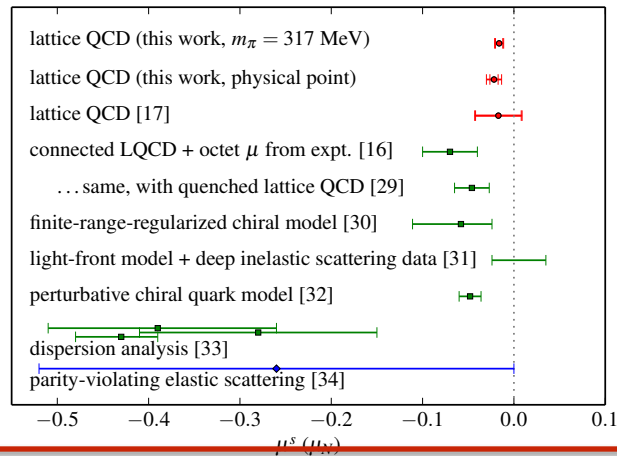
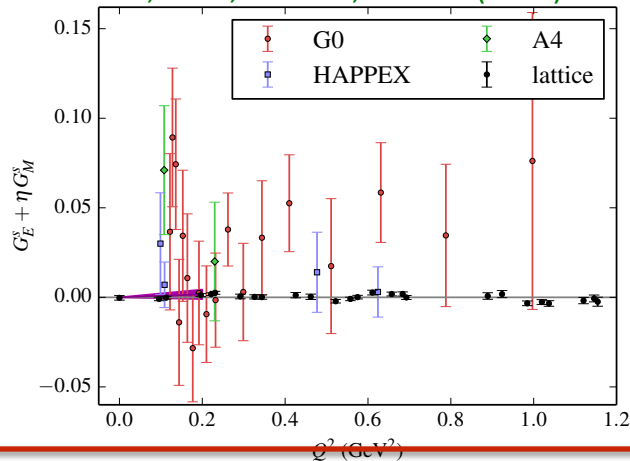
State-of-the-art lattice QCD calculations set the scale of what is interesting.

Recent lattice predictions for the strange magnetic moment:

$$G_M^s(Q^2 = 0) \equiv \mu_s = -0.07 \pm 0.03 \mu_N \quad \text{Green, et al., arXiv:1505.01803}$$

$$G_M^s(Q^2 = 0) \equiv \mu_s = -0.022 \pm 0.004 \pm 0.004 \pm 0.006 \mu_N \quad \text{Shanahan, et al., PRL 114, 091802 (2015)}$$

Shanahan, et al., PRL 114, 091802 (2015)



