

# Propagating Higher Order Mode Simulations in SRF Multi-Cavity Cryomodules Using ACE3P

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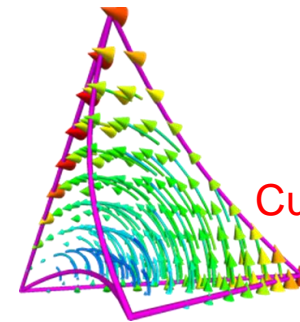
Computational Electrodynamics Department  
*SLAC National Accelerator Laboratory*

ICFA Mini Workshop on Higher Order Modes in Superconducting Cavities  
(HOMSC18) 1-3, October 2018, Cornell University

- **ACE3P and Omega3P**
- **Nonlinear Eigensolver CORK in Omega3P**
- **Benchmark of CORK Eigensolver**
- **3<sup>rd</sup> dipole band trapped mode simulations in TDR CM using Omega3P**
- **Summary**

**ACE3P** is a comprehensive set of parallel multi-physics codes

- Based on *high-order curved finite elements* for high-fidelity modeling
- Implemented on *massively parallel computers* for increased problem size and speed



$$\mathbf{E} = \sum_i x_i \mathbf{N}_i$$

Curved finite element



**NERSC Cori:** Cray XC40

- 632,672 compute cores
- 1 PB of memory
- peak performance of 27.9 Pflops/sec

## **ACE3P (Advanced Computational Electromagnetics 3P)**

Frequency Domain:

**Omega3P**

– Eigensolver (damping)

**S3P**

– S-Parameter

Time Domain:

**T3P**

– Wakefields and Transients

Particle Tracking:

**Track3P**

– Multipacting and Dark Current

EM Particle-in-cell:

