

# 20<sup>th</sup> Advanced Accelerator Concepts Workshop

## Near-Field CTR beam focusing and its application to Strong Field QED

Pablo San Miguel Claveria (LOA, IP Paris / X-GolP, IST Lisbon)

On behalf of:

E332 Collaboration (SLAC, UCLA, CU Boulder, U. Oslo, Stony Brook U., MPIK, CEA, LOA)

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Numerical simulations were performed using HPC resources from PRACE (Grant No. 2017174175) and GENCI-TGCC (Grants No. 2018- A0040507594, No. 2019-A0060510786, 2020-A0080510786 and 2021- A0100510786 to run PIC simulations.



# Overview

1. Near-Field Coherent Transition Radiation – modelling and simulations.
2. Analytical estimates and observables of NF-CTR beam focusing with FACET-II parameters.
3. Probing Strong-Field regime of QED via NF-CTR process.

PHYSICAL REVIEW LETTERS 126, 064801 (2021)

## Extremely Dense Gamma-Ray Pulses in Electron Beam-Multifoil Collisions

Archana Sampath,<sup>1</sup> Xavier Davoine,<sup>2,3</sup> Sébastien Corde,<sup>4</sup> Laurent Gremillet,<sup>2,3</sup> Max Gilljohann,<sup>4</sup> Maitreyi Sangal,<sup>1</sup> Christoph H. Keitel,<sup>1</sup> Robert Ariniello,<sup>5</sup> John Cary,<sup>5</sup> Henrik Ekerfelt,<sup>6</sup> Claudio Emma,<sup>6</sup> Frederico Fuza,<sup>6</sup> Hiroki Fujii,<sup>7</sup> Mark Hogan,<sup>6</sup> Chan Joshi,<sup>7</sup> Alexander Knetsch,<sup>4</sup> Olena Kononenko,<sup>4</sup> Valentina Lee,<sup>5</sup> Mike Litos,<sup>5</sup> Kenneth Marsh,<sup>7</sup> Zhan Nie,<sup>7</sup> Brendan O'Shea,<sup>6</sup> J. Ryan Peterson,<sup>6,8</sup> Pablo San Miguel Claveria,<sup>4</sup> Doug Storey,<sup>6</sup> Yipeng Wu,<sup>7</sup> Xinlu Xu,<sup>6</sup> Chaojie Zhang,<sup>7</sup> and Matteo Tamburini<sup>1,\*</sup>

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Ⓞ (Received 25 September 2020; revised 18 December 2020; accepted 8 January 2021; published 12 February 2021)

Sources of high-energy photons have important applications in almost all areas of research. However, the photon flux and intensity of existing sources is strongly limited for photon energies above a few hundred keV. Here we show that a high-current ultrarelativistic electron beam interacting with multiple submicrometer-thick conducting foils can undergo strong self-focusing accompanied by efficient emission of gamma-ray synchrotron photons. Physically, self-focusing and high-energy photon emission originate from the beam interaction with the near-field transition radiation accompanying the beam-foil collision. This near field radiation is of amplitude comparable with the beam self-field, and can be strong enough that

## Beam focusing by near-field transition radiation

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## Probing strong-field QED in ultrarelativistic beam-plasma collisions

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Ongoing progress in laser and accelerator technology opens new possibilities in high-field science, notably for the study of the largely unexplored strong-field QED regime where electron-positron pairs can be created directly from light-matter or even light-vacuum interactions. Laserless strategies such as beam-beam collisions<sup>1</sup> have also been proposed with the prospect of pushing strong-field quantum electrodynamics (SFQED) in the nonperturbative regime. Here we report on an original concept to probe SFQED by harnessing the interaction between an electron beam and a solid target. When a high-density, ultrarelativistic beam impinges onto an even denser plasma, the beam self fields are reflected at the plasma boundary: in the rest frame of the beam electrons, these fields can exceed the Schwinger field, leading to SFQED effects such as quantum nonlinear inverse Compton scattering and nonlinear Breit-Wheeler electron-positron pair creation. We show that such beam-plasma collisions can produce results similar to beam-beam collisions with the advantage of a much simpler experimental setup. This scenario opens the way to precision studies of strong-field QED, with measurable clear signatures in terms of gamma-ray photon and pair production, and thus is a very promising milestone on the path towards laserless studies of nonperturbative QED.

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See talk by D. Storey on Wed. 13:30 (WG 4+7) for first experimental results!

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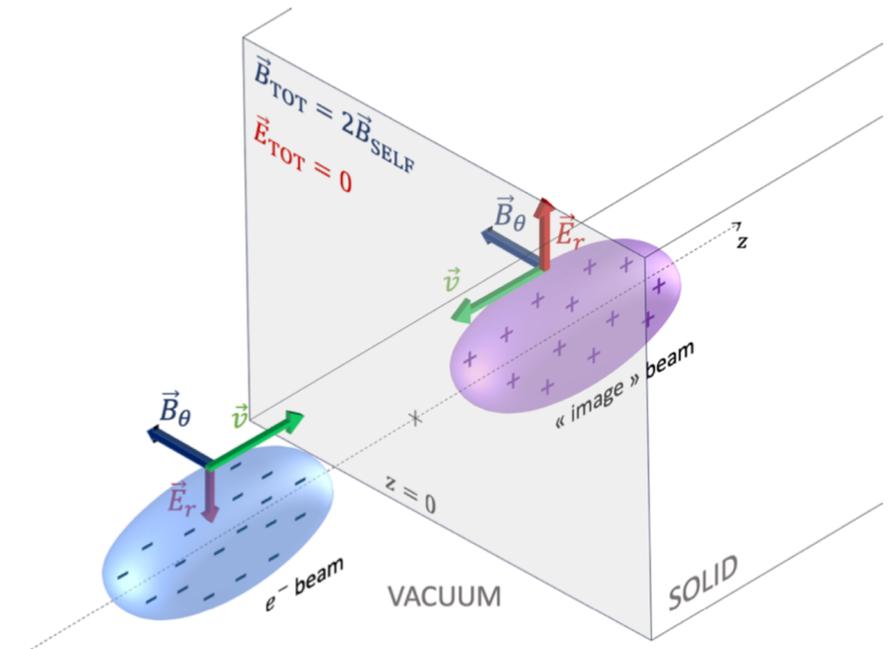
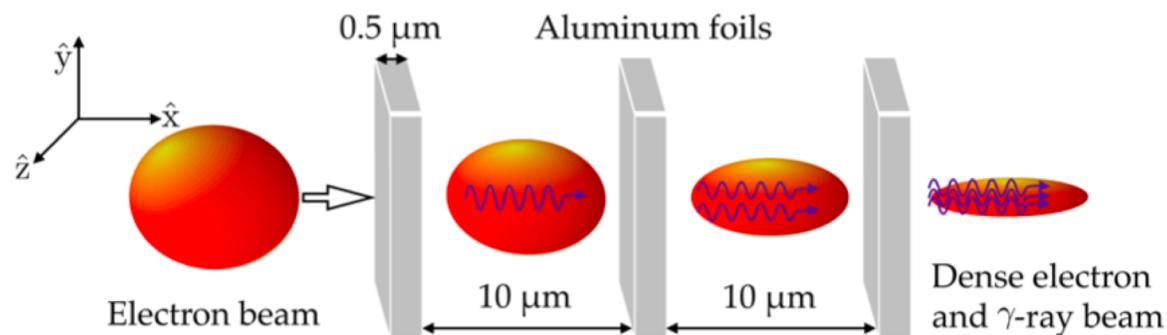
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# 1. Near-Field Coherent Transition Radiation – modelling and simulations

- Transition radiation emitted by a charged beam passing through a conducting foil has a net focusing effect on the beam.
- Initially suggested in 80's for beam transport<sup>1</sup> – image charges.
- Focusing element to reach solid-density electron beams and bright gamma-ray source<sup>3</sup>.

