



Plasma Wakefield Accelerator-Based Low Emittance Muon Source

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AAC'22

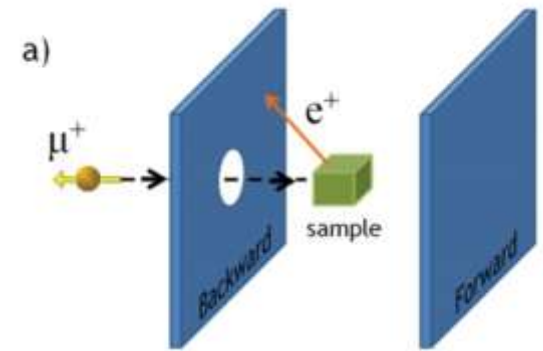
November 6-11, 2022



Promise of Muons

- μ SR for solid state research
- Screening NNSA/DARPA applications
- g-2 and mu2e experiments

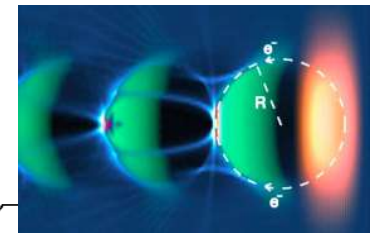
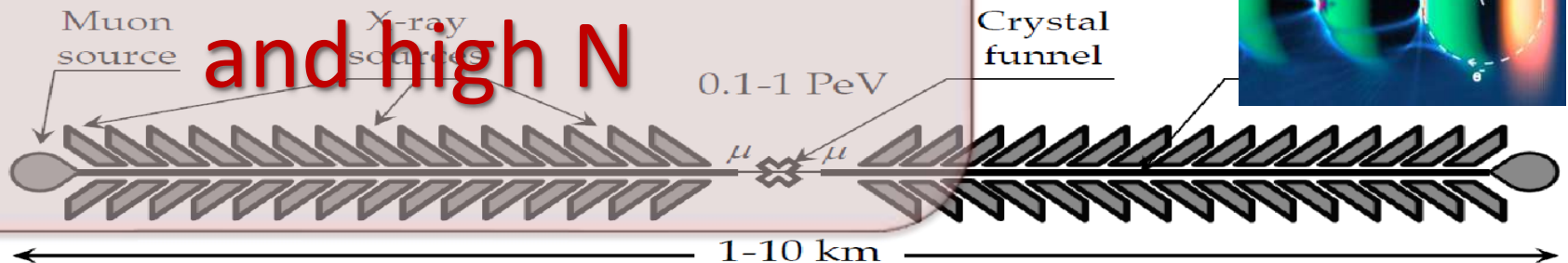
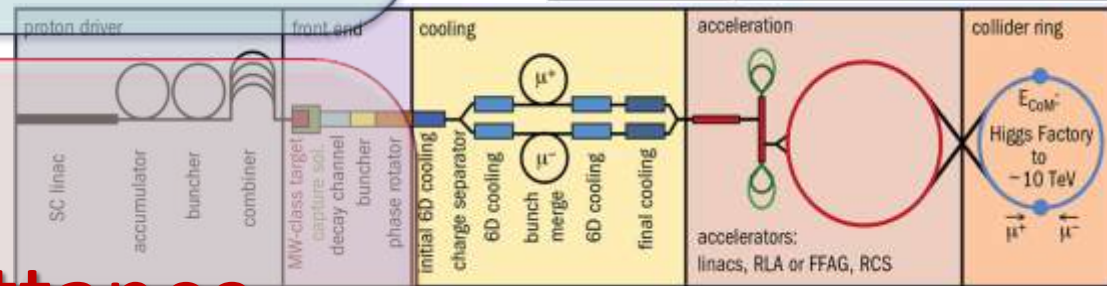
Intensity is the key



- Circular muon colliders
- Linear & crystal muon colliders

Need low emittance

and high N

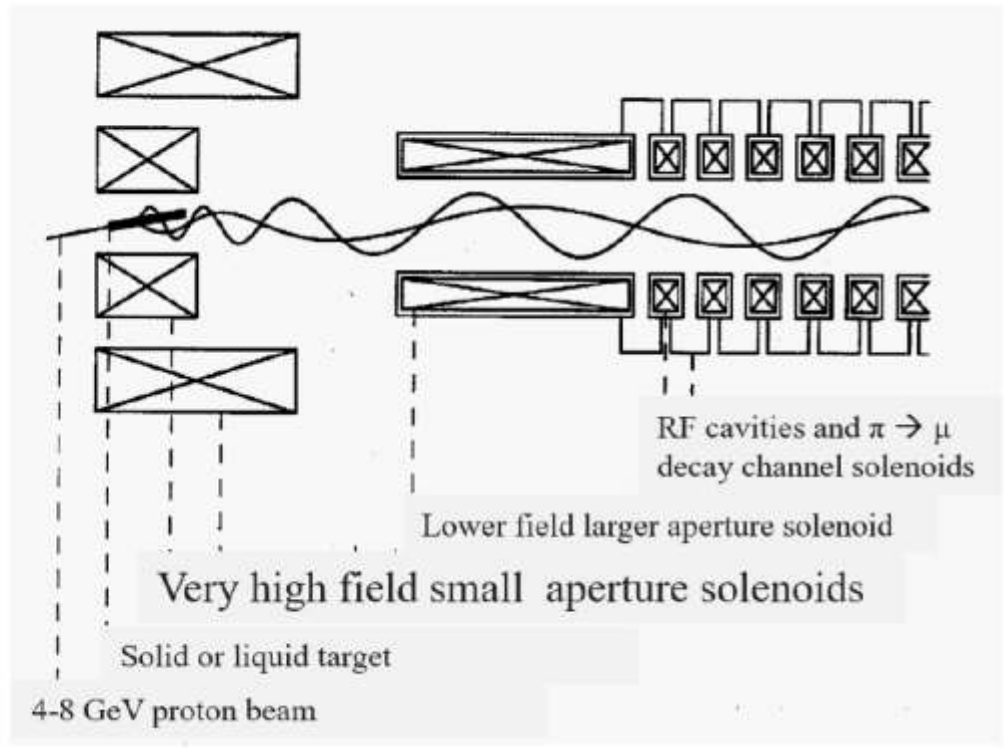
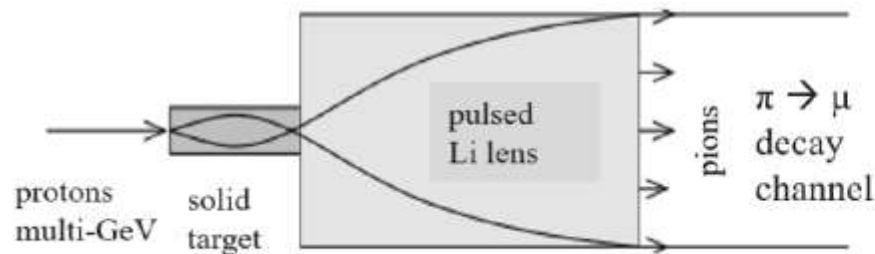
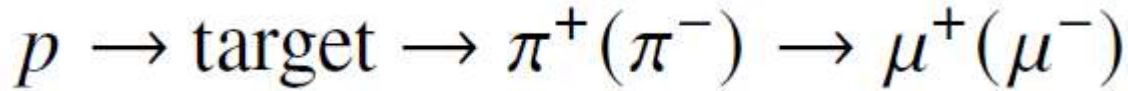