**Pic3P**

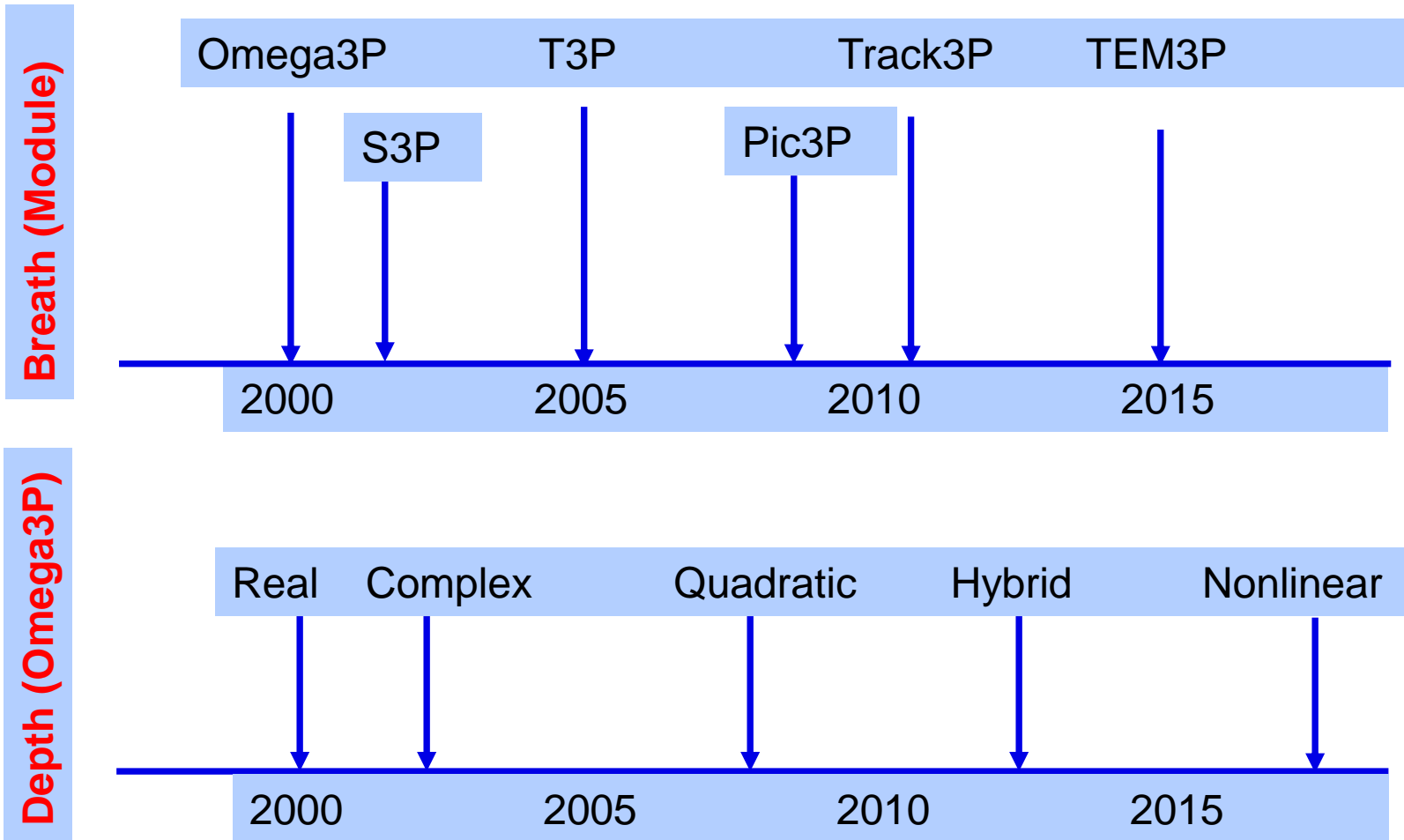
– RF guns & space charge effects

Multi-physics:

**TEM3P**

– EM, Thermal & Mechanical analysis

# ACE3P Development



# RF Eigenmode Analysis in Omega3P

## Maxwell's Equations in Frequency Domain in the E-formulation

$$\nabla \times \left( \frac{1}{\mu} \nabla \times \vec{\mathbf{E}} \right) - k^2 \epsilon \vec{\mathbf{E}} = 0 \text{ on } \Omega$$

$$\vec{\mathbf{n}} \times \vec{\mathbf{E}} = 0 \text{ on } \Gamma_E$$

$$\vec{\mathbf{n}} \times \frac{1}{\mu} \nabla \times \vec{\mathbf{E}} = 0 \text{ on } \Gamma_M$$

$$\vec{\mathbf{n}} \times \frac{1}{\mu} \nabla \times \vec{\mathbf{E}} + i \sqrt{k^2 - k_{c1}^2} \vec{\mathbf{n}} \times \vec{\mathbf{n}} \times \vec{\mathbf{E}} = 0 \text{ on } \Gamma_{wg}$$

PEC

PMC

WG (eg. TE mode)

## Discretizing electric field using high-order Nedelec type basis functions $\mathbf{N}_i$

$$\mathbf{E} = \sum_i x_i \mathbf{N}_i$$

## Eigenvalue problem

$$F(\lambda)\mathbf{x} = 0, \quad \lambda (= k^2) \text{ the eigenvalue, } \mathbf{x} \text{ the eigenvector}$$

$$k_{\text{real}} = \frac{2\pi f}{c}, \quad Q = \frac{K_{\text{real}}}{2K \text{Im}_{ag}}$$

# Omega3P Capabilities

Lossless Cavity

Cavity w/ Lossy Materials

Cavity w/ SIBC

Cavity w/ WG ports (one cutoff)

Cavity w/ WG ports (multi-modes)

$$F(\lambda)x = 0,$$



$$Kx = k^2Mx$$

Real/Complex eigenvalue problem



$$Kx + ikWx = k^2Mx$$

$$Kx + i\sqrt{k^2 - k_c^2}Wx = k^2Mx$$

Quadratic eigenvalue problem



Nonlinear eigenvalue problem

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# Nonlinear RF Eigenvalue Problem