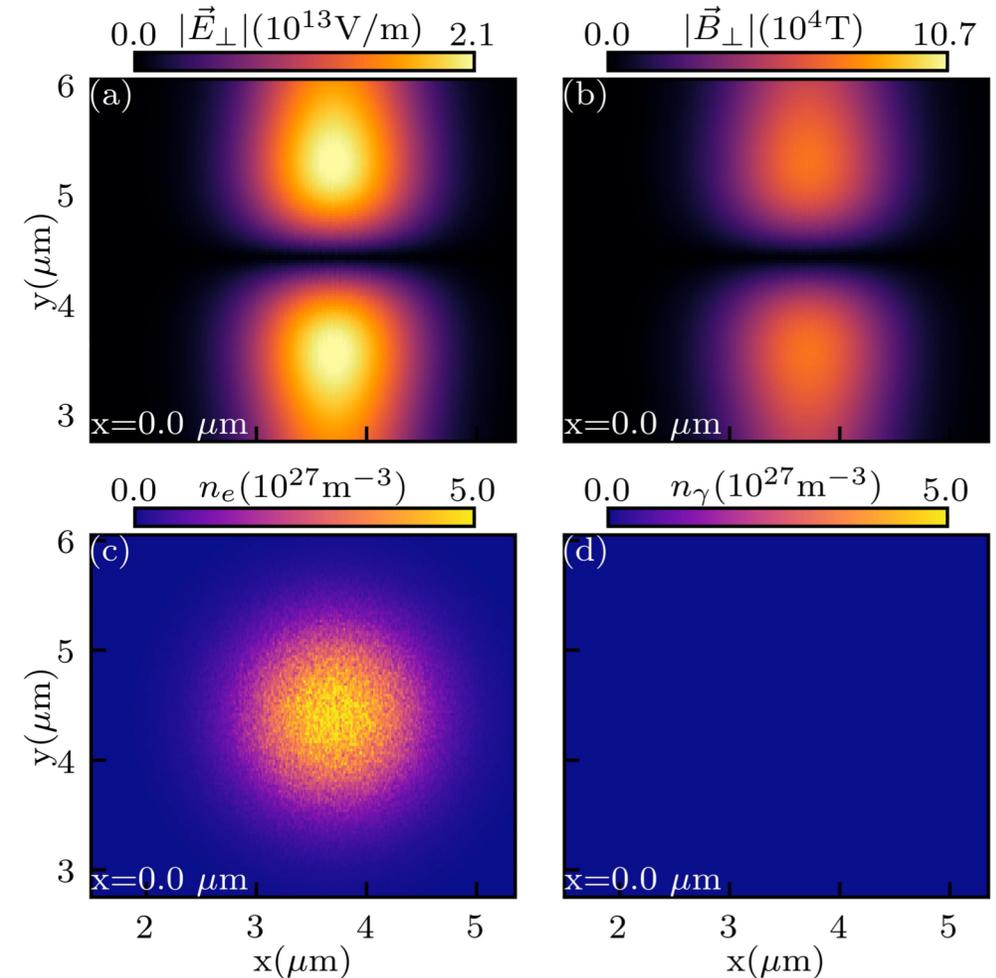
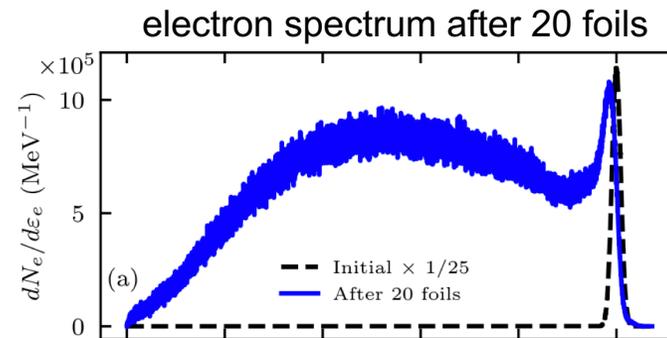
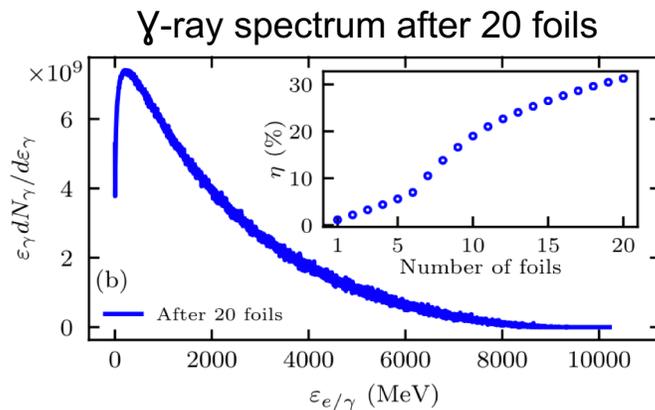


Conceptual representation of NFCTR image focusing<sup>2</sup>

<sup>1</sup> R. J. Adler, Part. Accel. 12, 39 (1982). <sup>2</sup> A. Matheron *et. al.*, submitted to Comm. Phys. <sup>3</sup> A. Sampath *et. al.*, PhysRevLett.126.064801 (2021)

# 1. Near-Field Coherent Transition Radiation – modelling and simulations

- 3D PIC simulation of e<sup>-</sup> beam interaction with stack of Al foils of 0.5 μm thickness.
  - Total charge  $Q = 2 \text{ nC}$ .
  - Initial beam energy  $E = 10 \text{ GeV}$ .
  - Initial beam size  $\sigma_{\parallel} = \sigma_{\perp} = 0.55 \mu\text{m}$ .
  - Initial beam normalized emittance  $\varepsilon_n = 3 \mu\text{m}$
- Conversion efficiency of  $\approx 30\%$  after 20 foils



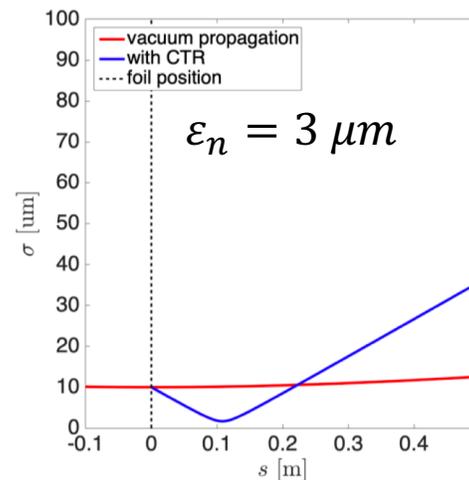
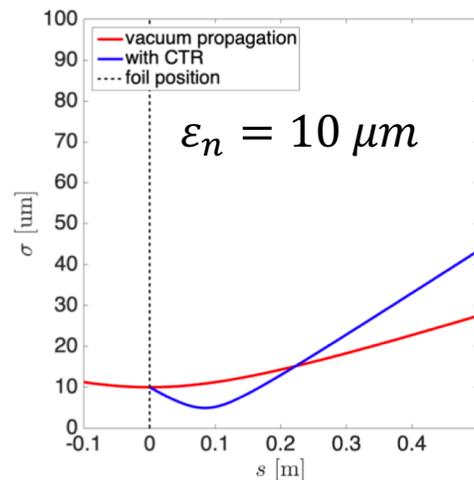
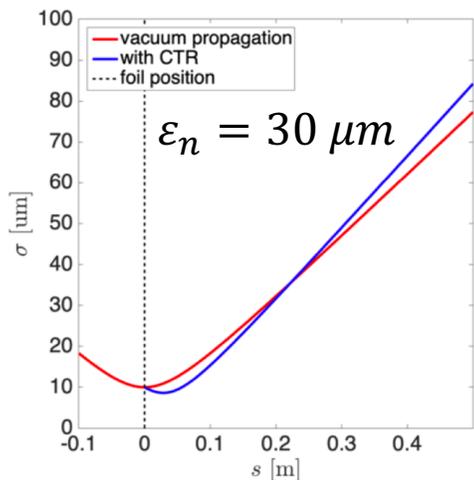
## 2. Analytical estimates and observables of NF-CTR beam focusing

- Motivation for E332 experimental collaboration at FACET-II.