Challenges

I. Muon are unstable

- 2.2 us \rightarrow almost no time for collection, cooling and acceleration... all should be superfast
- Muons are secondary/tertiary
 - Need high intensity primary (depends on efficiency)
 - Come in large angular spread $\sim m_{\text{muon}}/E_{\text{primary}}$

Traditional Muon Sources



pulsed short-focus Lithium lens

$$\theta \simeq m_{\pi} / P_{\pi}$$

FNAL g-2 experiment:

116 kA \rightarrow 232 T/m

2 GeV p^+ \rightarrow 3.1 GeV muons

Eff $2e-7$ (dE/E $\sim 1\%$)

$O(30 \text{ T})$ target solenoid

Future 10-14 TeV muon collider/MAP:

About 0.1 muon/proton

Large emit $\sim 10^{-3}$ rad

ionization cooling channel $\epsilon \sim 25 \mu\text{m}$

LEMMA

M. Boscolo, M. Antonelli, et al
NIM A **807** (2016) 101–107

Low emittance muon pair

at threshold $e^+e^- \rightarrow \mu^+\mu^-$

need 45 GeV positrons

Findings wrt MAP scheme:

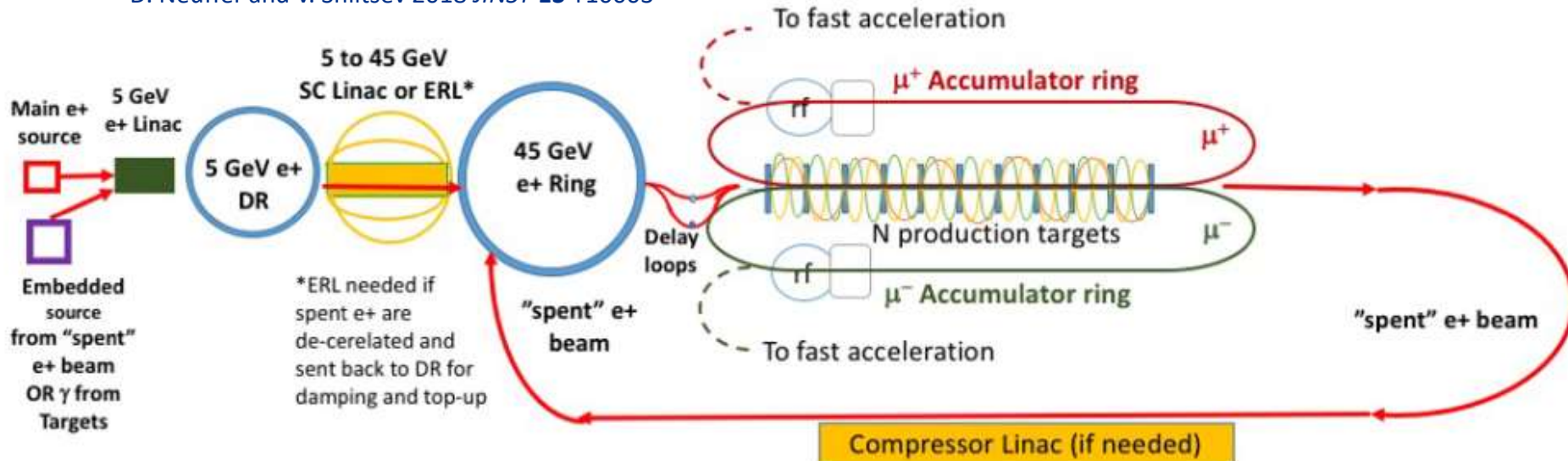
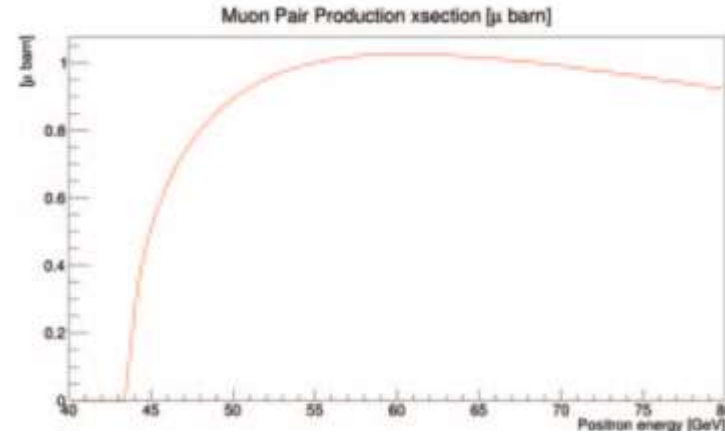
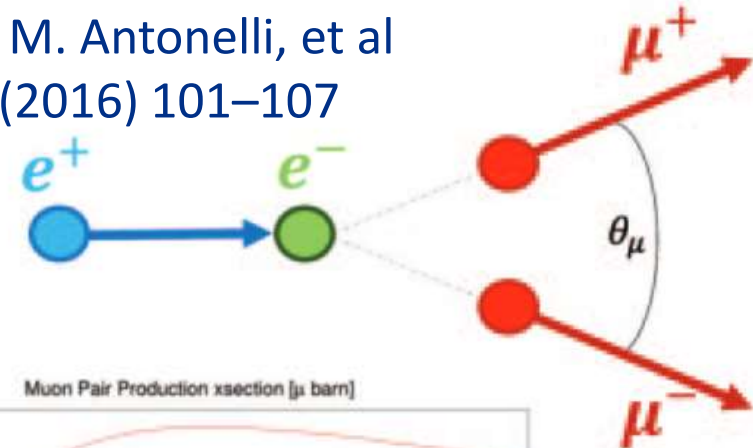
Emittance $\sim 1/500$