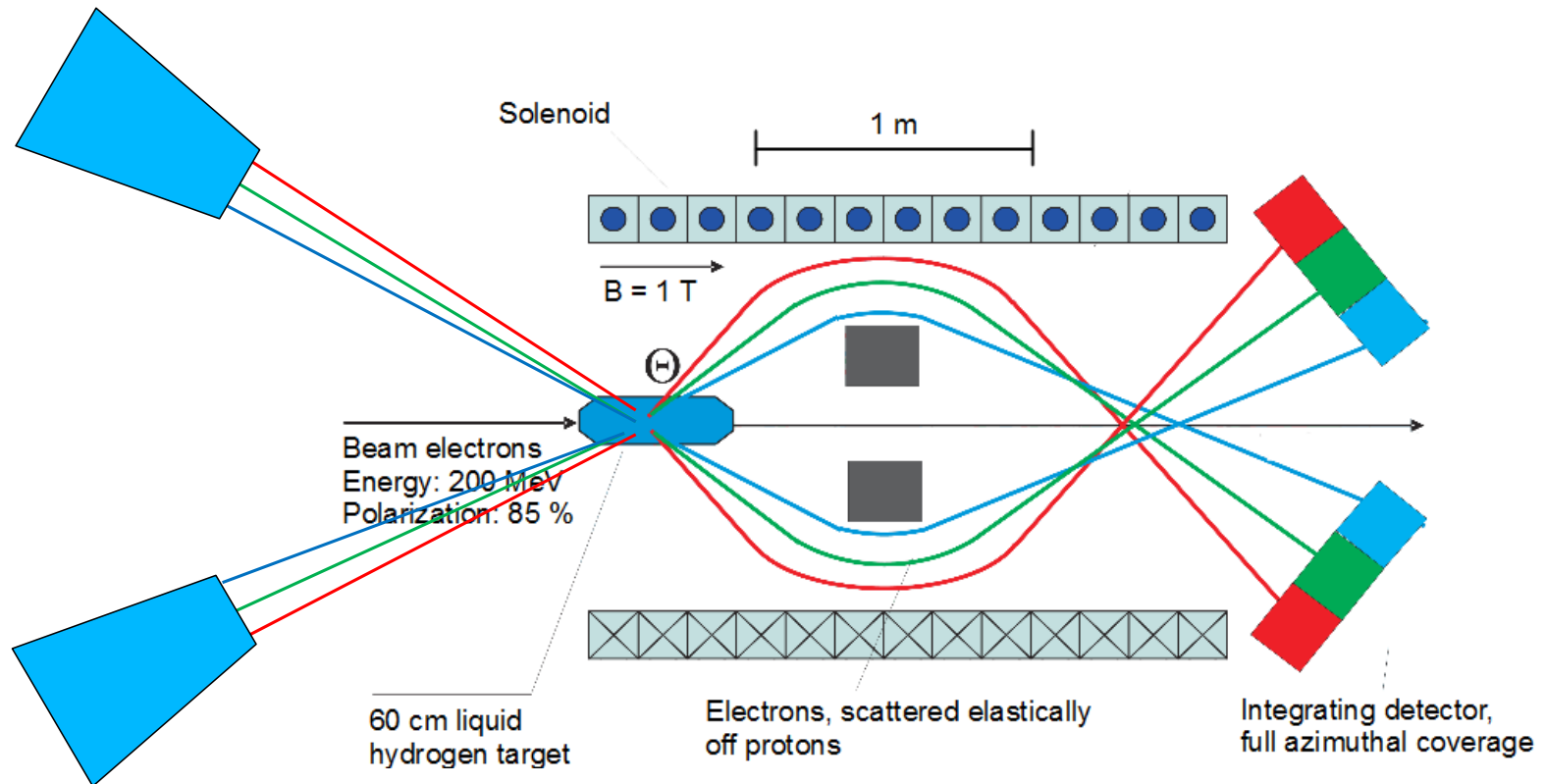
## Possible backangle measurements at low energies?

- "A4 style" fast calorimeter during P2,  $\theta \sim 140 - 150^\circ$ , 150 MeV, 150  $\mu\text{A}$ , 60 cm LH/D<sub>2</sub> targets, 1000 hours each  $\rightarrow \delta G_M^s \sim \pm 0.05 \mu_N$  (Baunack, PEB2013)
- "SAMPLE" style air Cerenkov,  $\theta \sim 130 - 170^\circ$ , not yet estimated