- First goal: measure focusing effect on the beam.

Focusing effect depends on beam size  $\sigma_{\perp}$ , beam energy  $E$  and beam charge  $Q$

$$\text{Effective focal length}^4: f = \frac{8\pi\epsilon_0\sigma_{\perp}^2 E}{eQ}, \text{ to be compared to } \beta\text{-function } \beta/f = \frac{e^2 N}{8\pi\epsilon_0 m c^2 \epsilon_n} \propto \frac{Q}{\epsilon_n}$$



Normalized emittance $\epsilon_n$	$\beta/f$
100 $\mu\text{m}$	0.176
30 $\mu\text{m}$	0.586
10 $\mu\text{m}$	1.76
3 $\mu\text{m}$	5.86

$$Q = 2 \text{ nC}, E = 10 \text{ GeV}, \sigma_{\perp} = 10 \mu\text{m}$$

<sup>4</sup> S. Corde *et. al.*, Beam focusing by near-field transition radiation. [Research Report] IP Paris; CEA; MPIK. 2020. hal-02937777v2

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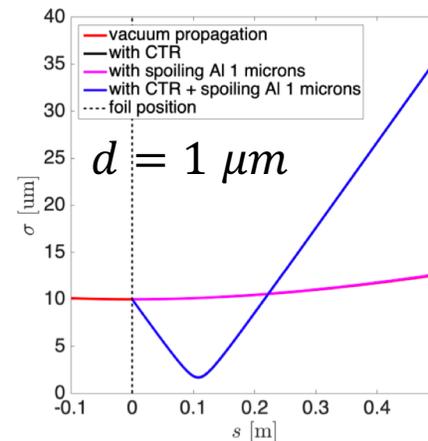
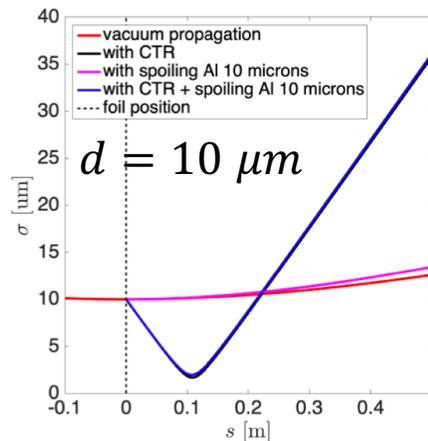
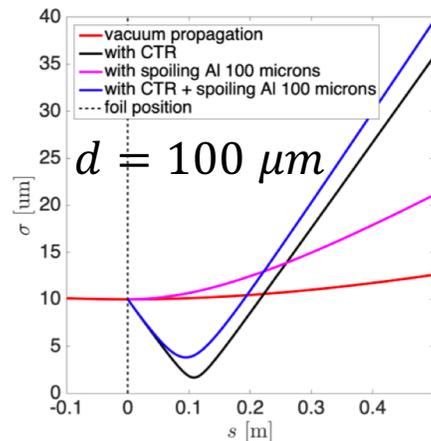
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Divergence increase due to multiple scattering in the foil: depends on foil thickness  $d$ .



$$\epsilon_n = 3 \mu\text{m}$$

$$Q = 2 \text{ nC}, E = 10 \text{ GeV}, \sigma_{\perp} = 10 \mu\text{m}$$

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- Optimized experimental parameters: **low beam emittance, high beam charge, thin foils.**

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Focusing effect depends on beam size  $\sigma_{\perp}$ , beam energy

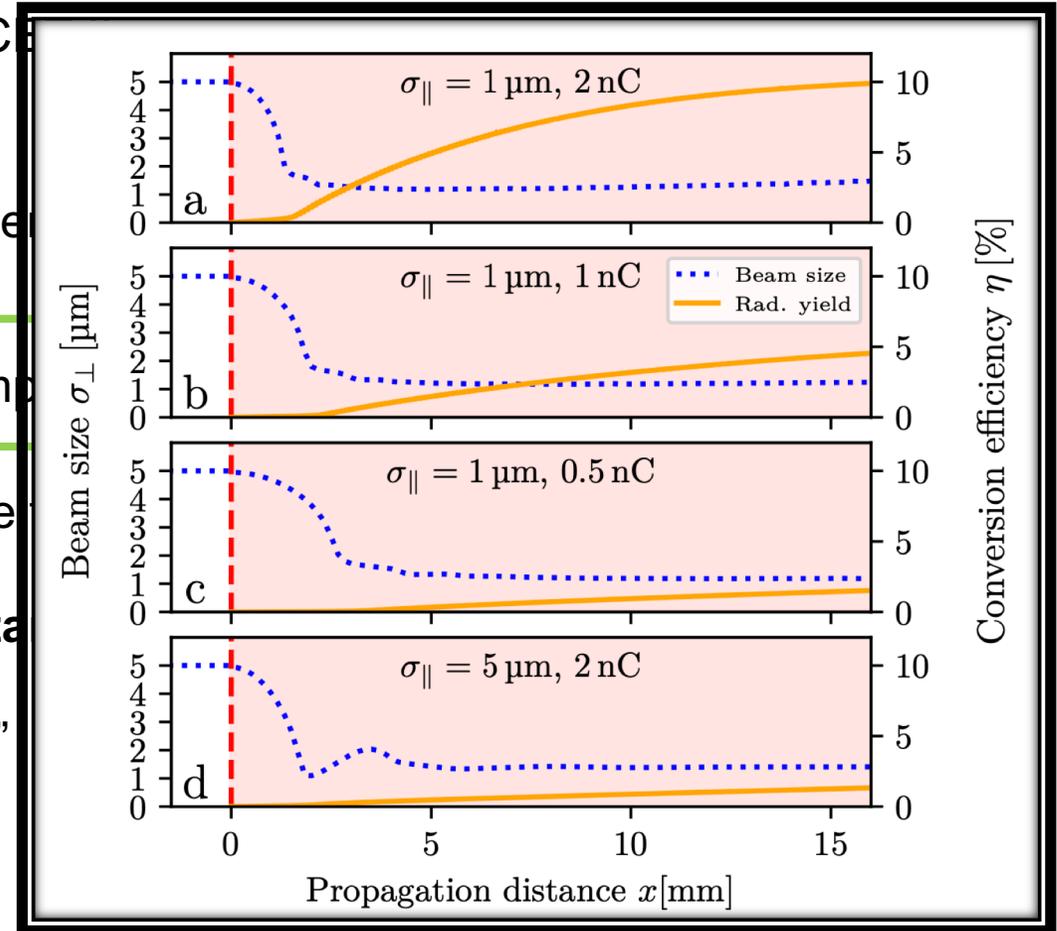
Effective focal length<sup>4</sup>:  $f = \frac{8\pi\epsilon_0\sigma_{\perp}^2 E}{eQ}$ , to be compared with

Divergence increase due to multiple scattering in the

- Optimized experimental parameters: **low beam emittance**
- Measure energy conversion to gammas with “relaxed” beam parameters?

