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Physics & Astronomy



Sub-fs electron beam generation at UCLA Pegasus Laboratory

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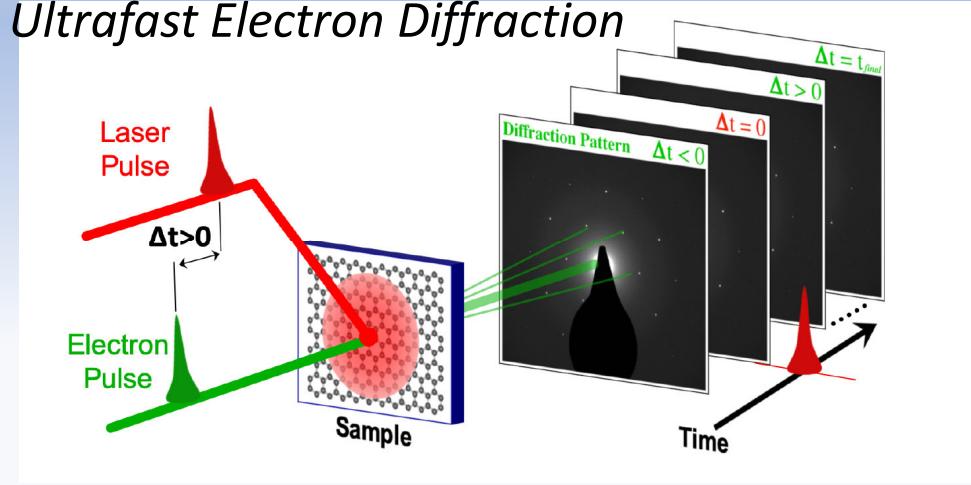
- Funding sources

- DOE Accelerator Stewardship DE-SC0009914.
- NSF Accel Science PHY-1734215
- GBMF4744 Accelerator-on-a-Chip

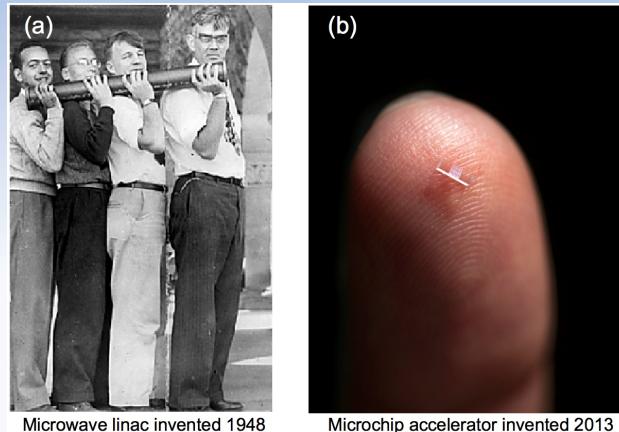


Application of sub-fs electron bunches

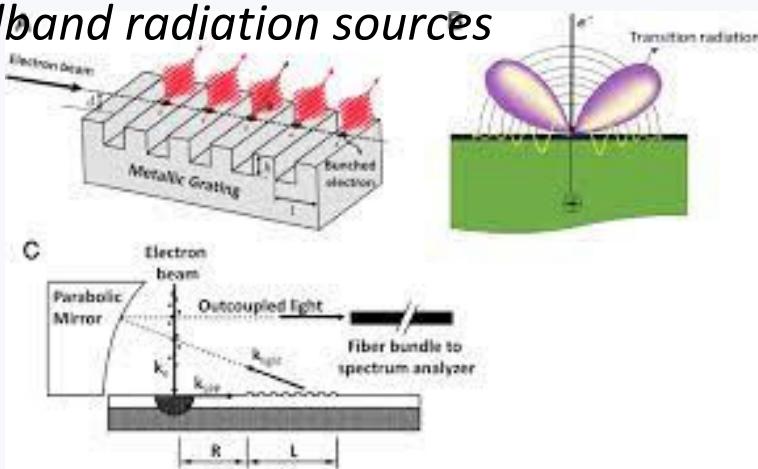
Ultrafast Electron Diffraction



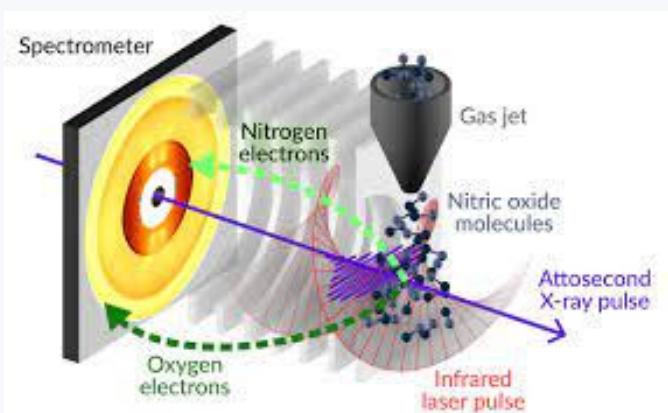
Injection into optical-scale accelerators



Broadband radiation sources



Electron-based attoscience



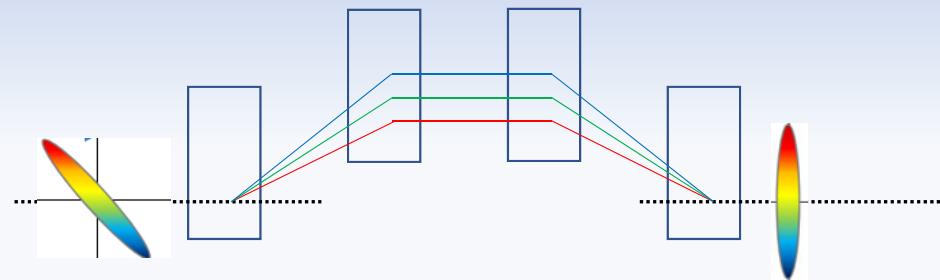
Outline

- Velocity bunching compression for low energy beamlines
- Non linear effects in ultrashort bunch generation
 - RF and space charge. (P. Denham et al. <https://arxiv.org/abs/2106.02102>)
- Compensation of non-linearities
 - Laser shaping.
 - Higher harmonic RF compensation.
- Start-to-end simulations of PEGASUS beamline.
- Sub-fs bunch length measurement techniques.