See K. Kumar talk for possibilities for strange radius at low energies

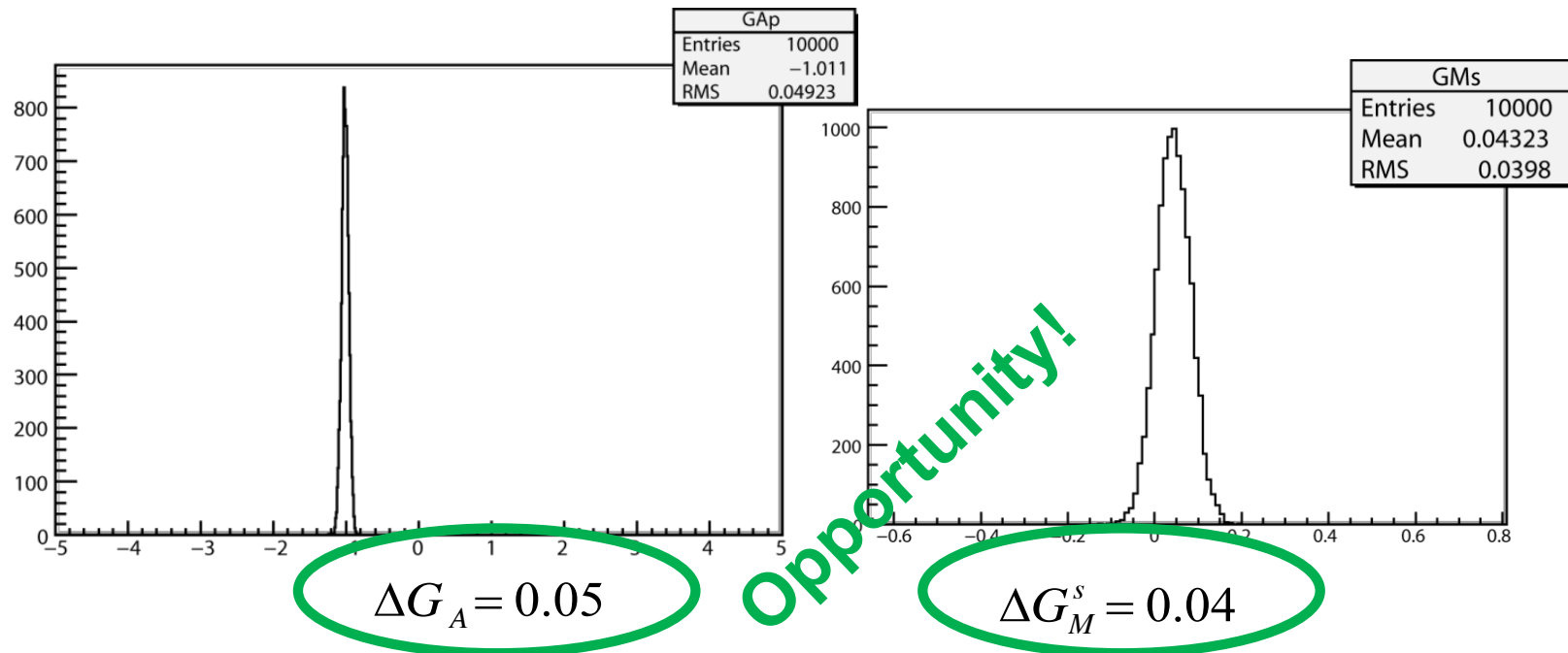
# P2 back angle measurement!

Back angle measurements: Determination of  $G_M^S$  and  $G_A$



# Possible uncertainties of $G_A$ and $G_M^s$ with P2 back angle measurement

- $Q^2=0.06 \text{ GeV}^2$
- Numerical determination of precision
- Choose randomly EM form factors and asymmetries according to their uncertainties and calculate  $G_A$  and  $G_M^s$
- Correlation of electromagnetic form factors input taken into account

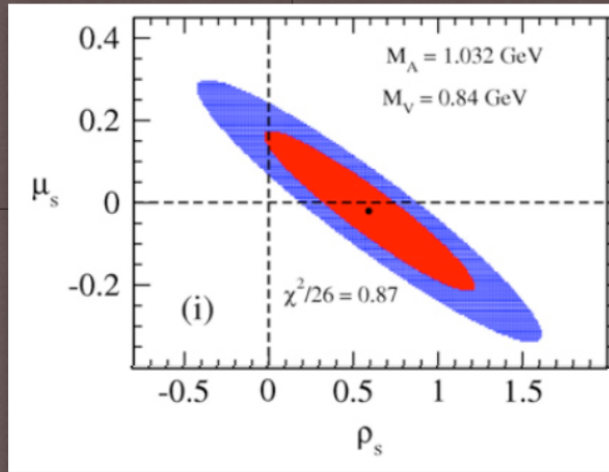


# Strangeness using isoscaler nucleus: $^{12}\text{C}$

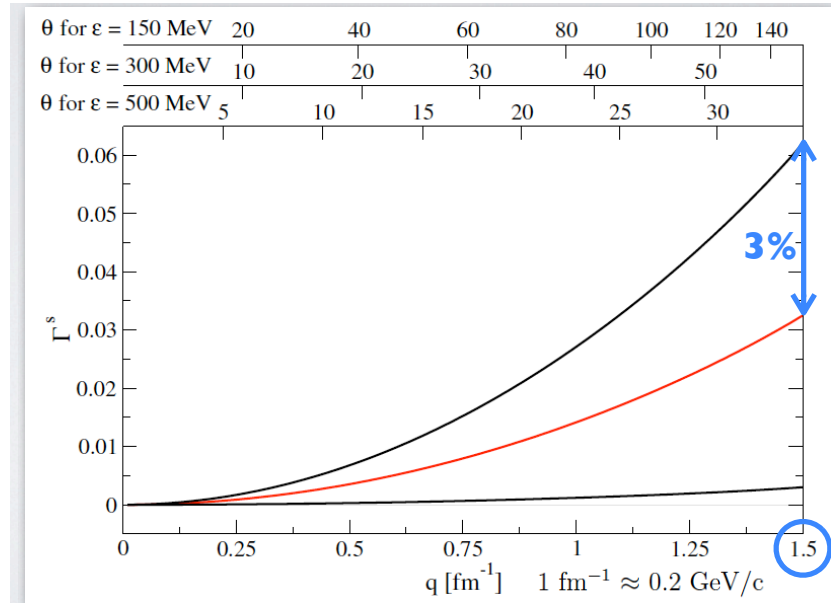
$$G_E^{(s)} = \frac{|Q^2|}{4m_N^2} \frac{\rho_s}{(1 + |Q^2|/M_V^2)^P}$$