$$\sigma_{\perp} = 5 \mu\text{m}, \epsilon_n = 3 \mu\text{m}, E = 10 \text{ GeV}$$

with stack of  $1 \mu\text{m}$  Al foils with  $100 \mu\text{m}$  spacing



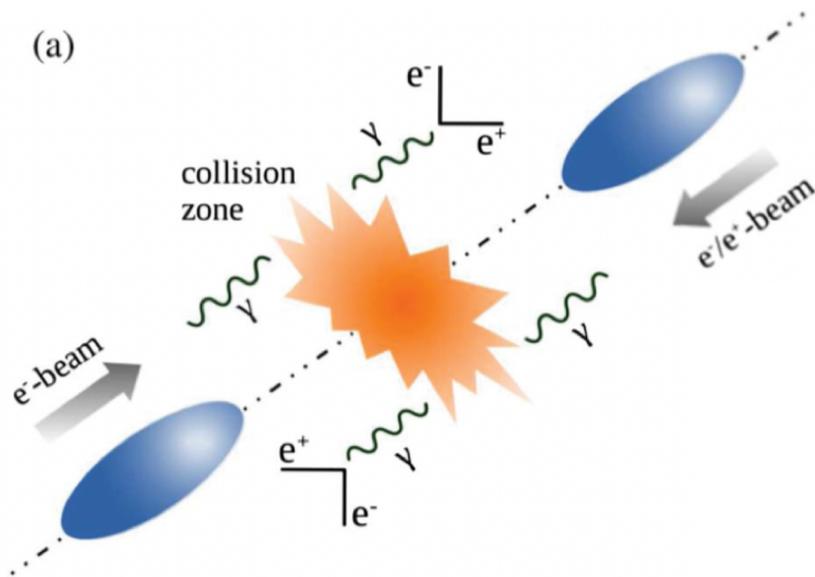
Rotationally symmetric PIC simulation results

<sup>4</sup> S. Corde *et. al.*, Beam focusing by near-field transition radiation. [Research Report] IP Paris; CEA; MPIK. 2020. hal-02937777v2

### 3. Probing Strong-Field regime of QED via NF-CTR process.

$E^*$  = Electric field in the  $e^-$  rest frame

- Strong-Field regime of QED:  $\chi = E^*/E_{cr} > 1 \rightarrow$  QED effects (pair production) become prominent.
- Beam-beam collisions were proposed<sup>5</sup> as a laser-less scenario to study the nonperturbative regime ( $\alpha\chi^{2/3} > 1$ )



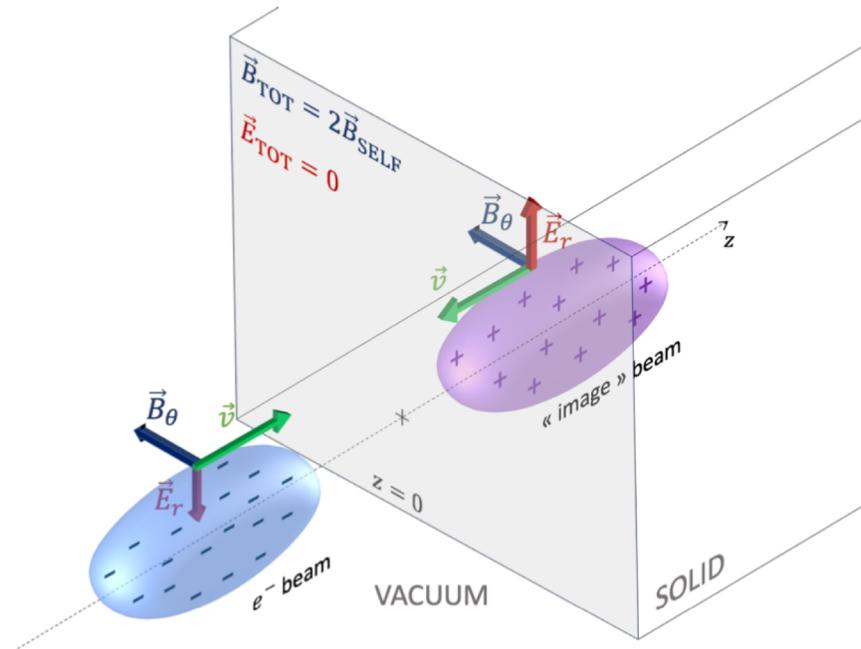
- The beam-beam scenario benefits from:
  - The large Lorentz factor  $\gamma$  of the relativistic  $e^-$  beams.
  - The high amplitude electromagnetic fields generated by tightly focused bunches.

<sup>5</sup> V. Yakimenko *et. al.*, Phys. Rev. Let. 122, 190404 (2019)

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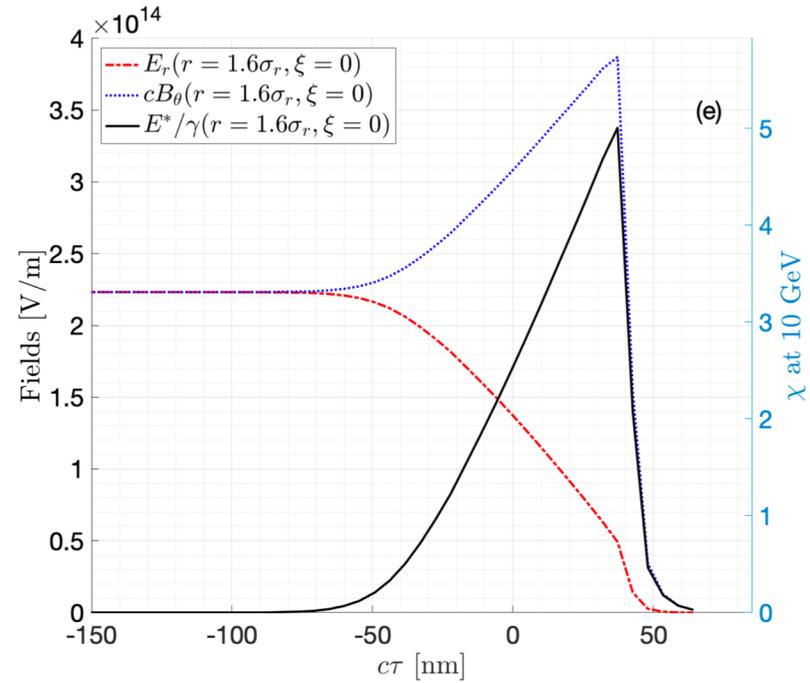
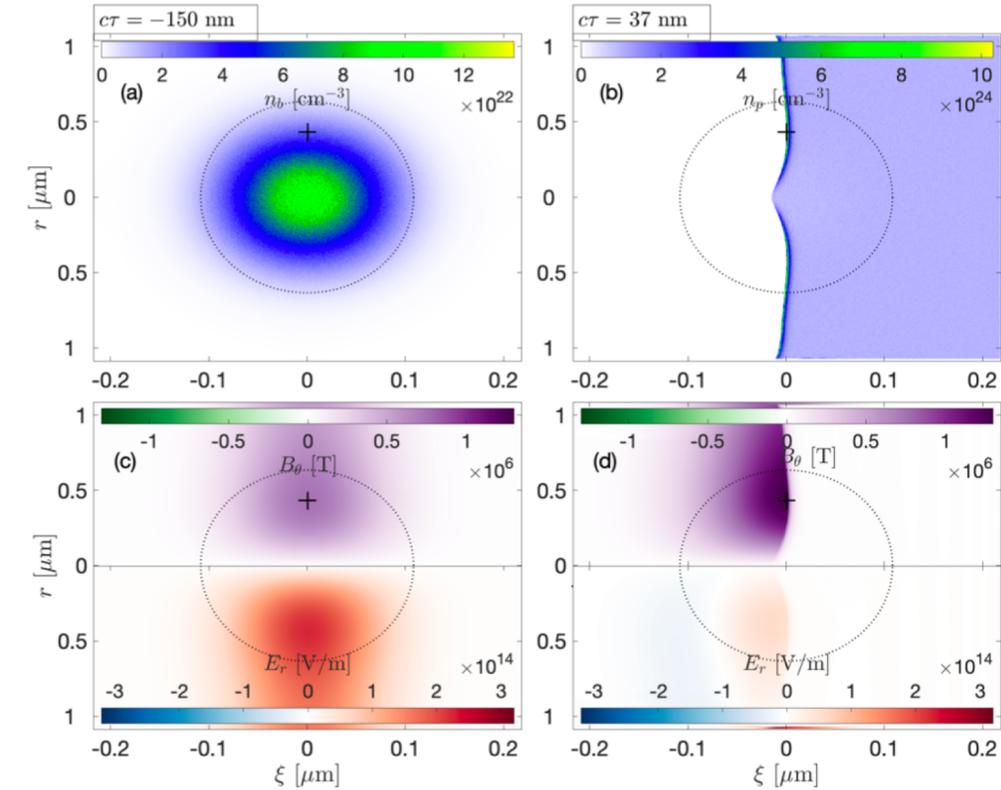
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- The beam-beam scenario benefits from:
  - The large Lorentz factor  $\gamma$  of the relativistic  $e^-$  beams.
  - The high amplitude electromagnetic fields generated by tightly focused bunches.
- **Idea:** Via the NFCTR process, one  $e^-$  beam can be replaced by a conductor.
  - The beam experiences its own self-fields.
  - Offers a significant simplification of the experiment.