Intensity $\sim 1/10,000$

Luminosity $\sim 1/1000$

Cost +30%, Site power $\sim x2$

D. Neuffer and V. Shiltsev 2018 *JINST* **13** T10003



Recent Idea

A muon source based on plasma accelerators

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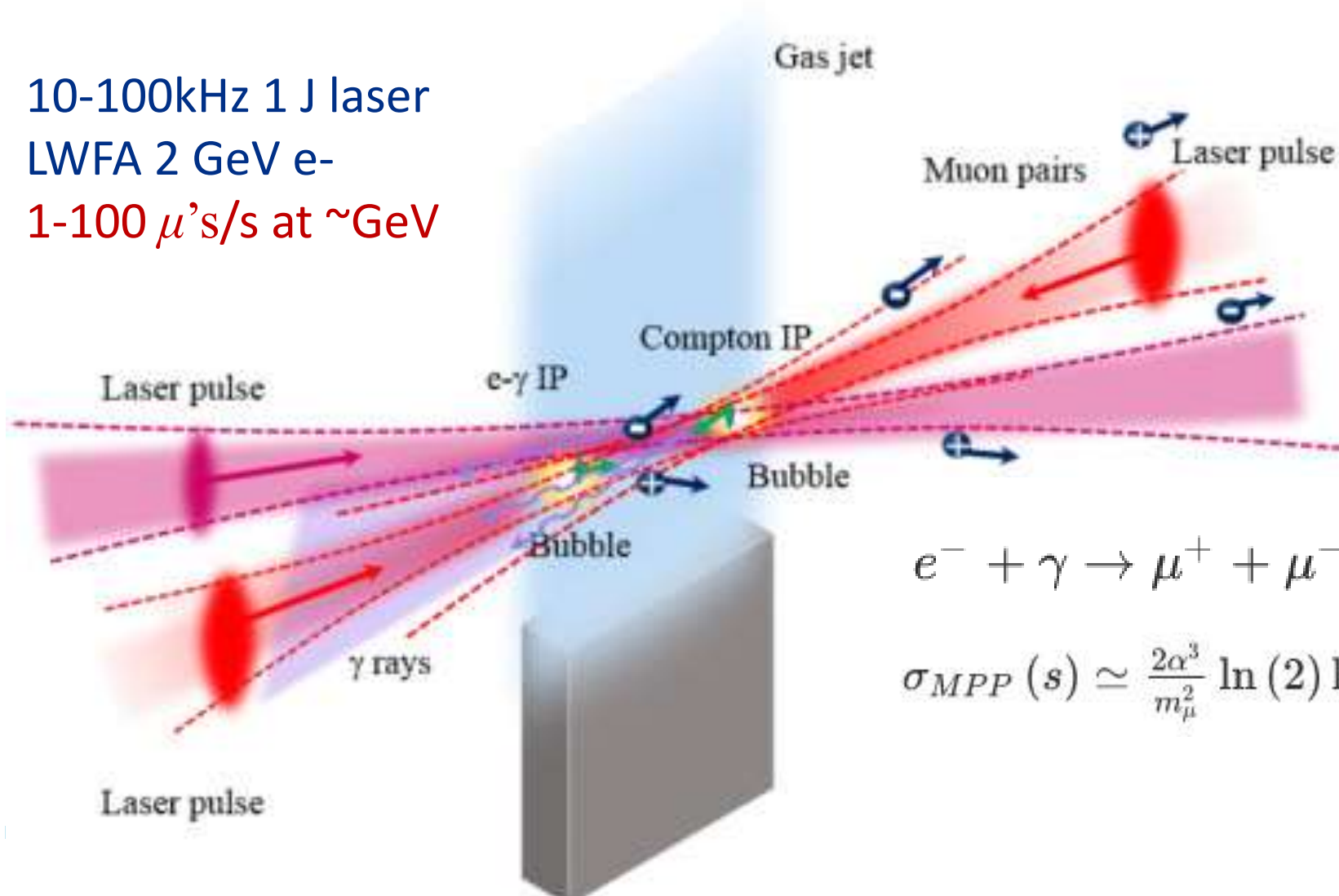
- e-gamma collider source

NIM A **909** (2018) 309–313

10-100kHz 1 J laser

LWFA 2 GeV e-

1-100 μ 's/s at \sim GeV



$$\sigma_{MPP} (s) \simeq \frac{2\alpha^3}{m_{\mu}^2} \ln(2) \ln\left(\frac{s}{m_e^2}\right)$$

New (Another) Approach

I. Employ Plasma-Wakefields (PWA cells)

- They are very strong transversely – can capture lots of muons
- They are very strong longitudinally – very fast acceleration to high $\gamma = E/mc$ and long muon lifetime

II. Presumably compact

EM Fields in Plasma Waves

Eg in the bubble regime, forces:

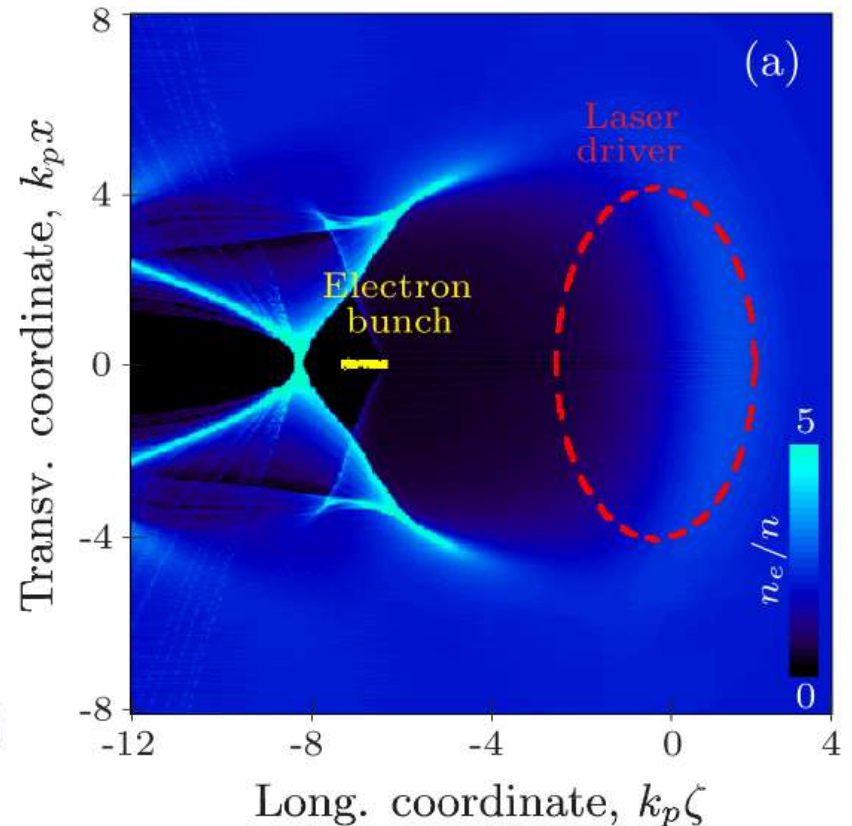
$$F_z = eE_0 k_p \zeta / 2$$

$$F_r = e(E_r - B_\theta) = -eE_0 k_p r / 2$$

where

$$E_0 = mc^2 k_p / e, \text{ or } E_0 [\text{V/m}] \approx 96 (n [\text{cm}^{-3}])^{1/2}$$

Electron plasma density

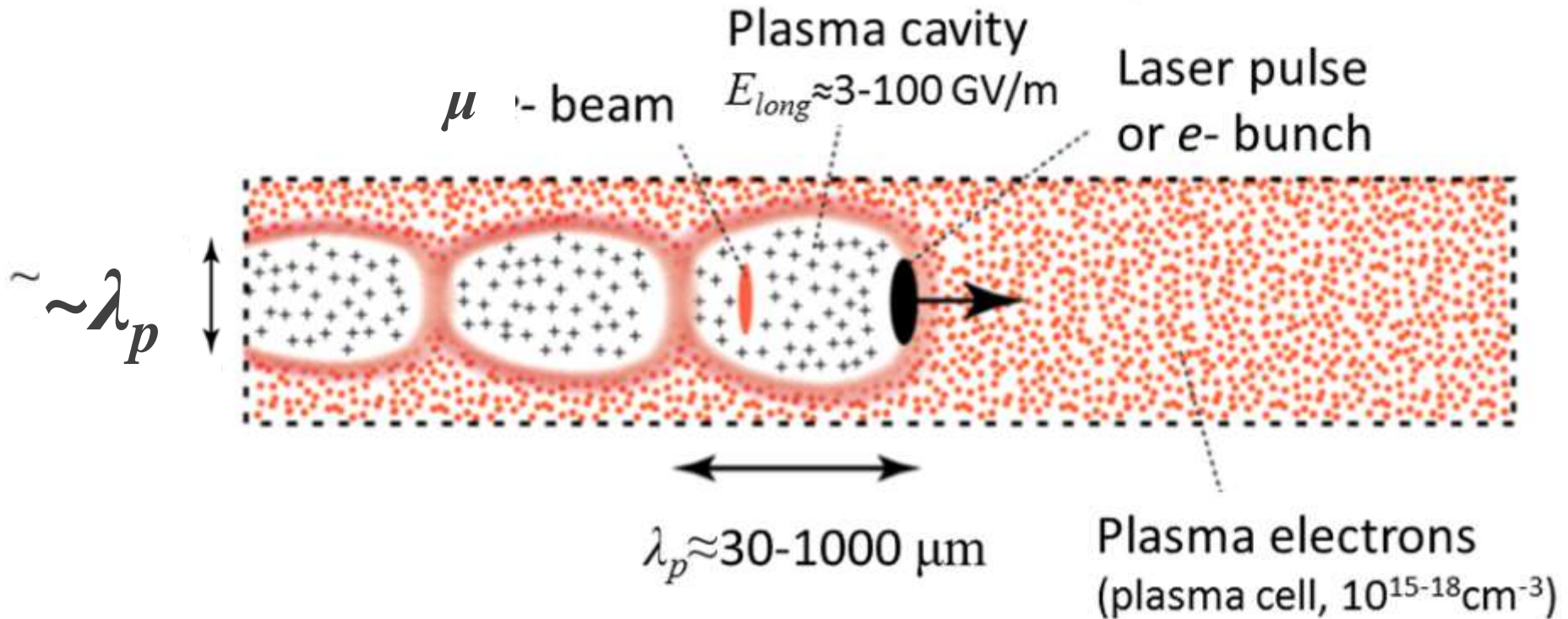


That's exactly what's needed for a muon source

Conceptual Scheme

- **Muons are born in the plasma WFA channel**
 - Either via photoproduction (need $O(1+ \text{ GeV})$ e- beam)
 - Or via pion decay
- Muons originally have large angular spread and small transverse dimension:
 - $\theta \approx m_\mu / E_\mu$
 - r smaller than the radius of the bubble $\sim 3/k_p = \lambda_p/2$
- **Muons get quickly accelerated** in the plasma while being super-strongly focused by the focusing fields of the plasma:
 - Come out with small emittance
 - Come out with very high energy $O(10 \text{ GeV})$

Plasma WFA Muon Source Concept



$$\lambda_p = c/\omega_p \approx 1 \text{ mm} \times \sqrt{10^{15} \text{ cm}^{-3} / n_0}$$