$$G_M^{(s)} = \frac{\mu_s}{(1 + |Q^2|/M_V^2)^P}$$

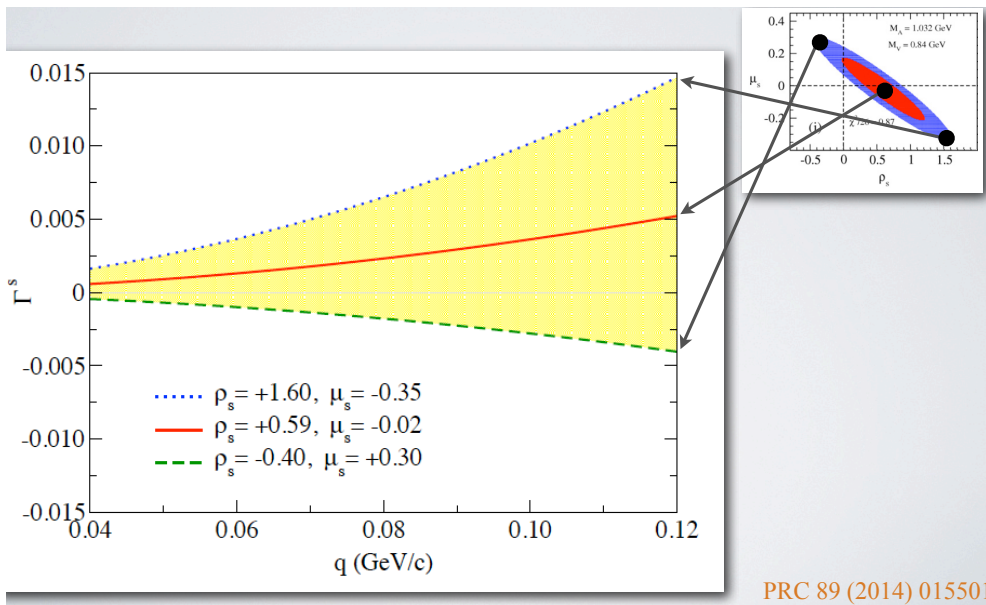
$$\chi^2 = \sum_{j=1}^{28} \left( \frac{\mathcal{A}_j^{\text{exp}} - \mathcal{A}_j^{\text{th}}}{\Delta \mathcal{A}_j^{\text{exp}}} \right)^2$$



R. González-Jiménez, J. A. Caballero, T. W. Donnelly, Phys. Rep. 524 (2013) 1



Strangeness uncertainty for  $^{12}\text{C}$  measurement



PRC 89 (2014) 015501

Measurement of  $^{12}\text{C}$  at higher  $q$  pins strangeness radius ( $G_E^s$ ), calibrates low- $q$  Standard Model study