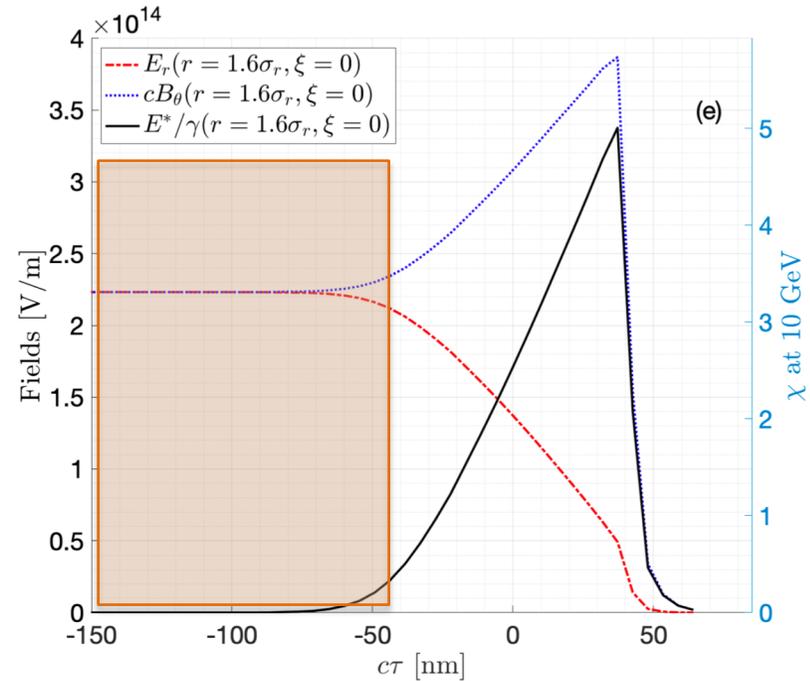
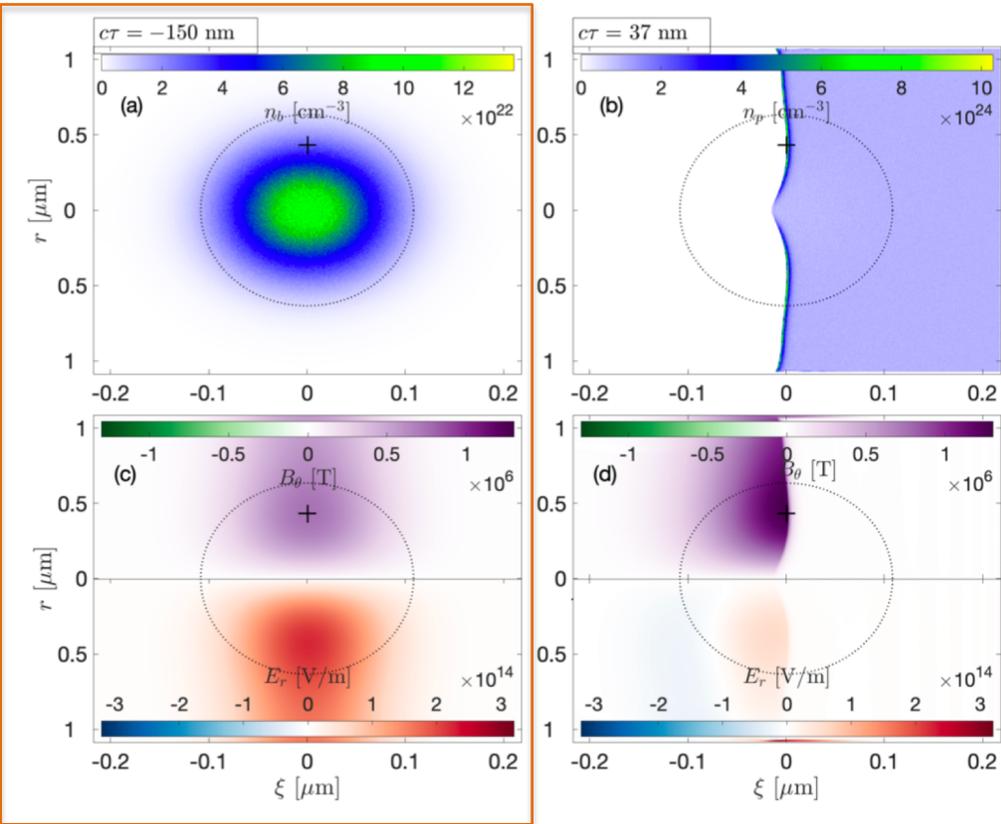
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**PIC simulation parameters:**  
 Beam charge  $Q = 2$  nC.  
 Beam energy  $E = 10$  GeV.  
 Bunch length  $\sigma_{\parallel} = 54$  nm.  
 Beam size  $\sigma_{\perp} = 271$  nm.  
 Plasma density  $n_p = 10^{24}$  cm $^{-3}$ .