Generation of Muons: Two Schemes

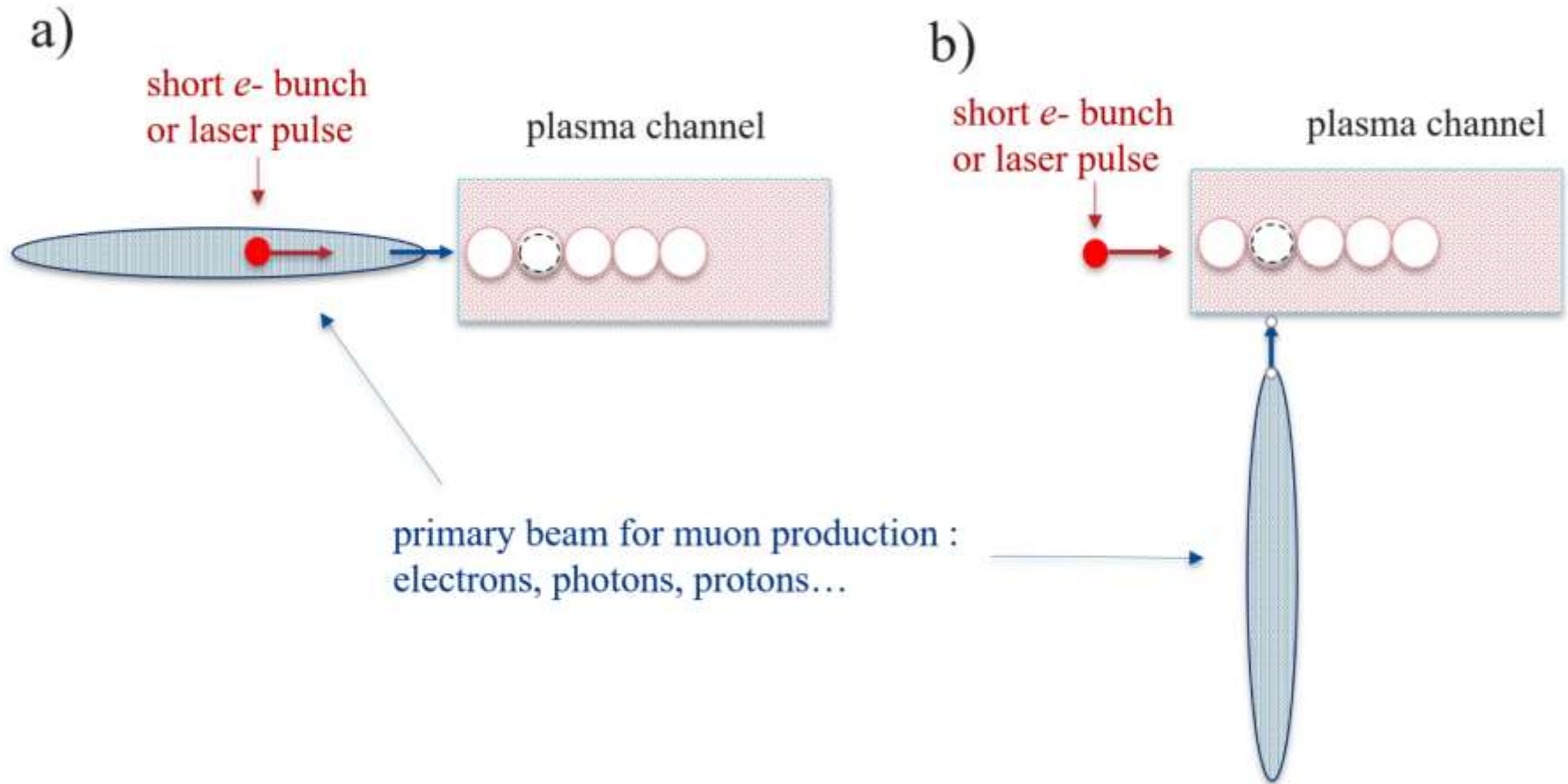


Figure 3. Muon PWA source: possible schemes for simultaneous injection of plasma wakefield drivers (laser pulses or short electron bunches) and primary beams needed for muon generation: a) collinear; b) orthogonal. Muons are produced in interaction of the primary beam with plasma, or, if suitable, with specially arranged dedicated target at the intersection of the beam with the plasma channel — see dashed circle.

Equations of Motion

$$\frac{dp_z}{dt} = F_z$$

$$\gamma_\mu = 1 + z\kappa eE_0/m_\mu c^2 \approx z\kappa k_p (m_e/m_\mu)$$

$$\frac{dp_r}{dt} = m_\mu c^2 \frac{d}{dz} \left(\gamma_\mu(z) \frac{dr}{dz} \right) = -eE_0 \frac{k_p r}{2}$$

$$r(z) = c_1 J_0 \left(2 \sqrt{\frac{z\pi}{\kappa\lambda_p}} \right) + c_2 Y_0 \left(2 \sqrt{\frac{z\pi}{\kappa\lambda_p}} \right)$$

The betatron motion is fast

$$\beta_\mu = 1/k_\mu \approx \lambda_p \sqrt{2\gamma_\mu m_\mu/m_e}$$

So the amplitude drops adiabatically

$$|r| \approx r_{\max} \sqrt{\beta_\mu/\gamma_\mu} \propto r_{\max}/\gamma_\mu^{1/4}$$