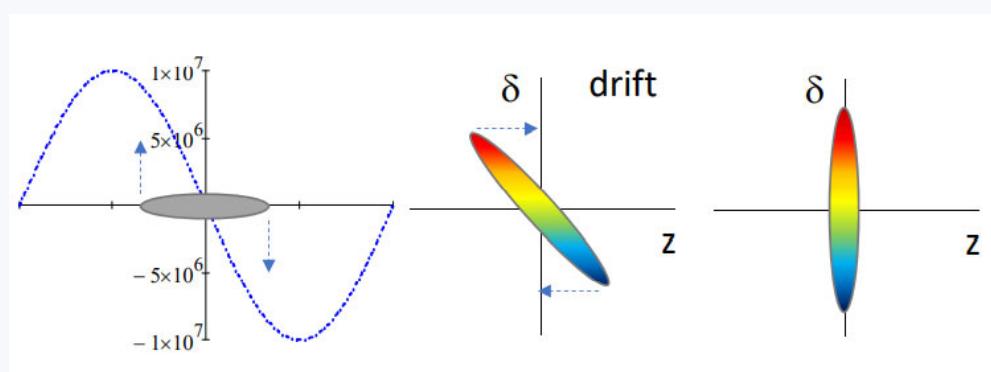
How to obtain ultrashort electron bunches

- Need a strongly chirped longitudinal phase space.
- **Magnetic compression:** higher energy electrons take a longer path through the bends, while lower energy take a shorter path. Longitudinal focus is achieved in a shorter distance. Needed for ultra-relativistic beams.
- **Ballistic compression:** RF cavity imparts negative chirp on beam then beam drifts to longitudinal focus. Favored at lower energy.

- **Magnetic compression**

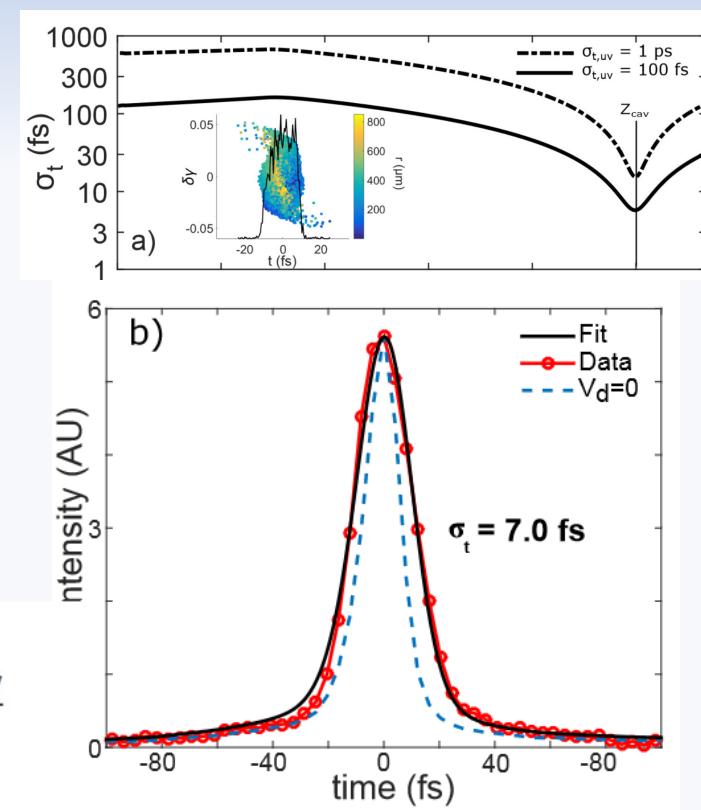
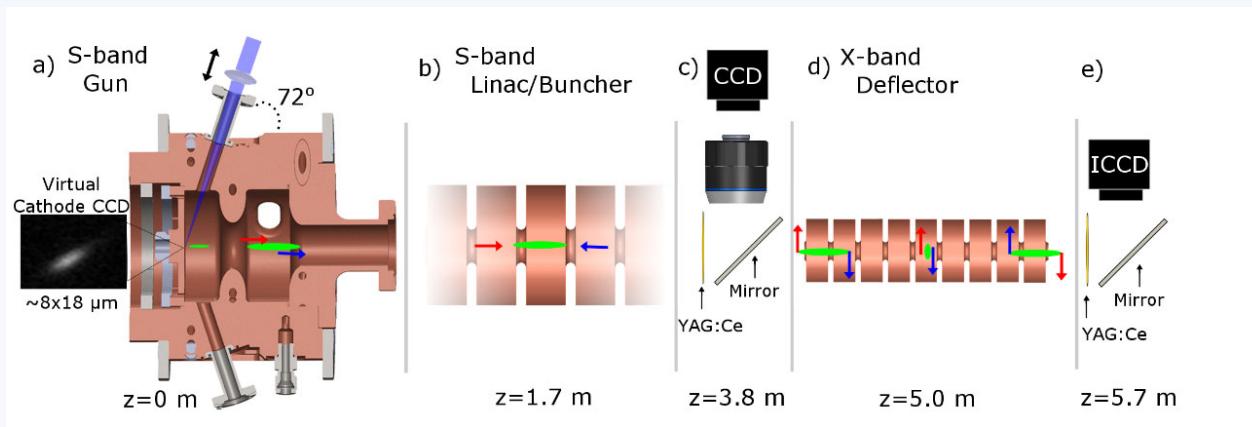


- **Ballistic compression**



Sub 10fs RF compression at few MeV energies

- Demonstrated sub 10fs RF compression at Pegasus.
- Shorter initial pulse achieved shortest focus.
- **What are the limitations?**
 - How do optimal bunch lengths scale with beam and RF parameters?



Non linear effects in RF compression

Ballistic bunching is inherently non-linear.

$$f = \frac{1}{\eta_1 \alpha k} = \frac{m_0 c^2 \gamma^3 \beta^2}{e V_0 k}$$

RF curvature

$$\Delta\gamma = -\alpha \sin(kz_0)$$

Vacuum dispersion

$$z = z_0 + s \frac{\Delta\beta}{\beta} = z_0 + s \sum_{n=1}^{\infty} \eta_n \Delta\gamma^n$$

$$\left. \begin{array}{l} \eta_1 = \frac{1}{\beta^2 \gamma^3} \\ \eta_2 = \frac{2 - 3\gamma^2}{2\gamma^6 \beta^4} \\ \eta_3 = \frac{2 - 5\gamma^2 + 4\gamma^4}{2\gamma^9 \beta^6} \end{array} \right\}$$

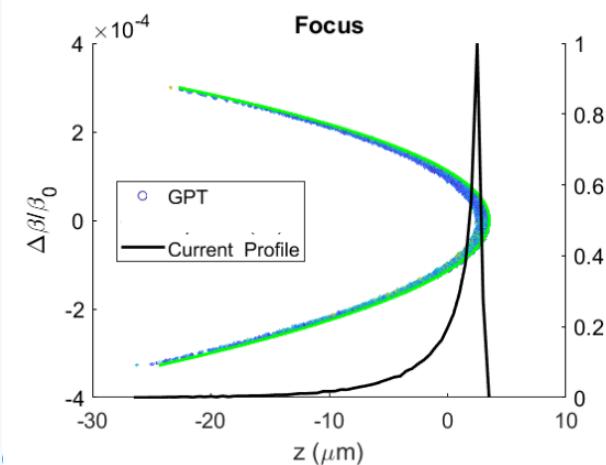
- These cause emittance growth

$$\epsilon_{RF}^2 \approx \sigma_{z0}^2 \left[2\eta_2^2 \alpha^4 k^4 \sigma_{z0}^4 + \frac{1}{6} (\eta_1 \alpha - 6\eta_3 \alpha^3)^2 k^6 \sigma_{z0}^6 \right]$$