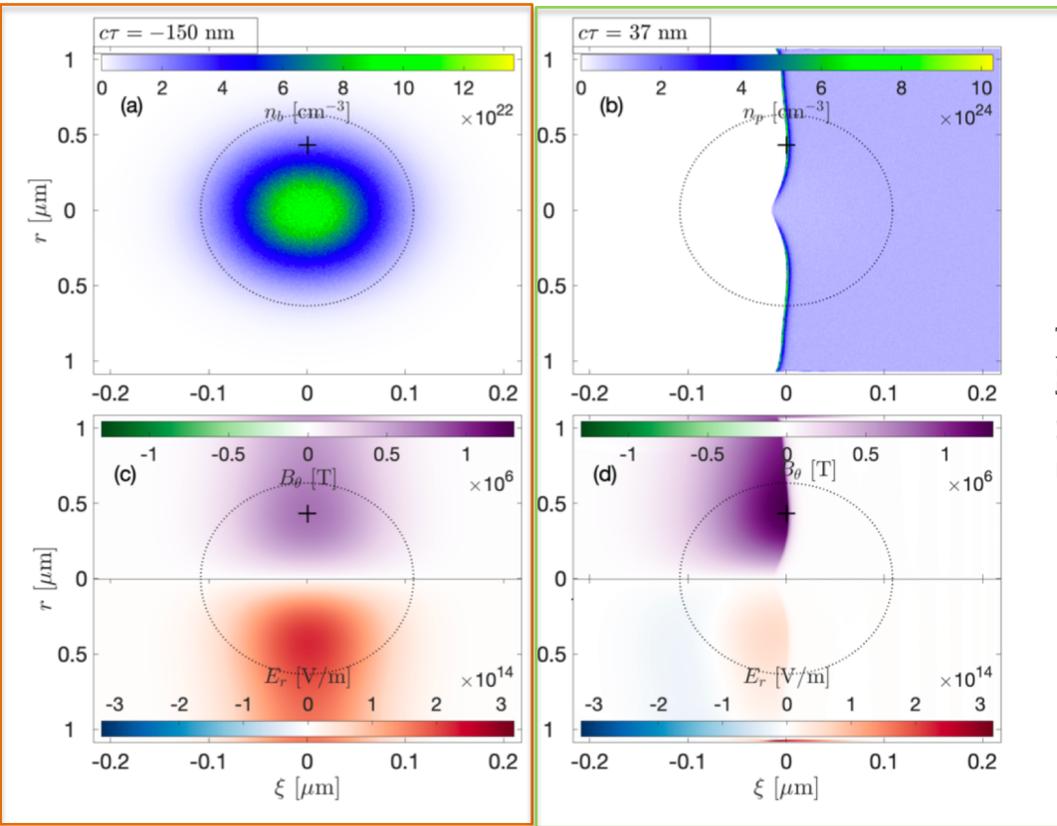
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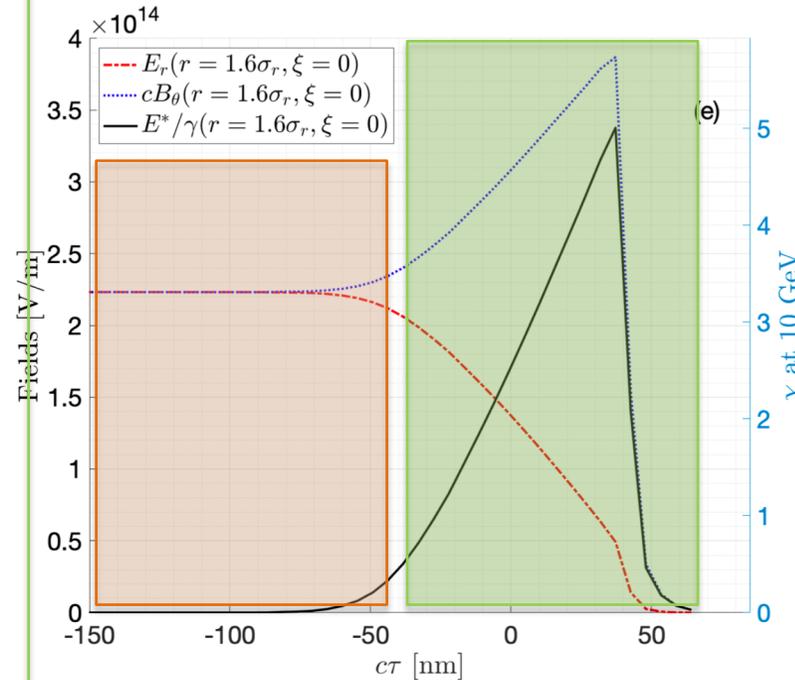
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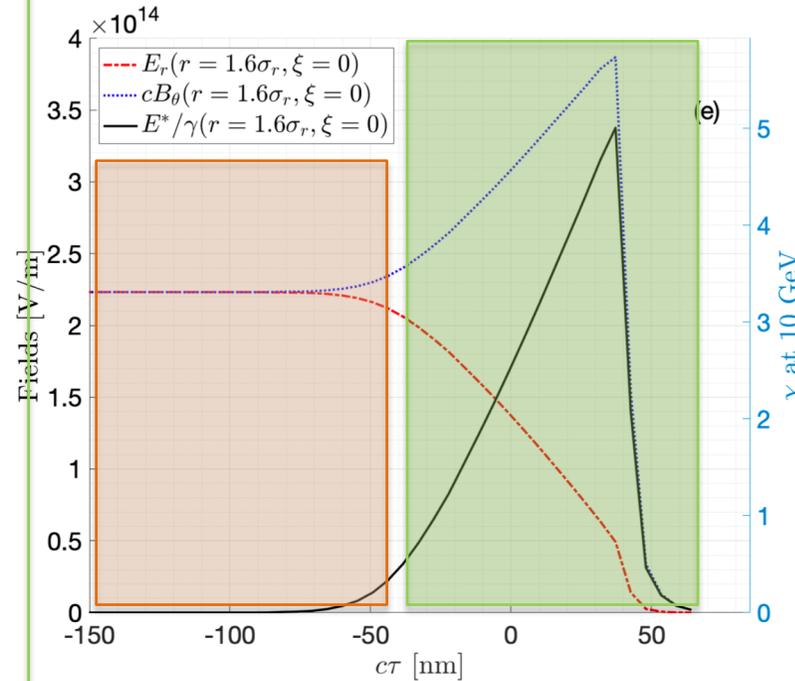
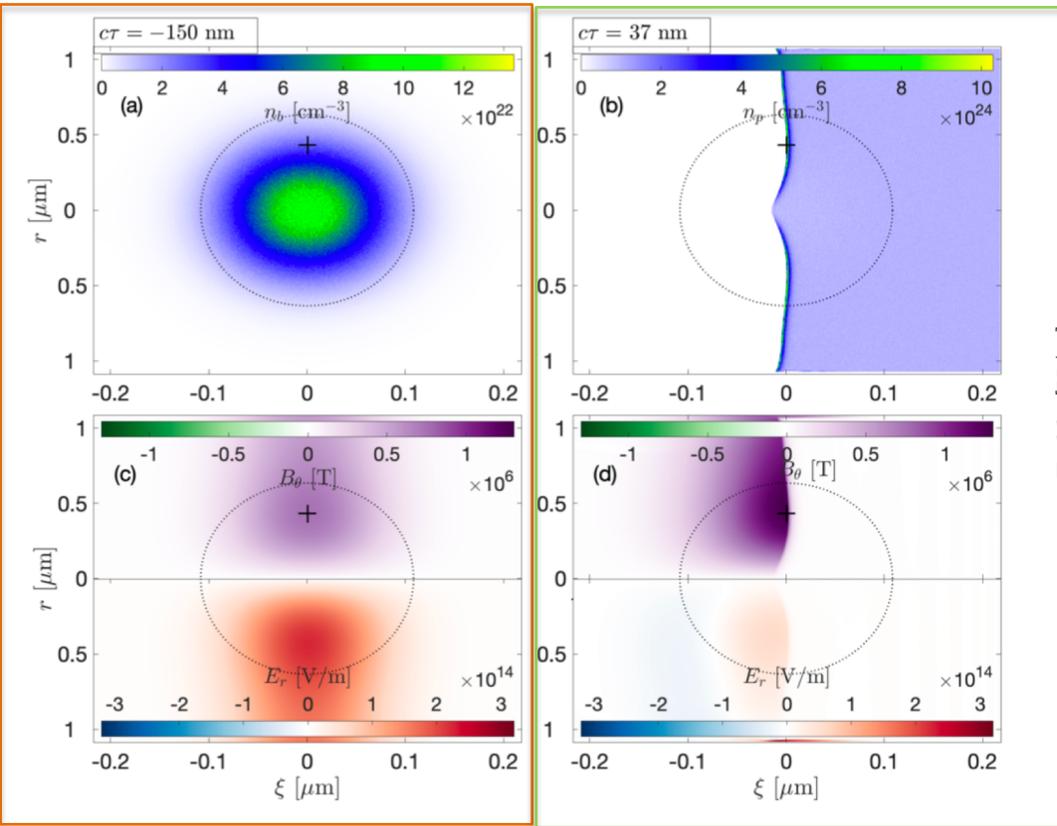


**PIC simulation parameters:**  
 Beam charge  $Q = 2 \text{ nC}$ .  
 Beam energy  $E = 10 \text{ GeV}$ .  
 Bunch length  $\sigma_{\parallel} = 54 \text{ nm}$ .  
 Beam size  $\sigma_{\perp} = 271 \text{ nm}$ .  
 Plasma density  $n_p = 10^{24} \text{ cm}^{-3}$ .