**Moderate running time, needs 300 MeV or more**

O. Moreno

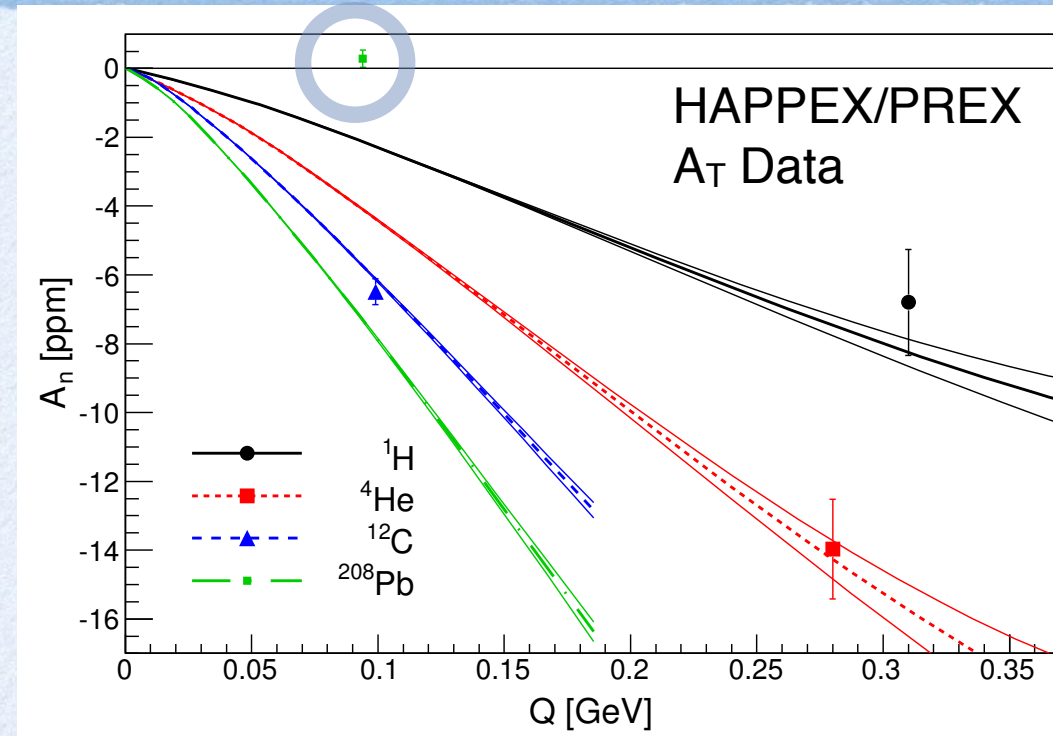
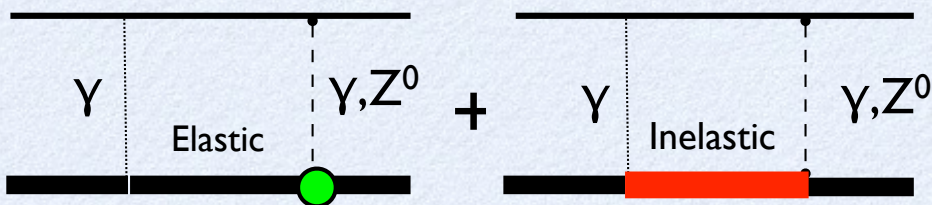
## Summary: Weak Charge

- Proton weak charge - hard to beat P2
- $^{12}\text{C}$  can provide powerful SM test (2500 hrs)
- $^{12}\text{C}$  requires additional precision on GEs.
  - Hard to do at MESA, needs 300 MeV

# Vector Analyzing Power

$$A_T \equiv \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \propto \vec{S}_e \cdot (\vec{k}_e \times \vec{k}'_e)$$

We measured this, in part, because it is a possible systematic error for the PV measurements.



- **What does the Pb-208  $A_T$  result imply?**
  - *dispersion corrections on top of Coulomb distortions?*
  - *What if it is a very sensitive cancellation?*
    - *What happens when we run again at slightly different kinematics?*
    - *What if Ca-48 doesn't have this accidental cancellation?*
  - *should other electroweak corrections be revisited?*
  - *Motivates more  $A_T$  measurements at different energies*

# Experimental Requirements

$$\mathcal{L} \sim 10^{39} / (\text{s-cm}^2)$$

Charge-normalized yield :  $Y = \frac{S}{I}$       S: Integrated detector signal  
 I: Integrated charge measurement