- Distortions manifest as a ramped current profile.
- To compute final bunch length we can include emittance growth in envelope analysis

$$\epsilon_{zz'}^2 = \epsilon_{z_0 z'_0}^2 + \epsilon_{RF}^2$$

$$\sigma_{zf} = \frac{1}{\sqrt{\frac{1}{\sigma_{z0}^2} + \frac{\sigma_{z0}^2}{\epsilon_{zz'}^2 f^2}}} \approx \frac{f \epsilon_{z,z'}}{\sigma_{z0}}$$

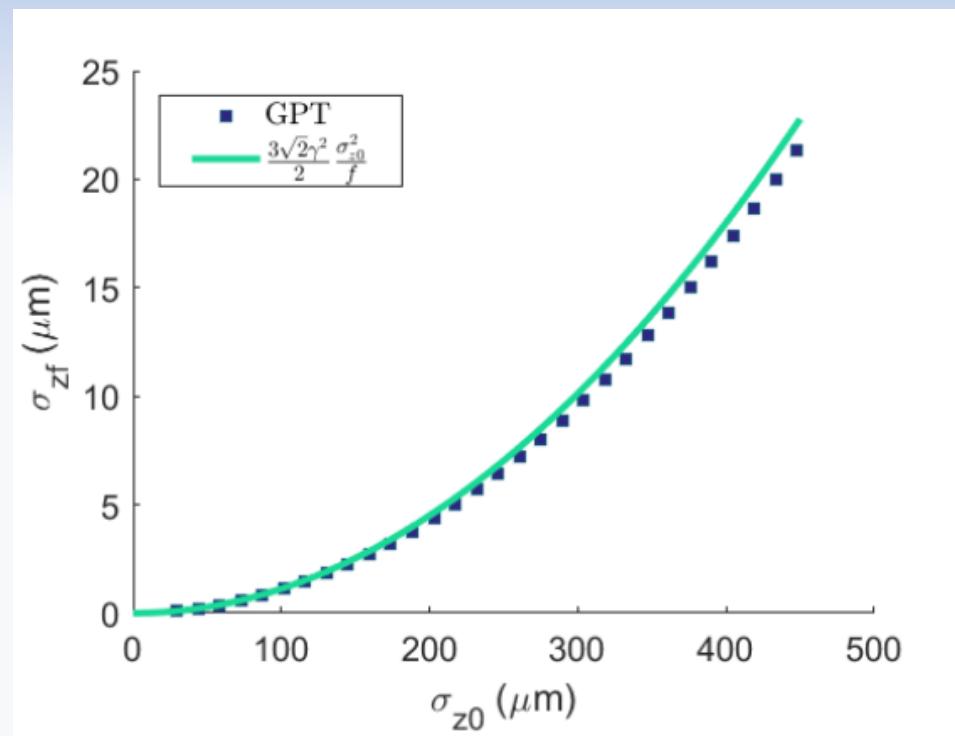


Analytical estimate of bunch length from RF non-linearities (Ignoring space charge).

At few MeV energies, bunch length dominated by quadratic distortions at the ballistic focus

$$\sigma_{zf} \propto \frac{\gamma^2 \sigma_{z0}^2}{f}$$

- Smaller longitudinal emittance enables better compression.
- Few electron pulses (stroboscopic) can achieve attosecond durations.
- What about single shot? Space charge limits?



Limitations set by space charge (linear and nonlinear)

- Space charge at first order is included using longitudinal perveance

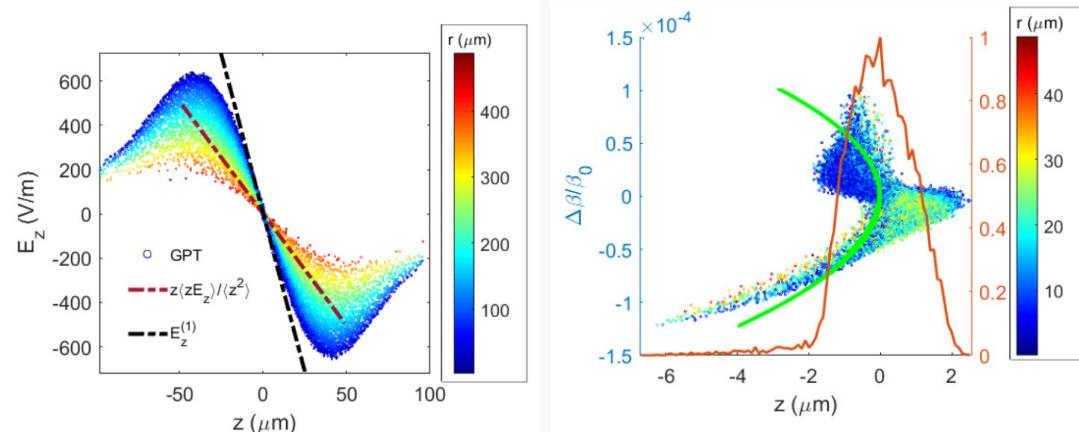
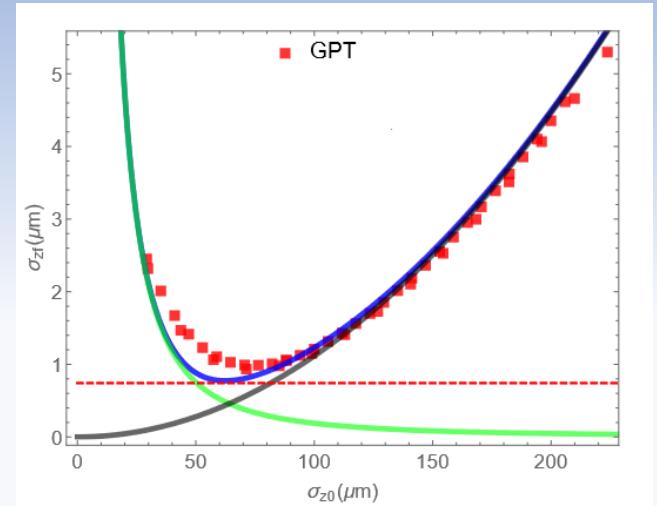
$$\sigma_z'' = \frac{K_L}{\sigma_z^2} + \frac{\epsilon_{z,z'}^2}{\sigma_z^3}$$

- Non-linear space charge forces: e.g., Gaussian induces emittance grows up to ballistic focus. Effect amplified for shorter initial bunches.