Equations on Emittance

Max amplitude and angular spread define max acceptable emittance

$$r_{\max} \simeq \lambda_p / 2,$$

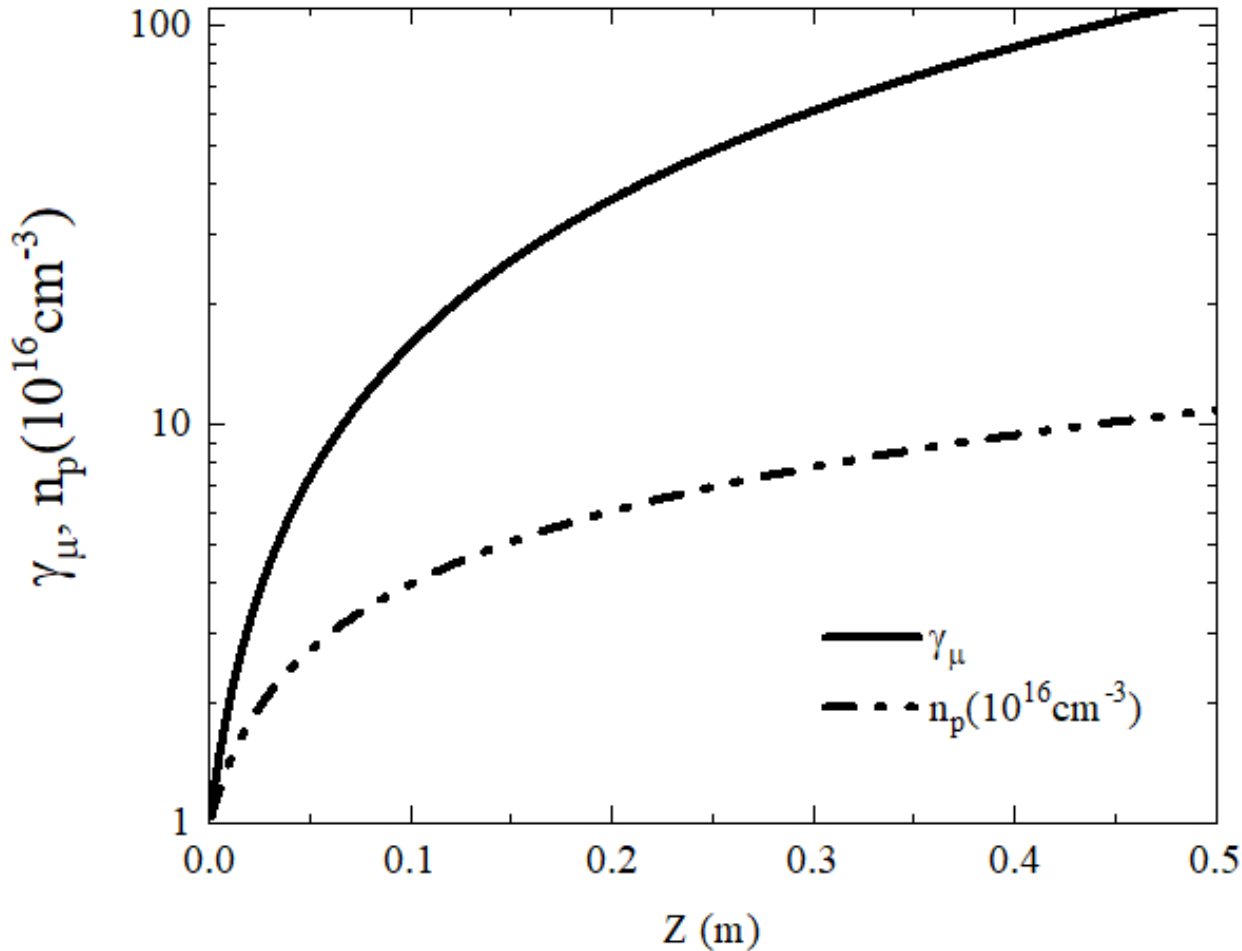
$$\theta_{\max} = r_{\max} / \beta_{\mu} = \pi \sqrt{\frac{m_e}{2\gamma_{\mu} m_{\mu}}} \approx \frac{0.15}{\sqrt{\gamma_{\mu}}}$$

$$\epsilon_{\mu}^{\max} = \gamma_{\mu} \theta_{\max} r_{\max} = \lambda_p \gamma_{\mu}^{1/4} \sqrt{\frac{m_e \pi^2}{8m_{\mu}}} \approx 0.078 \lambda_p$$

Assuming at birth $\gamma \sim 1$, if max accepted emittance is $\epsilon \sim 25 \mu\text{m}$ for $n = 10^{16} \text{ cm}^{-3}$, in that case the length of a 10 GeV source is about 1 m. For denser plasma $n = 10^{18} \text{ cm}^{-3}$ the max accepted emittance is $\epsilon \sim 2.5 \mu\text{m}$ and 10 GeV source is about 10 cm.

Plasma density rampup can make the source shorter without loss of the acceptance (see next slide)