- When waveguide ports have different cutoff frequency modes, the RF system is a nonlinear eigensolver problem;

$$Kx + i\sqrt{k^2 - (k^c)^2} W^{TEM} x + i\sum_m \sqrt{k^2 - (k_m^c)^2} W_m^{TE} x + i\sum_m \frac{k^2}{\sqrt{k^2 - (k_m^c)^2}} W_m^{TM} x = k^2 Mx$$

$$K = \int_{\Omega} (\nabla \times N_i) \cdot \frac{1}{\mu} (\nabla \times N_j) d\Omega$$

$$M = \int_{\Omega} N_i \cdot \epsilon N_j d\Omega$$

$$W^{TEM} = \int e_0^{TEM} \cdot \vec{N}_i ds \int e_0^{TEM} \cdot \vec{N}_j ds$$

$$W^{TE} = \int e_0^{TE} \cdot \vec{N}_i ds \int e_0^{TE} \cdot \vec{N}_j ds$$

$$W^{TM} = \int e_0^{TM} \cdot \vec{N}_i ds \int e_0^{TM} \cdot \vec{N}_j ds$$

M: the mass matrix, K: the stiffness matrix,  $e_0$ : the normalized waveguide mode



# Compact Rational Krylov for Nonlinear Eigenvalue Problem

- The development of a *nonlinear eigensolver CORK* with LBNL has been integrated in Omega3P;

Three steps for solving NLEPs using CORK:

1. Approximation of the scalar nonlinear functions by interpolating rational functions, yielding a rational eigenvalue problem;

$$F(\lambda) = \mathbf{0} \rightarrow \max \|F(\lambda) - R(\lambda)\| < \varepsilon_{approx}$$

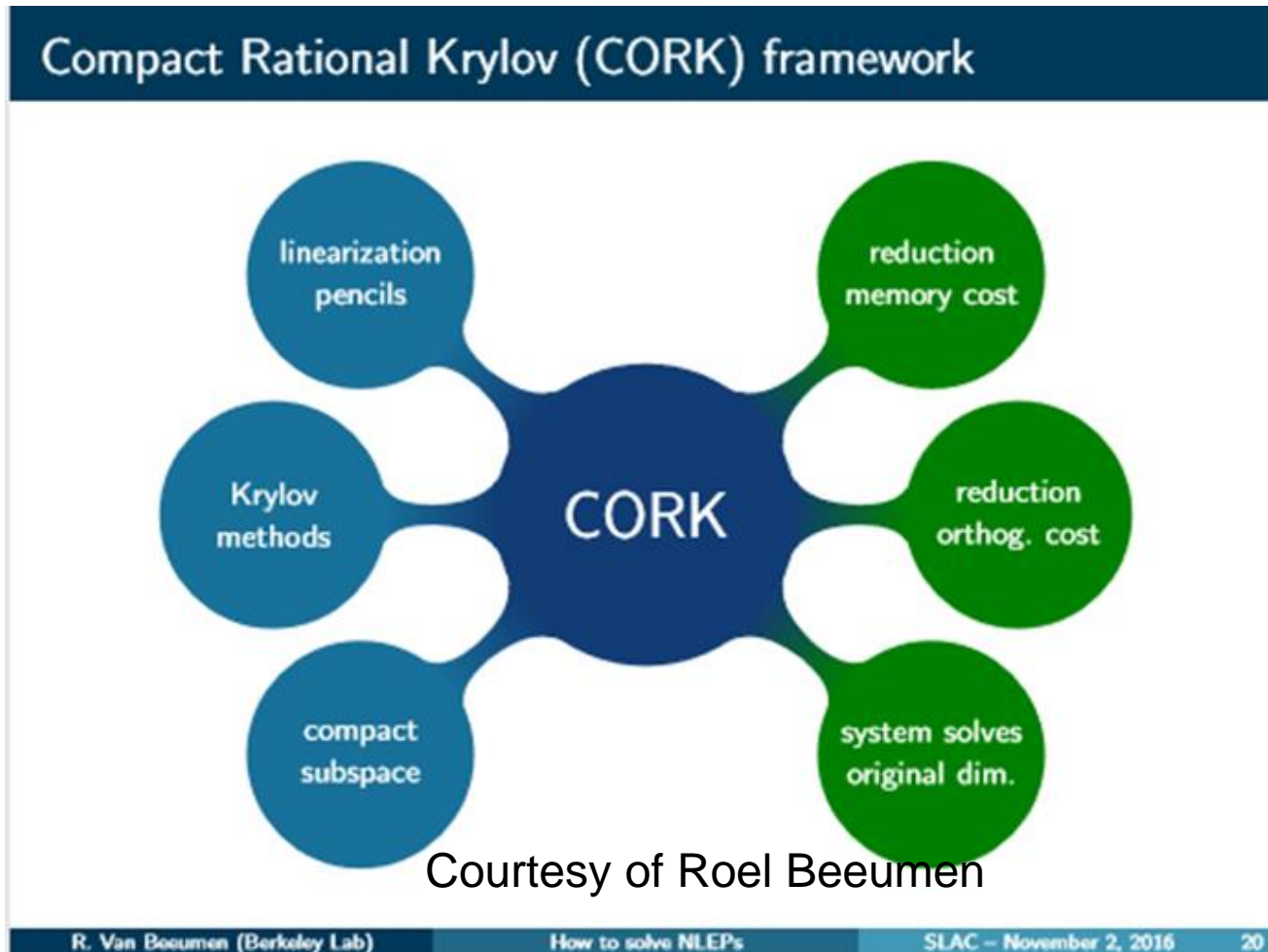
2. Linearization of the resulting rational eigenvalue problem, i.e., a reformulation of the rational eigenvalue problem as a generalized (linear) eigenvalue problem with the same eigenvalues but much larger in problem size;