Fine tune initial bunch length for optimum compression—Optimal trade off between space charge and RF effects-- Ideal working point.

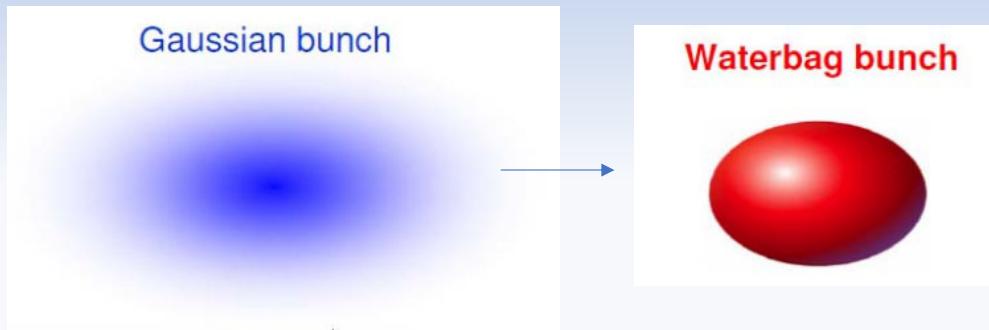
fixed focal length optimum $\sigma_{zf} \propto Q^{1/2} \gamma^{-3/2}$

Emittance growth sensitive to electron 3D distribution and its evolution.

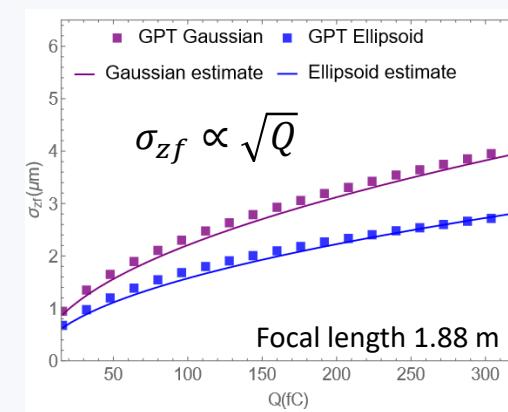
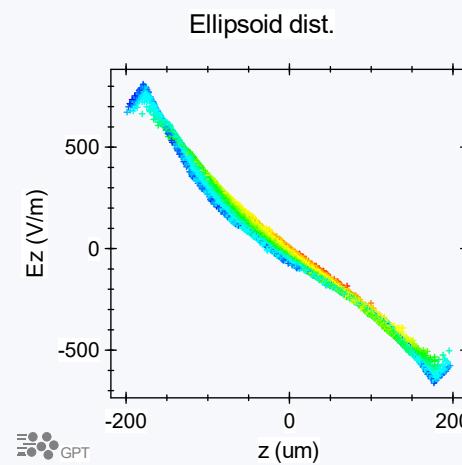
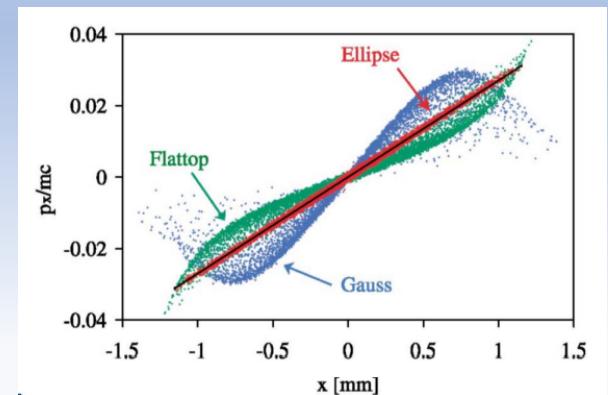


Electron beam shaping

- Shaping cathode drive laser allows control on initial configuration space (i.e. x,y,z distribution).



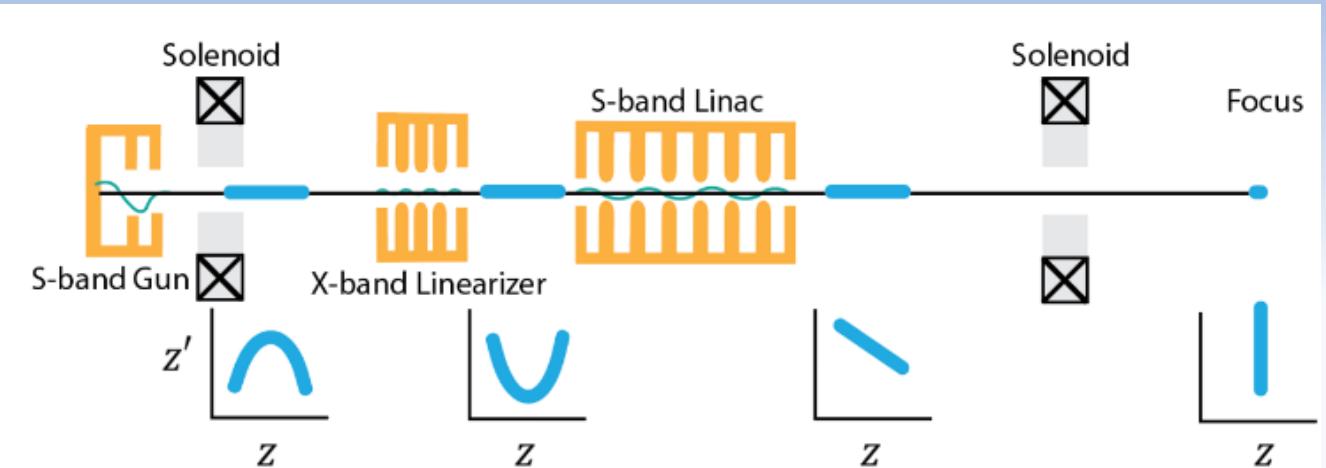
- Best case scenario, uniformly filled ellipsoid.
Stretch laser pulse and use long parabolic temporal profile.
- Linear space charge forces → linearly correlated phase space → shorter final bunch lengths.