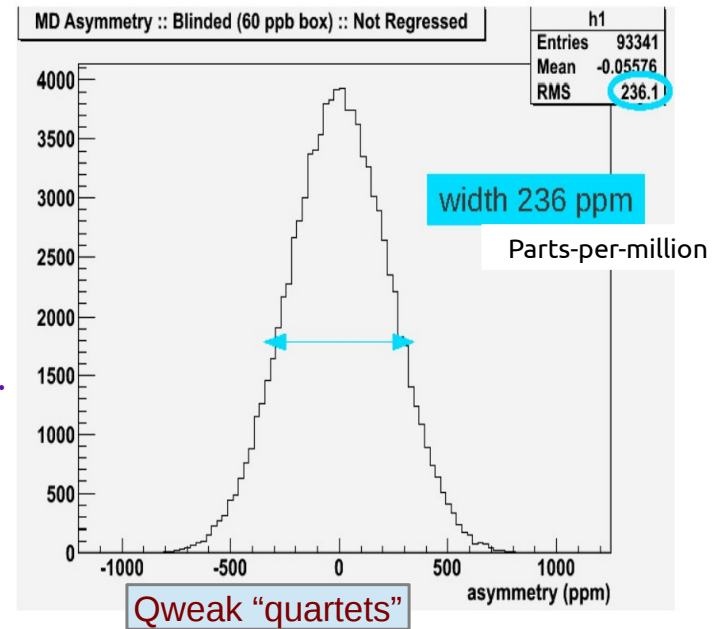
Requires precise (relative) charge measurement

Raw measured asymmetry :

$$A_{raw} = \frac{Y_+ - Y_-}{Y_+ + Y_-}$$

High precision (part-per-billion level) achieved through repeated measurements.

RMS of distribution important figure-of-merit.  
 Noise contributions must be suppressed, precision monitoring required

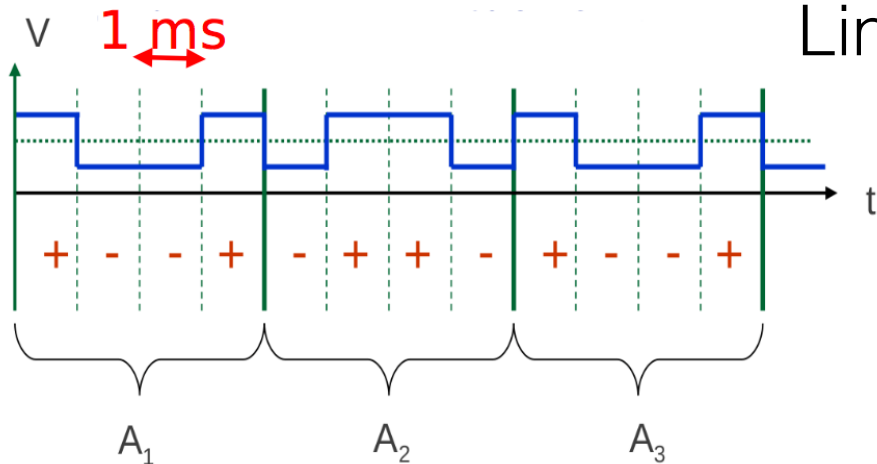


## M. Kargiantoulakis

Jun 17, 2015

Intense Electron Beams Workshop, Cornell University

4



Linear integration over helicity “window”

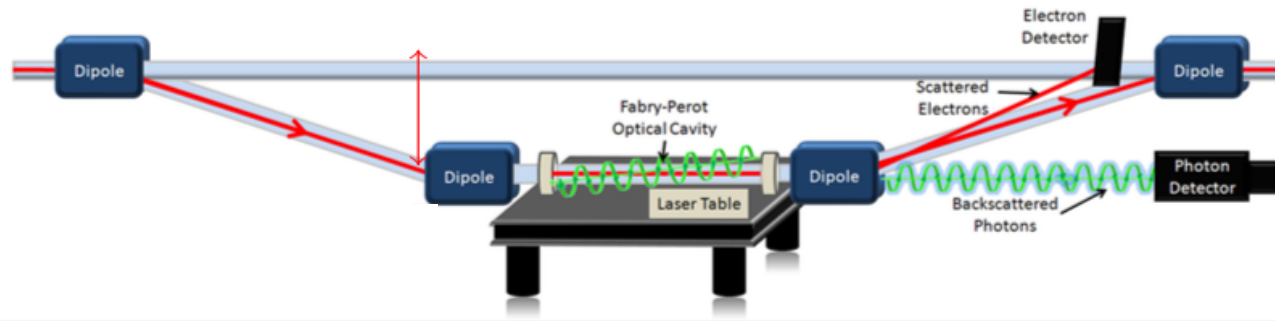
RF beam monitors - well known technology, but linear integration requirement is sometimes different  
 modest resolution  $\sim 1$  micron over 1 ms