$$R_d(\lambda) = A_0 + A_1\lambda + A_2\lambda^2 + \dots + A_d\lambda^d$$

3. Solving the generalized eigenvalue problem by a CORK method.

$$L(\lambda)x = (A - \lambda B)x = \mathbf{0}$$

# Compact Rational Krylov Framework



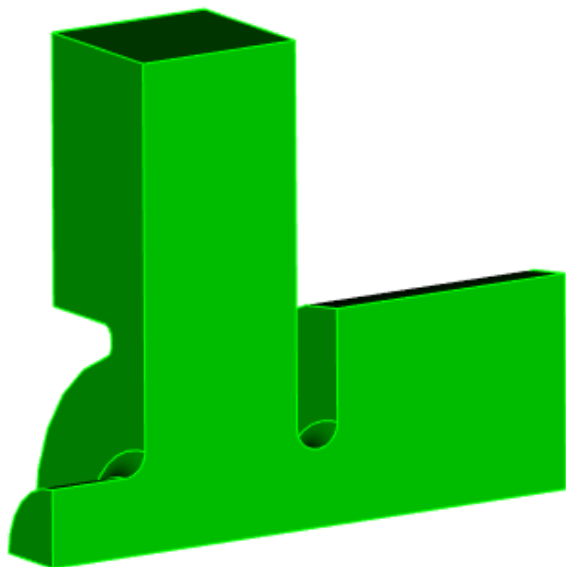
Roel Van Beeumen, et al, "Computing Resonant Modes of Accelerator Cavities by Solving Nonlinear Eigenvalue Problems via Rational Approximation", *Journal of Computational Physics*, Aug. 2018

- ACE3P and Omega3P
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# Benchmark of CORK Eigensolver

- Using a pillbox with circular & rectangular waveguides for the solver benchmark;
- It is a nonlinear eigenvalue problem, which can be solved using S3P or Omega3P.