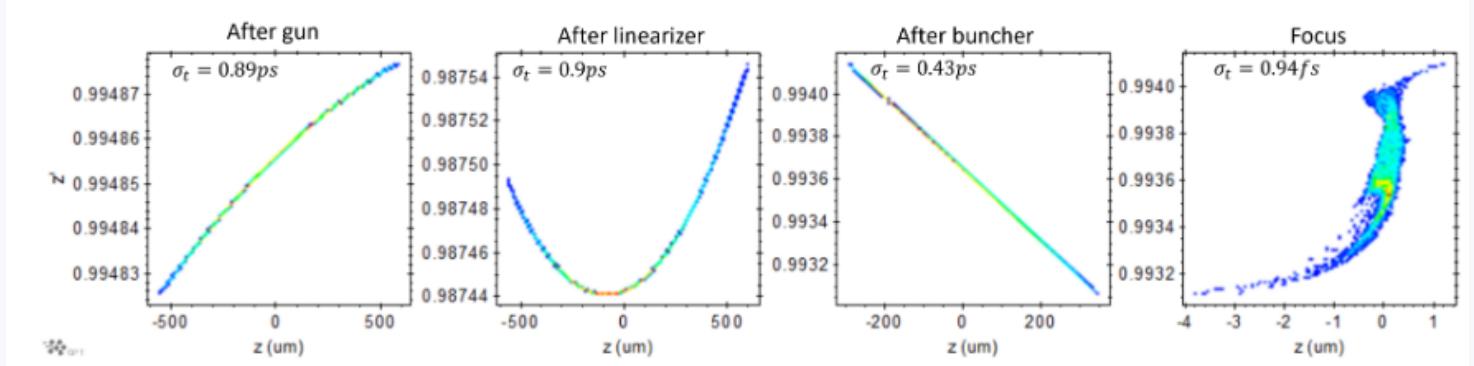
Higher harmonic compensation of emittance growth

- Deceleration in 3rd harmonic cavity removes second order curvature.

$$\alpha_x = 2 \frac{|\eta_2|}{\eta_1} \left(\frac{k}{k_x} \right)^2 \alpha^2$$

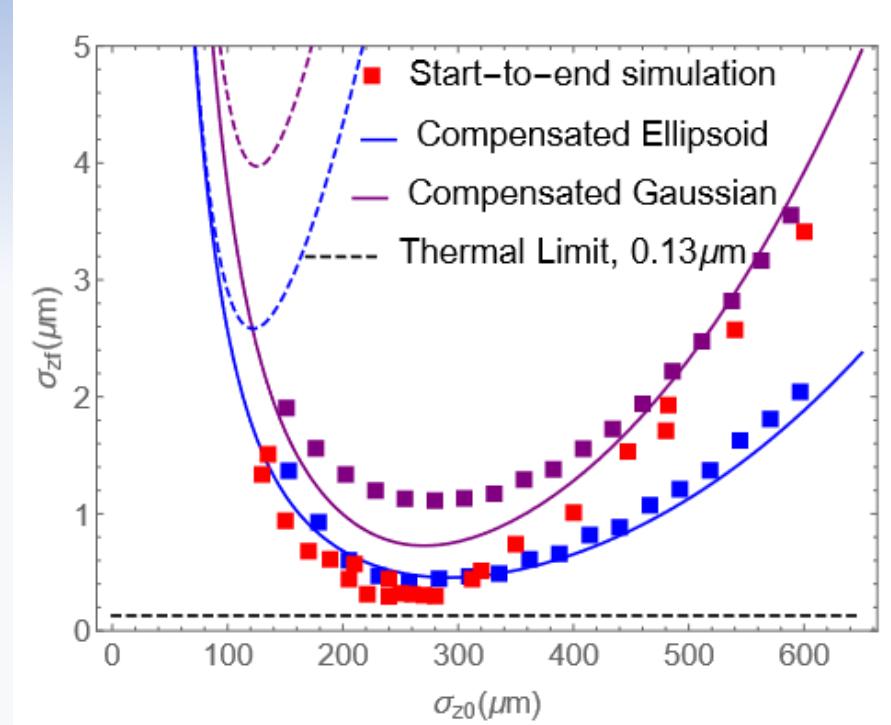
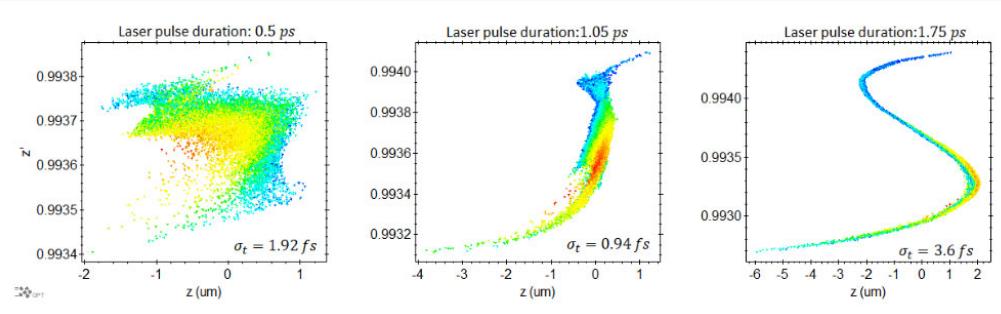


- Third order removal possible as well properly tuning Linac gradient and phase
- Compensation allows longer bunch lengths to be used to relax space charge effects.



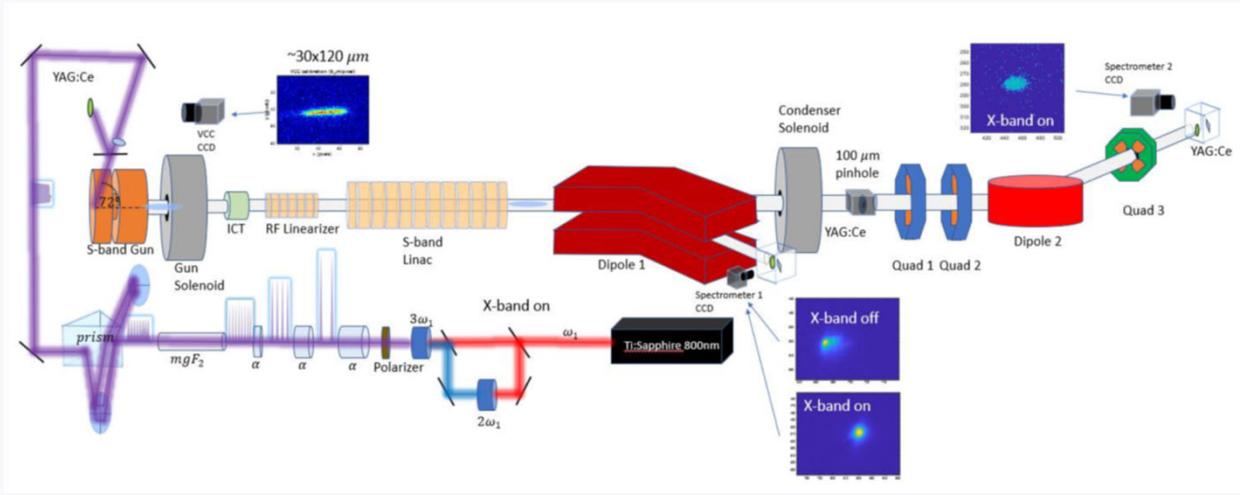
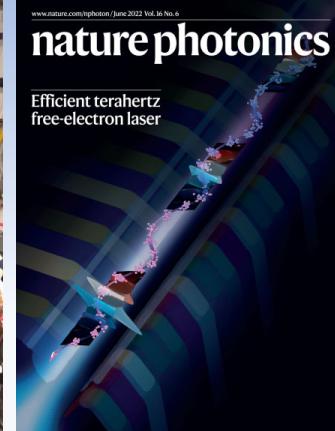
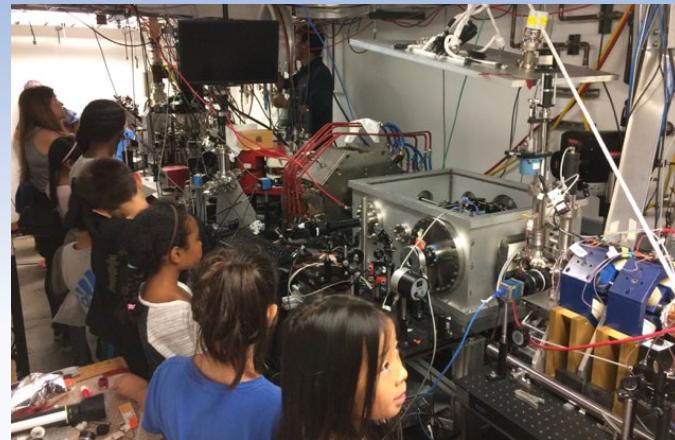
Start-to-end simulations for UCLA PEGASUS