# Polarimetry

## Compton

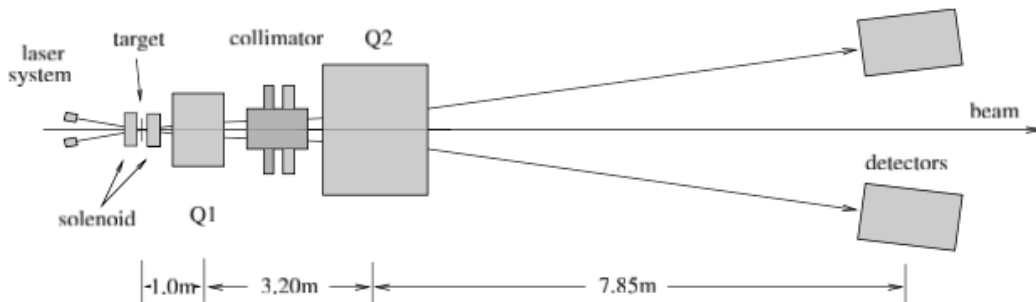
532 nm laser, 300 MeV:  
Compton edge  $\sim 3$  MeV



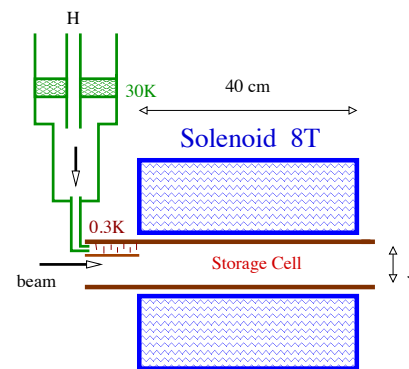
Require 2+ meter dispersion for  
electron measurement

Possible for 0.5% or better...  
but very hard

## Atomic hydrogen Moller



### Storage Cell



First: 1980 (I.Silvera, J.Walraven)  
 $\vec{p}$  jet (Michigan)  
Never put in high power beam

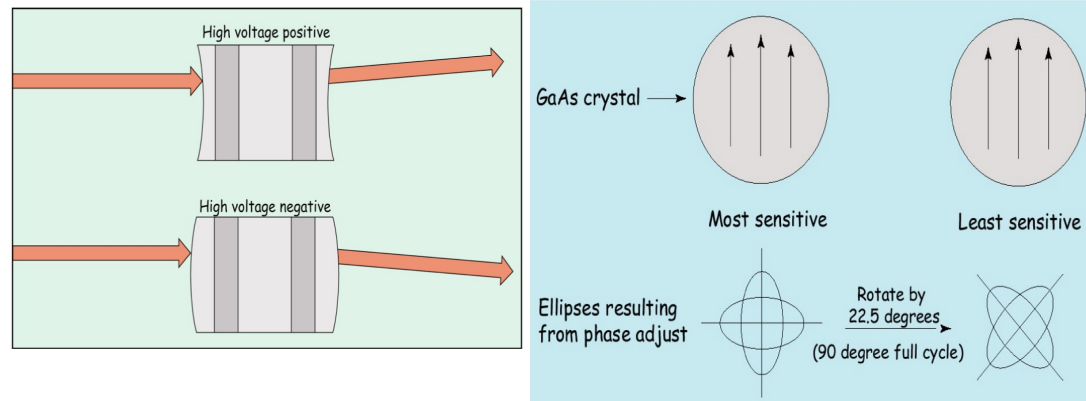
- $-\vec{\nabla}(\vec{\mu}_H \vec{B})$  force in the field gradient
  - pulls  $|a\rangle, |b\rangle$  into the strong field
  - repels  $|c\rangle, |d\rangle$  out of the field
- $H+H \rightarrow H_2$  recombination (+4.5 eV)  
high rate at low T
  - parallel electron spins: suppressed
  - gas: 2-body kinematic suppression
  - gas: 3-body density suppression
  - surface: strong unless coated  
 $\sim 50$  nm of superfluid  $^4\text{He}$
- Density  $3 \cdot 10^{15} - 3 \cdot 10^{17} \text{ cm}^{-3}$ .
- Gas lifetime  $> 1$  h.