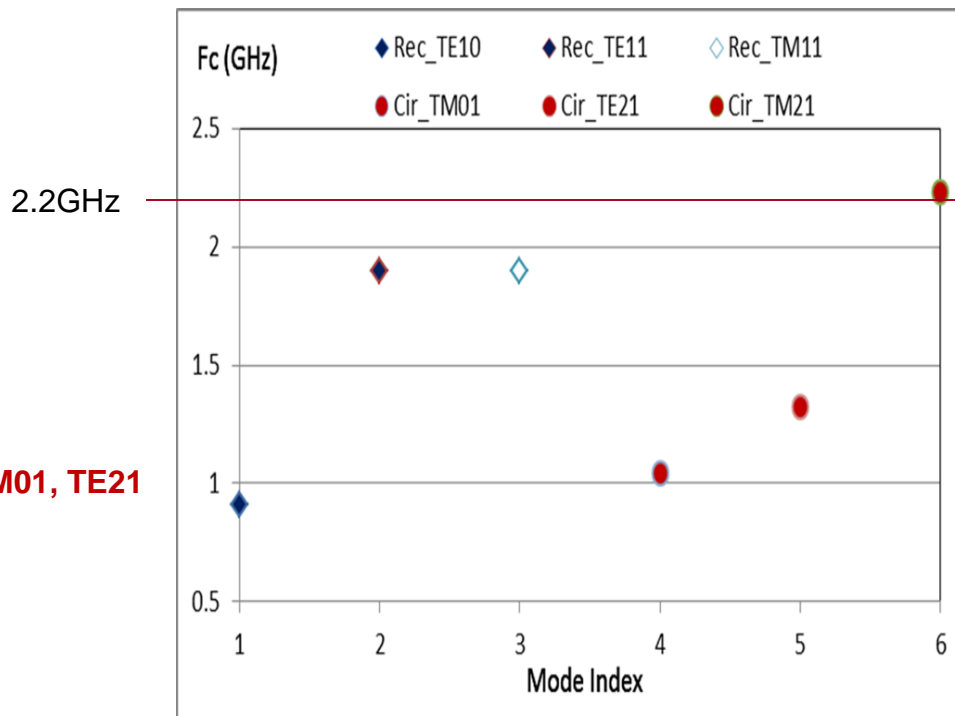
TE10, TE11, TM11



Two sym. Planes: PMC  
Beampipe End: PEC ( $f < f_c$ )

$f_c = 2.2\text{GHz}$

TM01, TE21



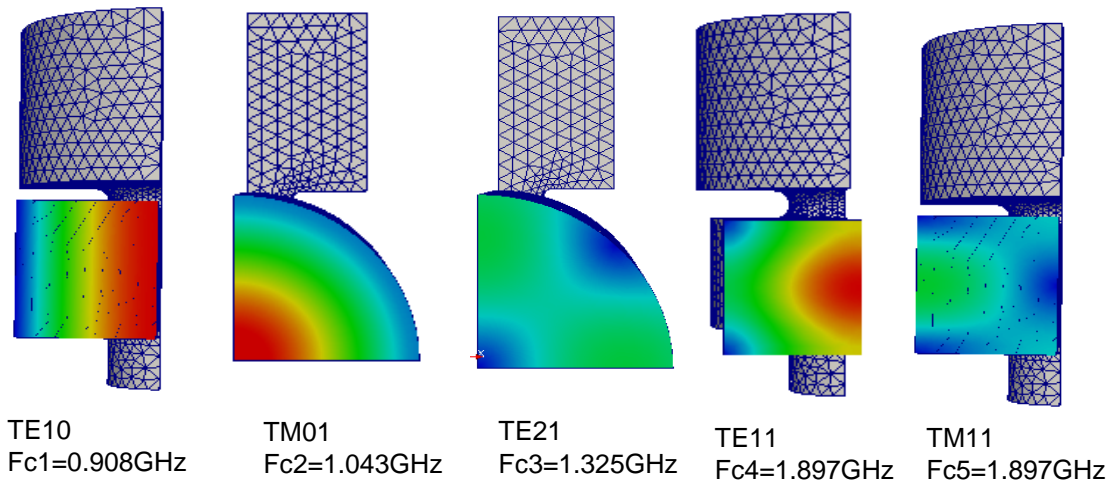
The lowest 6 eigenmodes  $f_c$  in the couplers