- 10^6 electrons – 160 fC bunches
- Shortest bunches achieved when 2nd and 3rd order RF non-linearities are compensated, and electronic distribution is optimally shaped into uniform ellipsoid.
- Approaching final limit -- set by initial beam thermal emittance.



High peak brightness beam research at UCLA Pegasus Laboratory

- 3-14 MeV student-run university-size accelerator beamline optimized for sub-pC beams
- High brightness electron beams:
 - High resolution ultrafast electron diffraction
 - Single shot imaging / UEM
 - THz acceleration
 - High efficiency THz FEL



J Maxson, et al., Direct measurement of sub-10 fs relativistic electron beams with ultralow emittance. *Phys. Rev. Lett.*, **118**, 154802 (2017)

D. Cesar, et al., Demonstration of single-shot picosecond time-resolved MeV electron imaging. *Phys. Rev. Lett.*, **117**, 024801 (2016)

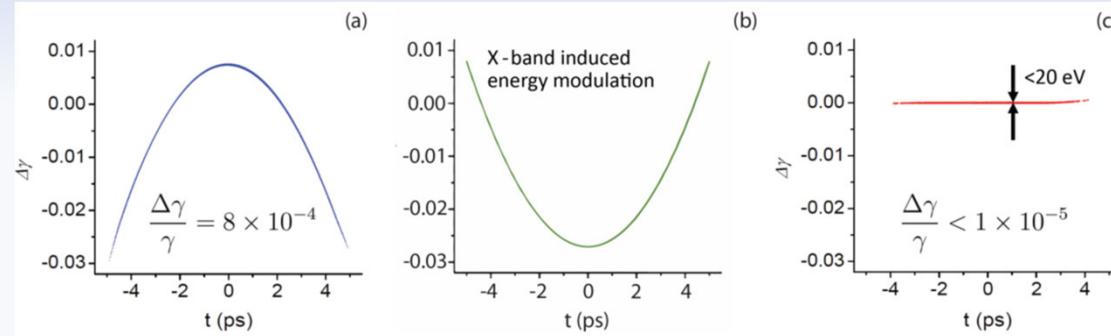
E. Curry et al. Meter-scale THz-driven acceleration of a relativistic electron beam. *Phys. Rev. Lett.* **120**, 094801 (2018)

A. Fisher, Y. Park et al. High efficiency single pass THz FEL. *Nature Photonics*, **16**, 441 (2022)

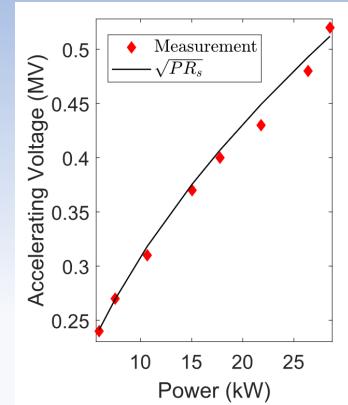
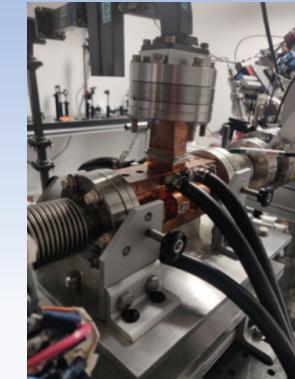
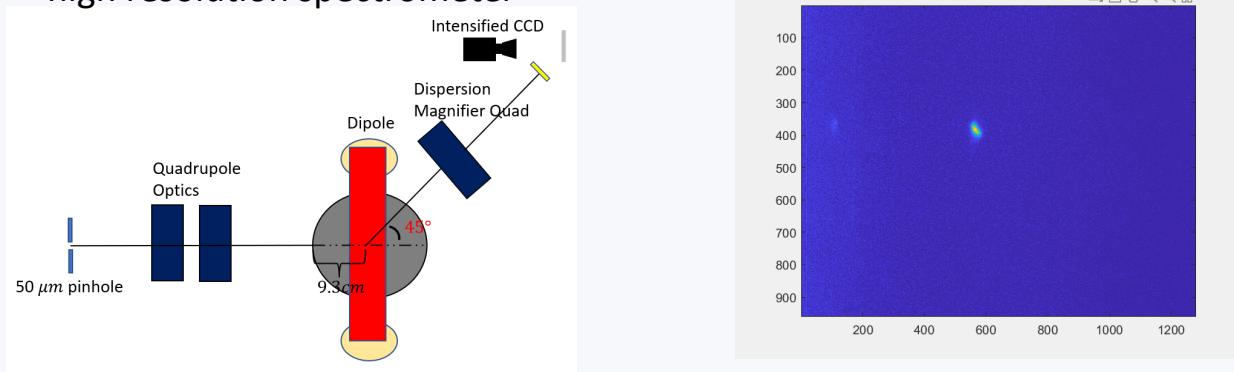
Demonstration of X-band LPS linearization

- Install and commission X-band cavity
- Removal of non-linearities in LPS for energy spread minimization as well.
- Application to minimize chromatic aberrations in ps-time resolved MeV TEM