At the plasma surface (NFCTR)

- Radial electric component vanish ( $E_r \rightarrow 0$ )
- Azimuthal magnetic component
  - Doubles (radiative regime,  $\sigma_{\parallel} \ll \sigma_{\perp}$ )
  - Remains constant (stationary regime,  $\sigma_{\parallel} \gg \sigma_{\perp}$ )

### 3. Probing Strong-Field regime of QED via NF-CTR process.



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Control beam parameters  
 ↓  
 Precision studies of SFQED

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### 3. Probing Strong-Field regime of QED via NF-CTR process.

Can we reach an arbitrarily high  $\chi$  ?

- Increase beam energy ( $\chi \propto \gamma$ )
- Increase beam self-fields amplitude  $\rightarrow$  Increase beam density  $n_b$**

However, NFCTR fails if  $n_b \gtrsim n_p$  (plasma is not able screen the beam self-fields).

**The maximum achievable plasma density  $n_p$  in an experiment limits the highest value of  $\chi$**

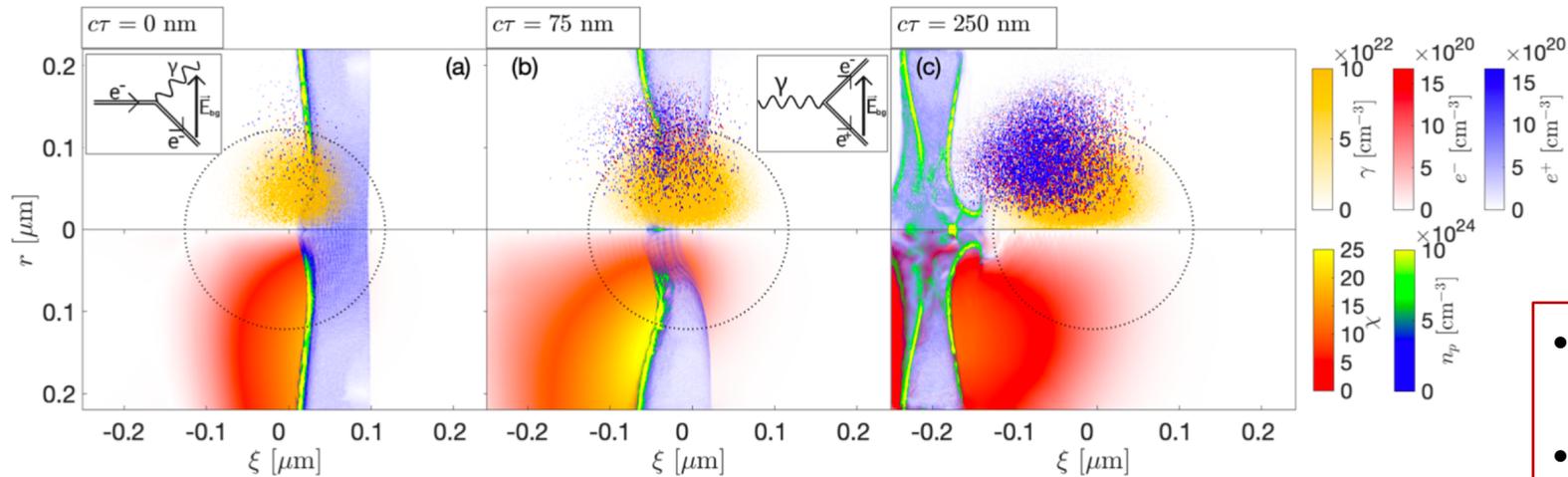
- Due to large field ionization of solid atoms and plasma surface compression<sup>6</sup>, PIC simulations show that for  $n_b \approx n_p$  the beam self-fields are reflected leading to  $\chi \approx \gamma E_{self} / E_{cr}$ .

Furthermore, the beam shape  $\sigma_{\parallel} / \sigma_{\perp}$  also affects the reflection (radiative vs. stationary regimes), and thus also limits the maximum  $\chi$  of a given experimental configuration.

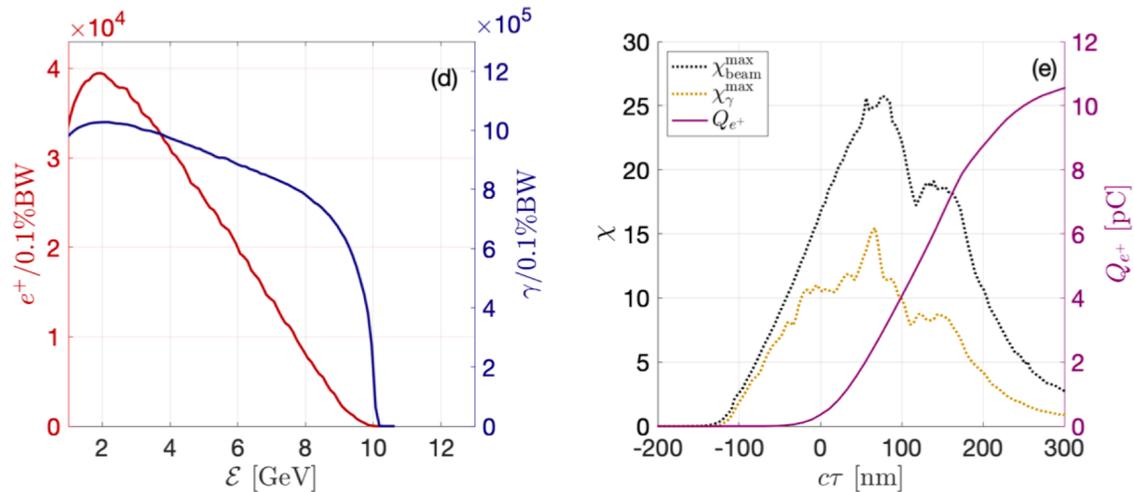
<sup>2</sup>A. Matheron *et. al.*, submitted to Comm. Phys. <sup>6</sup>X. Xu *et. al.* Phys. Rev. Lett. 126, 094801 (2021).

### 3. Probing Strong-Field regime of QED via NF-CTR process.

Taking into account all the limiting phenomena, a QED-PIC simulation<sup>7</sup> with an optimized set of parameters results in:



Beam charge  $Q = 2$  nC.  
 Beam energy  $E = 10$  GeV.  
 Beam size  $\sigma_{\parallel} = \sigma_{\perp} = 55$  nm.  
 Initial Au<sup>1+</sup> foil



- $\chi_{max} \approx 25$
- Abundant high energy gamma-ray and Breit-Wheeler pair production, well above competing processes such as Bremsstrahlung or Bethe-Heitler.
- 10 pC of positron charge above 1 MeV

<sup>2</sup> A. Matheron *et. al.*, submitted to Comm. Phys.

<sup>7</sup> M Lobet et al 2016 J. Phys.: Conf. Ser. 688 012058

# Summary and conclusions

The intense EM fields produced at the surface of a solid conducting foil when interacting a high peak-current beam (**Near-Field CTR**) are interesting as:

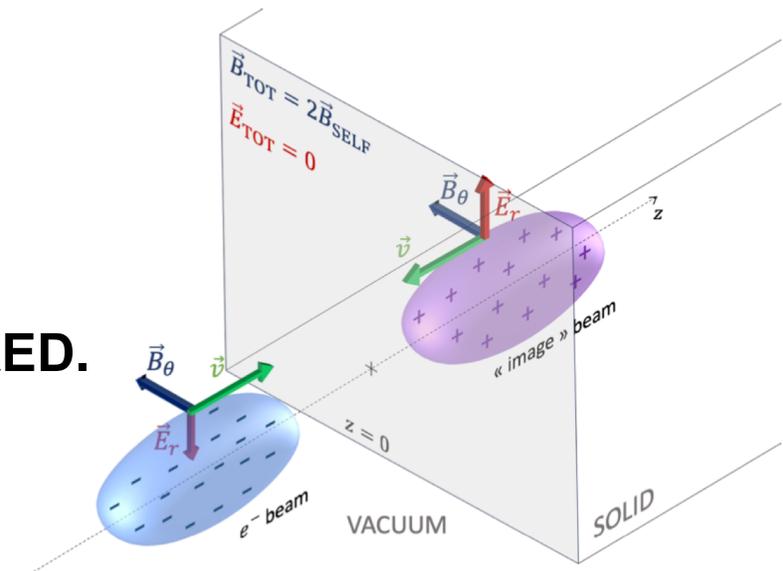
- **Beam focusing elements to reach solid-density beams → bright gamma-ray source.**

Main motivation of E332 experiment at FACET-II:

- need high charge, low emittance and thin foils.
- Simulations show that substantial conversion efficiency should be measured with FACET-II nominal parameters

- **Laser-less scenario to probe the Strong-Field regime of QED.**

- Precision studies of SFQED.
- QED signal above competing processes.
- Intermediate step before beam-beam collisions.