from E. Chudakov

## Generation of Helicity-Correlated differences in the source

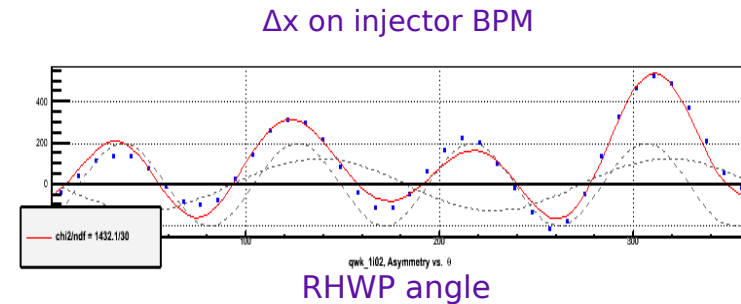
**Mechanical** PC steering

**Polarization** effects:  
PC birefringence gradients  
coupled with cathode  
analyzing power



Optimization strategies:

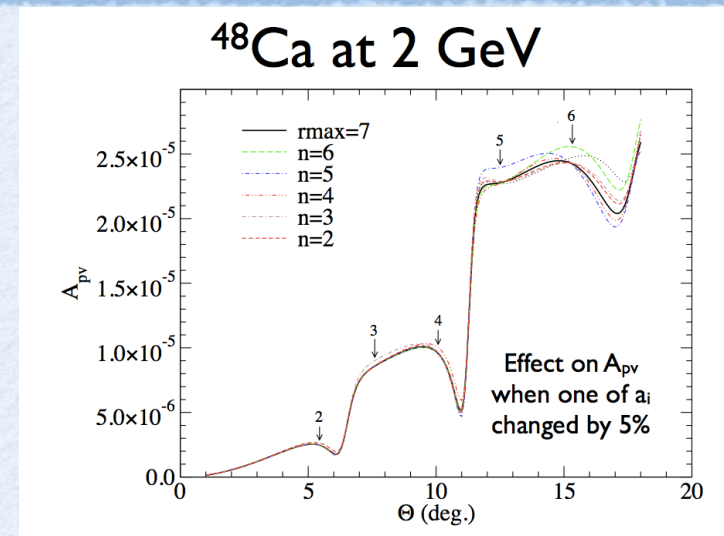
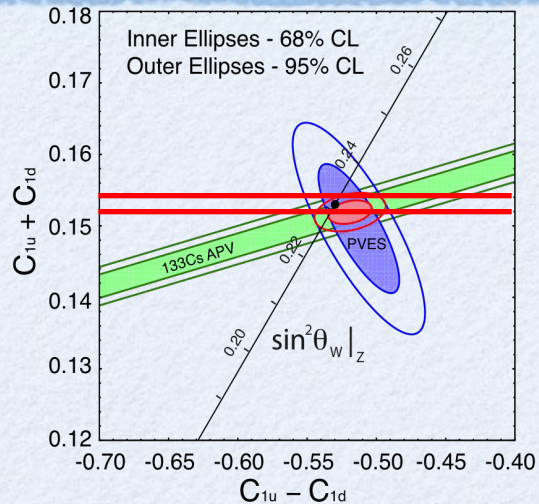
- Careful alignment on laser table
- Balance residual linear polarization from PC with vacuum window birefringence and cathode analyzing power



# Apparatus

- Assumption: 100 microAmp at 300 MeV
- Solenoid, not toroid (resolution to isolate elastic signal)
- Extracted beamline
  - Space for apparatus, diagnostic beamline, fast raster
  - beam height  $\sim 3\text{m}$
  - Space for polarimeters - atomic hydroMoller
  - Beam dump (with acceptance for disrupted beam)
- High dispersion point (few meters?) for E measurement
  - (and another?) for Compton Polarimetry (2+ meters)
- Linear integrating beam monitors, spanning phase space
- Special considerations in polarized source

# Potential of New Machine



- **Program with 300 MeV, 100 microamps**
  - *Higher current, higher E would help*
- **New measurements on Carbon-12**
  - *A Standard Model test extremely interesting if 0.3% can be reached*
  - *Must be coupled with higher  $Q^2$  measurements to constrain strange quark radius (strange quark contribution to charge radius)*
- **New measurements on neutron-rich nuclei**
  - *“Super” PREX/ CREX / Sn-REX*
  - *Higher  $Q^2$  measurements will provide a complete and model-independent distribution of neutrons in the ground state*