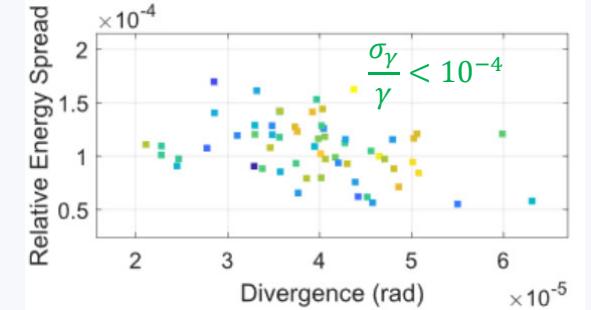
P. Denham, et al., "X-Band Harmonic Longitudinal Phase Space ... presented at the NAPAC'22



high resolution spectrometer

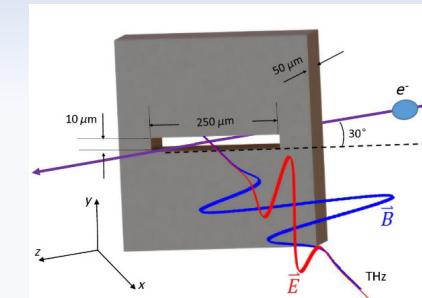
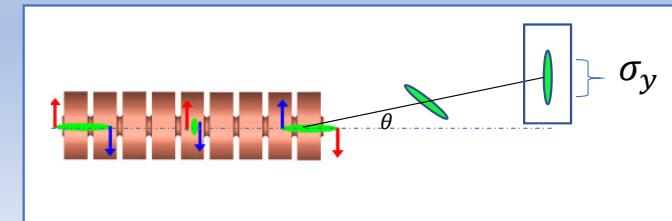


$\beta - \text{tron}$ contribution subtraction

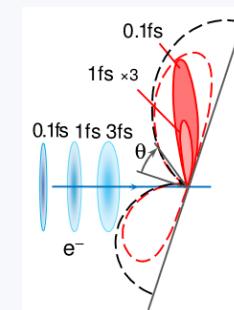


Options for sub-fs electron bunch measurements

- RF techniques.
 - RF deflecting cavity--Resolution limited by cavity frequency and screen resolution.
- THz streaking
 - Sensitive to timing jitter
 - Limited availability of THz power
- Laser techniques.
 - Laser Induced energy modulations in DLA or bend magnets?
- Radiative techniques.
 - Optical transition radiation—coherence factor suppressed by transverse extent of the bunch.
 - Different geometries enhance optical coherence.



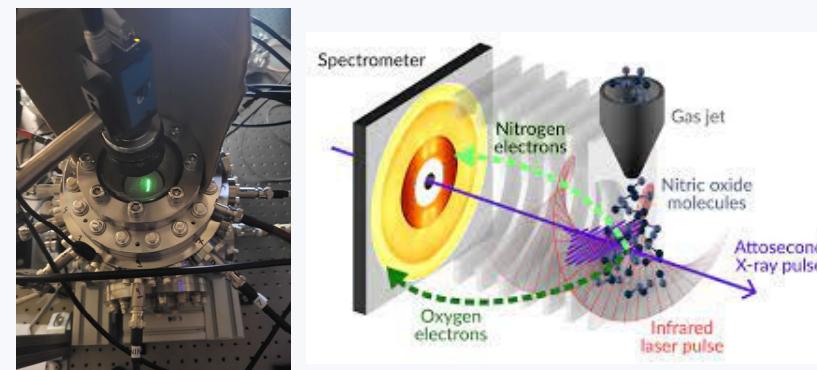
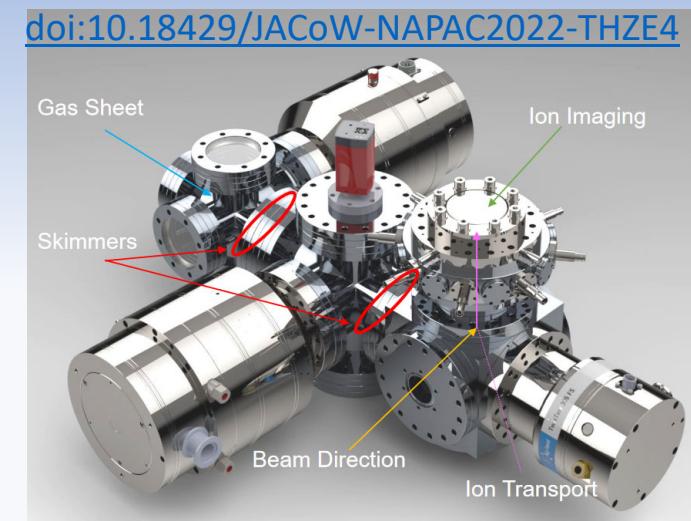
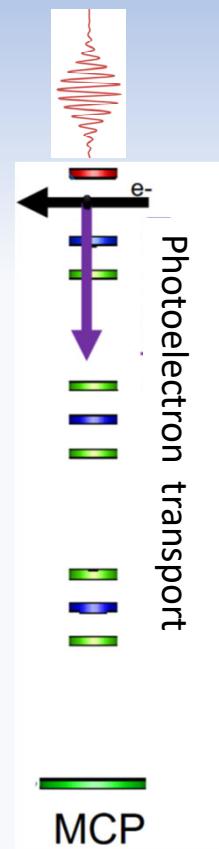
L. Zhao, et al. Phys. Rev. X 8, 021061