# Thanks: E332 Collaboration

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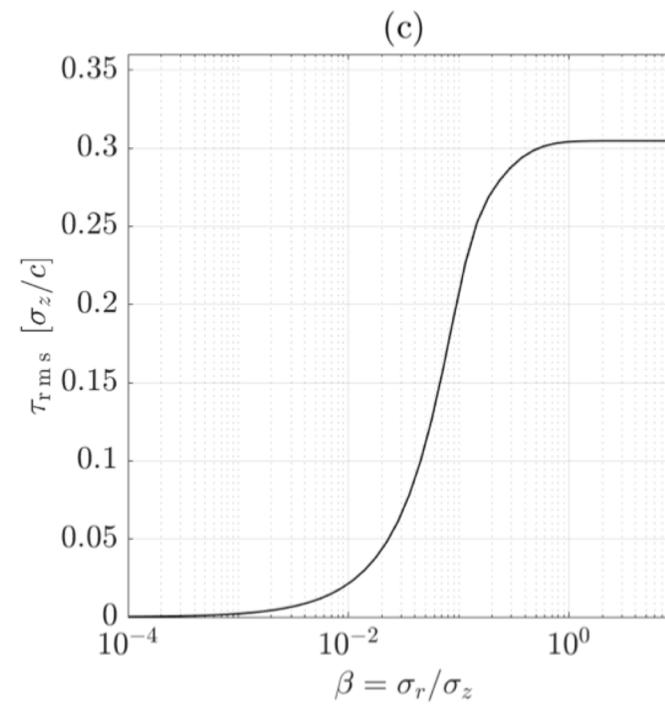
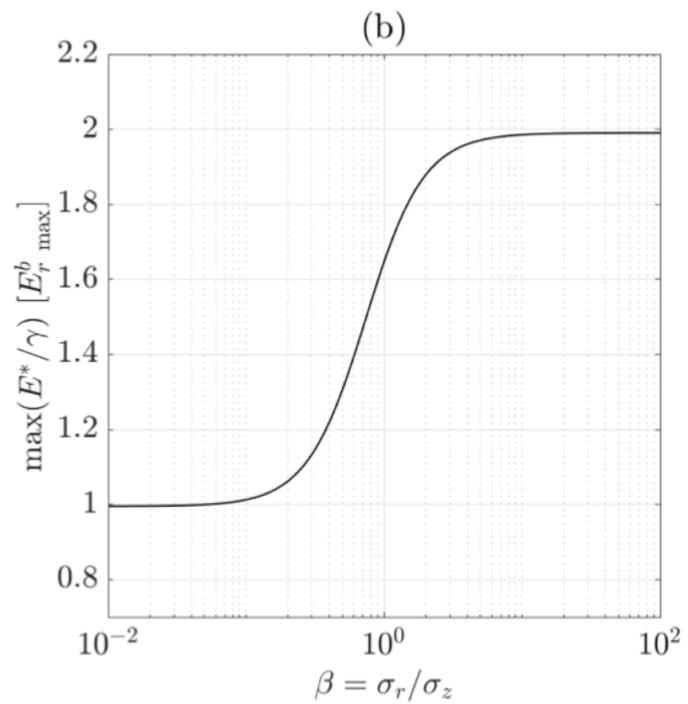
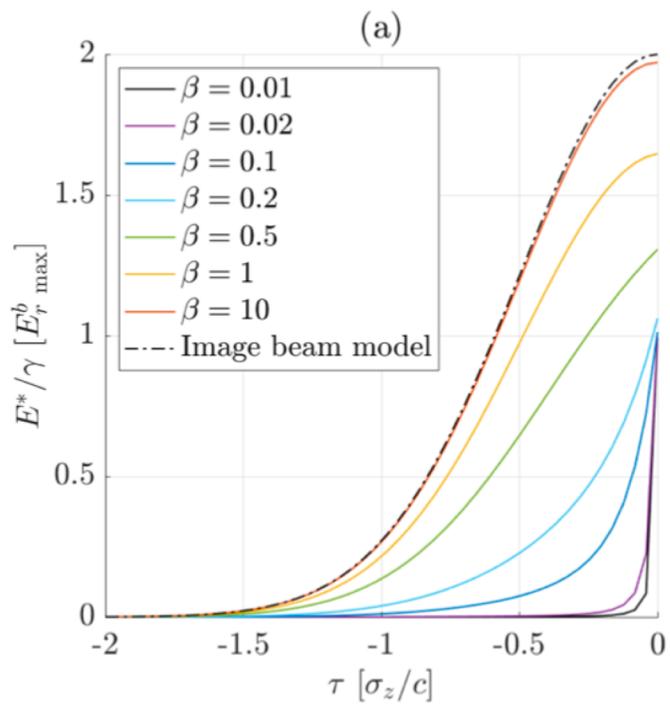




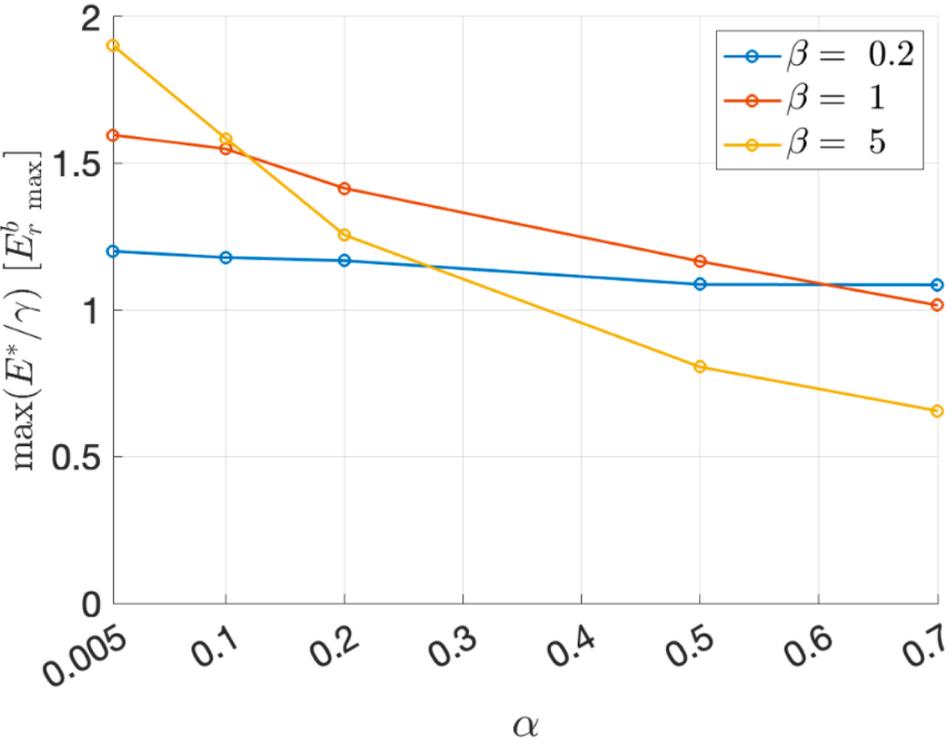
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