Tapered Plasma Density Rampup



$$\gamma_p \approx (1+az)^{4/3}$$

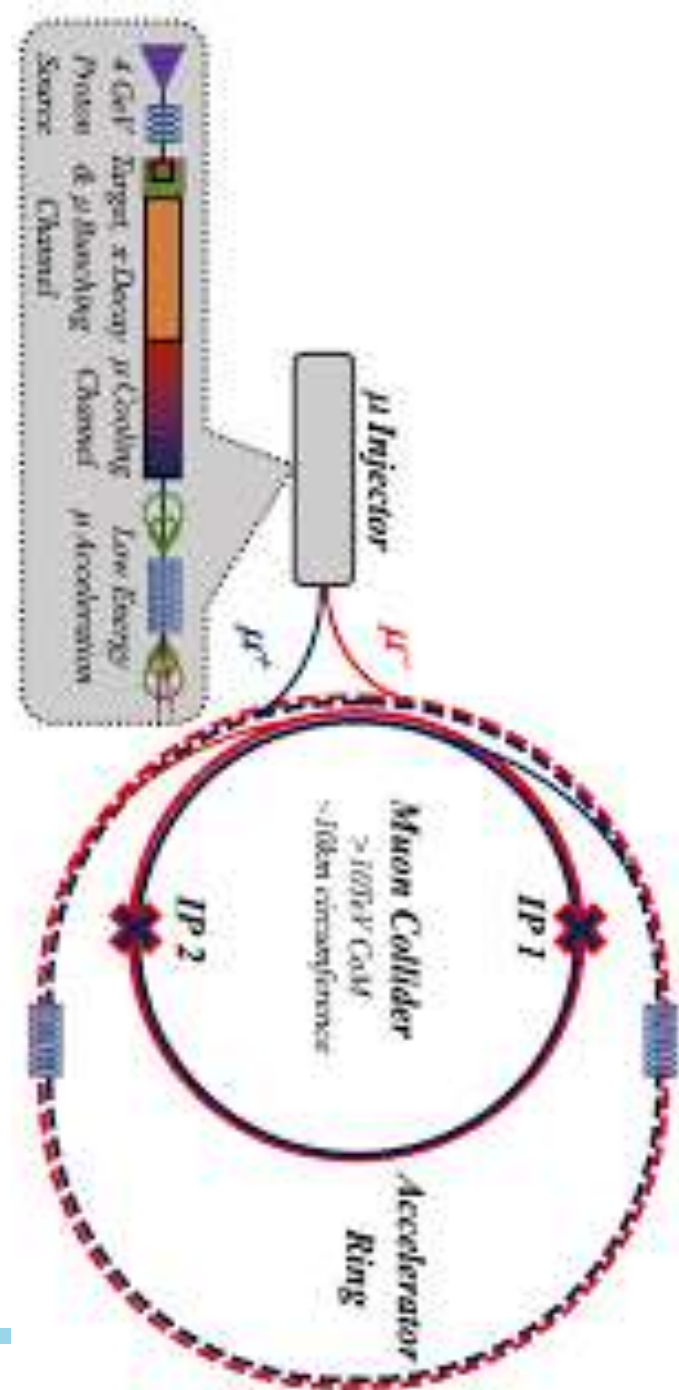
$$n_p = n_0 (1+az)^{2/3}$$

Plasma density and muon energy in tapered PWA-based 10 GeV muon source with normalized acceptance of $25 \mu\text{m}$ - corresponding to

$$\lambda_p^0 = 0.33 \text{ mm and } n_p^0 = 10^{16} \text{ cm}^{-3}$$

PWA μ Source MC

- May eliminate the most complex muon production and early acceleration part of multi-TeV muon colliders
- Very compact $O(1\text{ m})$ total
- $O(10\text{ GeV})$ beams with norm. emittance from few to tens μm
- Very large energy and angle acceptance
- *(solves so many problems of traditional muon ionization cooling scheme – from scattering and struggling to massive hardware needs, magnets and RF)*

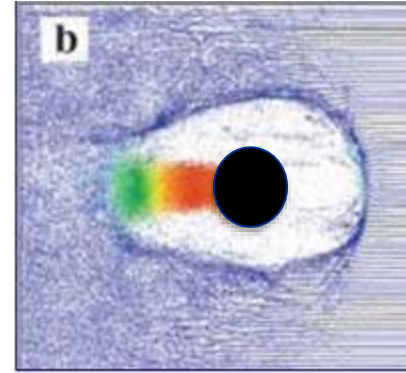


BUT : Challenges and Open Questions (1)

- Acceleration and focusing of μ^+ as challenging as for **positrons** in “traditional” LPWA and PWFA
- Collider application requires both small emittances and high intensity (ie brightness) μ beams – very hard, see next slide
- Plasma wakes are essentially low-Q, so **only few plasma periods/bubbles can be employed** → few short μ bunches
- Bubble wake will be loaded by **secondaries** (e^+e^- , pions, etc)

Muon Generation Challenges – Possible Schemes

- **Photoproduction:** need multi-GeV e^- beam, BH cross-section is $1/m_\mu^2$ small, $O(1\mu/\text{pulse})$, wake loading by e^+e^- pairs... solid target in the bubble?
- **e^+e^- annihilation:** even small cross-section
- **Proton-nucleon:** need 1-10 GeV protons, cross section is very high, but plasma density is low and pions need ~ 8 m to decay (in the wake)
- **Prompt muons:** either D or π^0 require $O(100-100)$ GeV p^+ , cross-section is high $O(\pi^0 \rightarrow \gamma\gamma)$ t plasma density is low
- **External injection:** **sub-mm short intense muon bunches!!!**



Can solid density object be inserted in the bubble?

Conclusions:

- Fast acceleration and strong focusing in plasma wakefields allow compact $O(10 \text{ cm})$ 10 GeV muon source
- If muons produced inside the bubble, then captured emittance is about $\approx 0.078\lambda_p$
- Difficulties (so far) : a) probably does not work for μ^+ b) muon generation schemes do not promise decent intensities (other schemes?); c) bubble wake will be loaded by secondaries
- (The only practical application might be) injection of mmxmmxmm muon bunches in the plasma WFA channel

Muons are the Particles of the Future!

Thanks for your attention!



- Special thanks to
 - *Philippe Piot for inviting me to submit this work to ICFA BD Newsletter 83 (JINST)*
 - *WG5 conveners for inviting me*
- ... see JINST Technical Report publication :

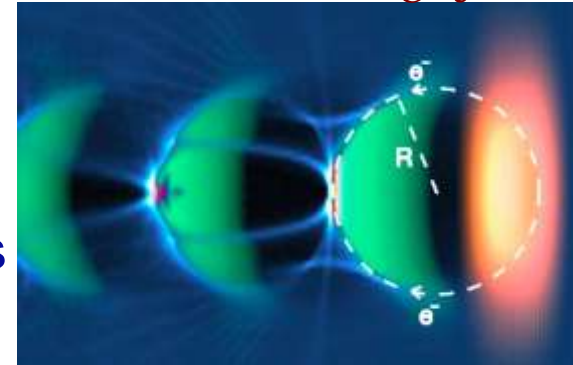
<https://iopscience.iop.org/article/10.1088/1748-0221/17/05/T05010/meta>