M.V. Tsarev and P. Baum 2018
New J. Phys. **20** 033002

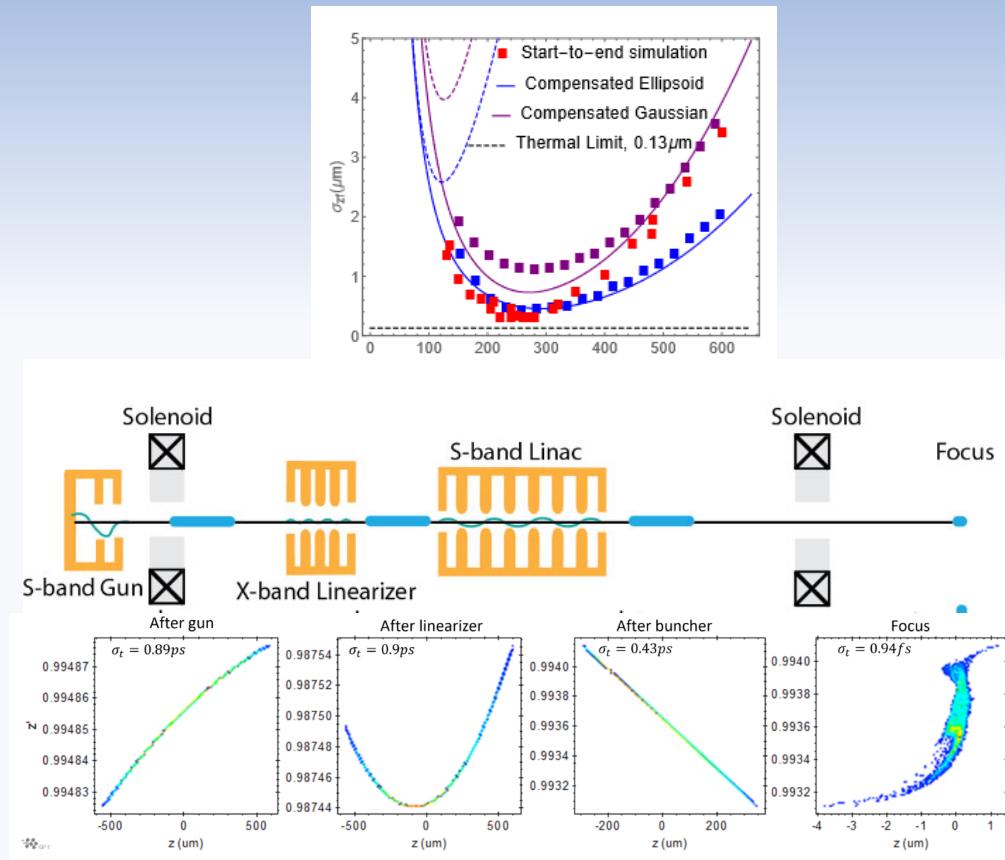
Ionized electron Velocity Map Imaging

- Learn a lesson from attosecond X-ray sources
- Set e-beam ballistic focus at gas sheet plane.
- Impact ionization of gas sheet.
- Laser streaking of “knocked” electrons.
- Ion microscope tuned to image velocity space at interaction point.
- Image time dependent streak profile—signature of e- beam duration.
- What is the signature exactly? What dilutes the signal dependence on bunch duration?



Summary: Ultrashort Beam generation at PEGASUS.

- Single shot UED temporal resolution determined by the bunch length at ballistic focus.
- Competing non-linear effects limiting the shortest achievable bunch length.
- Overcome non-linearities with laser shaping of electron beam and higher harmonic compensation.
- Bunch lengths in attosecond to sub-fs regime (for 10-100 fC charge) achievable.
- Implemented X-band and took measurements of LPS linearization—reduced energy spread.
- Seeking novel bunch length diagnostic to measure sub-fs bunches. Developments underway.



Laser Induced Energy Modulations in a Dipole

Plane wave Interaction

$$\vec{r}(t) = R(\cos(\omega_c t)\hat{x} + \sin(\omega_c t)\hat{z})$$

$$\vec{v}(t) = R\omega_c(-\sin(\omega_c t)\hat{x} + \cos(\omega_c t)\hat{z})$$

$$\vec{E} = \hat{x}E_0 \cos(k(z - ct) + \psi)$$

$$\phi = \omega_c t$$

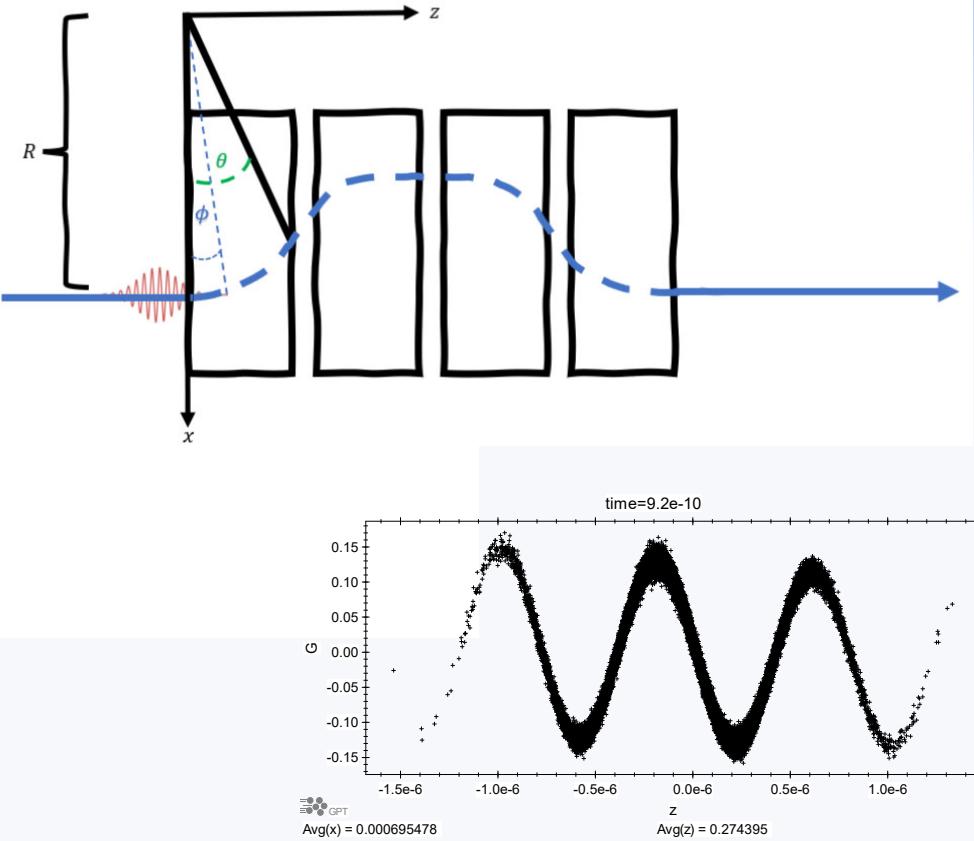
Neglecting interaction between laser and beam prior to chicane entrance

$$\frac{dE}{dt} = \vec{F} \cdot \vec{v} \rightarrow \frac{d\gamma}{dt} = \frac{qE_0}{mc^2} v_x \cos(k(z - ct) + \psi)$$

$$\frac{d\gamma}{dt} = -\frac{qE_0 R \omega_c}{mc^2} \sin(\phi) \cos\left(kR \left(\sin(\phi) - \frac{\phi}{\beta}\right) + \psi\right)$$

$$\therefore \frac{d\gamma}{d\phi} = -\frac{qE_0 R}{mc^2} \sin(\phi) \cos\left(kR \left(\sin(\phi) - \frac{\phi}{\beta}\right) + \psi\right)$$

$$\therefore \Delta\gamma = \frac{qE_0 R}{mc^2} \int_0^{\phi_{max}} \sin(\phi) \cos\left(kR \left(\sin(\phi) - \frac{\phi}{\beta}\right) + \psi\right) d\phi$$



How do we preserve temporal structure in transport?