Back up slides

Linear e^+e^- Colliders

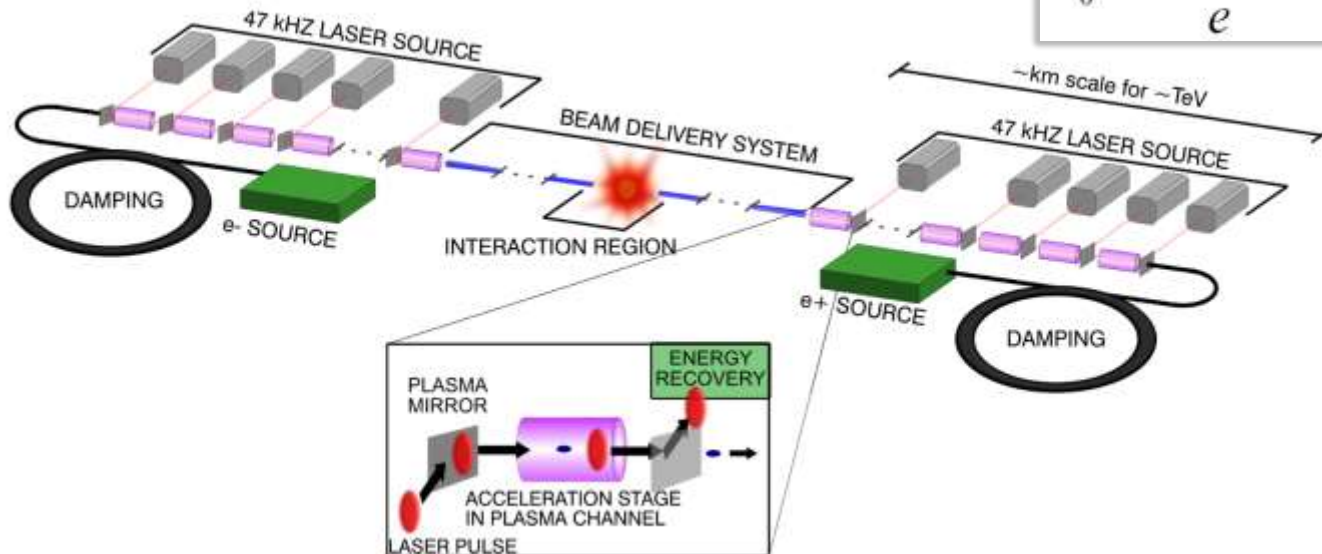
15 TeV e^+e^-
 100+ km
 10-15 km

Plasma sustains high fields



- Either RF acceleration (50-200 MeV/m) or wake-field acceleration in plasma (2-5 GeV/m)
- Major limitations:
 - 100% energy spread at IP (beamstrahlung)
 - One-time collisions ineffective $\rightarrow Lumi \sim P/(E\sigma)$
 - Very long/complex *Final Focus* to get nm IP size
 - Extreme sensitivity to nm jitters of linac elements
 - In plasma – ultra-strong focusing hurts staging, impossible(?) to accelerate positrons

$$E_0 = \frac{m_e c \omega_p}{e} \approx 100 \left[\frac{GeV}{m} \right] \cdot \sqrt{n_0 [10^{18} cm^{-3}]}$$

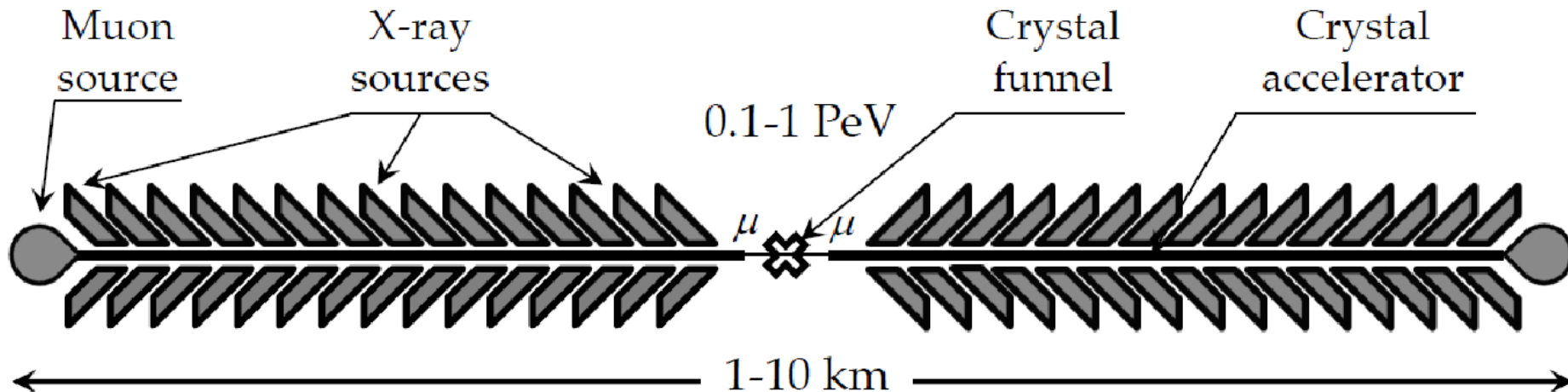
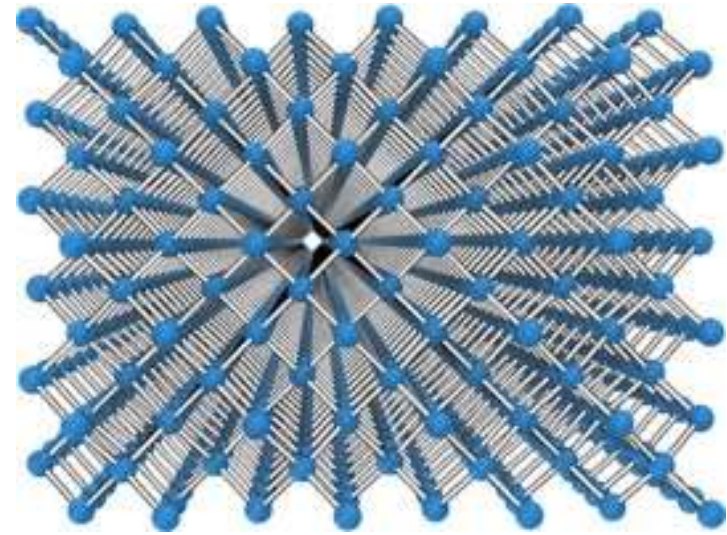


Takes about 1 GW electric power even for 15 TeV

Exotic Colliders

- Plasma-wakefield acceleration **and** channeling in structured media, eg CNTs or crystals (**only muons!!!**)
- Major advantages:
 - solid density \rightarrow 1-10 TV/m gradients
 - continuous focusing and acceleration (no cells, one long channel, particles get strongly cooled *betatron radiation*)
 - small size promises low cost
- *Lumi* $\sim 1/E^2$...totally unproven yet concept:
 - proof-of-principle experiment *E336* @ SLAC

$$E [\text{GV/m}] \approx 100\sqrt{n_0 [10^{18} \text{ cm}^{-3}]}$$



Xtal Collider

$n \sim 10^{22} \text{ cm}^{-3}$, $10 \text{ TeV/m} \rightarrow 1 \text{ PeV} = 1000 \text{ TeV}$

$n_{\mu} \sim 1000$, $n_B \sim 100$, $f_{rep} \sim 10^6$ $L \sim 10^{30-32}$

V.Shiltsev, Phys. Uspekhy 55 965 (2012)