# Nonlinear RF Eigenvalue Problem

- The eigenmodes in the RF cavity under 2.2GHz can couple to three rectangular & two circular waveguide modes;
- This configuration yields the nonlinear eigenvalue problem  $F(\lambda)x=0$ , where

$$F(\lambda) = K - \lambda M + i\sqrt{\lambda - k_1^2} W_{10}^{TE} + i\lambda/\sqrt{\lambda - k_2^2} W_{01}^{TM} + i\sqrt{\lambda - k_3^2} W_{21}^{TE} + i\sqrt{\lambda - k_4^2} W_{11}^{TE} + i\lambda/\sqrt{\lambda - k_5^2} W_{11}^{TM}$$

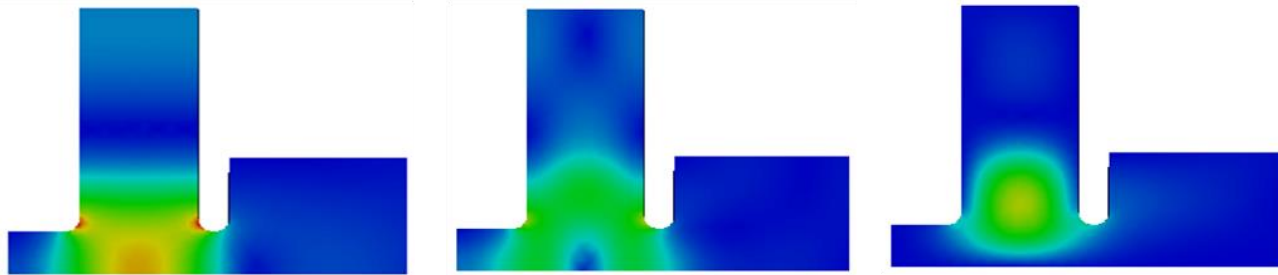
$$k_i = 2\pi f_{ci}/c$$



2D port modes from Omega3P

# Eigenmode Solutions from Omega3P

- CORK solver integrated into Omega3P allows correct calculation of mode damping with different cutoff frequencies at waveguide ports.
- There are three eigenmodes below beampipe cutoff (2.2GHz) from Omega3P.



Eigenmodes from Omega3P

## Error parameter

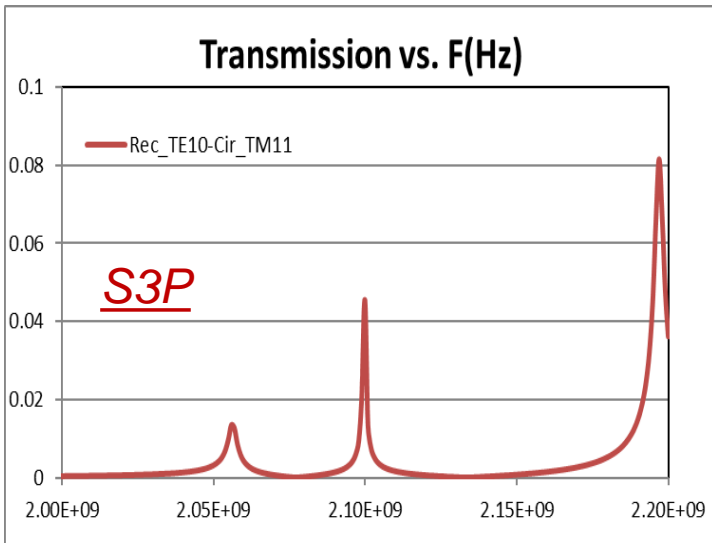
$$E(\lambda, x) = \frac{\|F(\lambda)x\|_2 / \|x\|_2}{\alpha(\lambda)}$$

$$\alpha(\lambda) = \|K\|_2 + |\lambda| \|M\|_1 + \sqrt{|\lambda - k_1^2|} \|W_{10}^{TE}\|_1 + \frac{|\lambda|}{\sqrt{|\lambda - k_2^2|}} \|W_{01}^{TM}\|_1 + \sqrt{|\lambda - k_3^2|} \|W_{21}^{TE}\|_1 + \sqrt{|\lambda - k_4^2|} \|W_{11}^{TE}\|_1 + \frac{|\lambda|}{\sqrt{|\lambda - k_5^2|}} \|W_{11}^{TM}\|_1$$

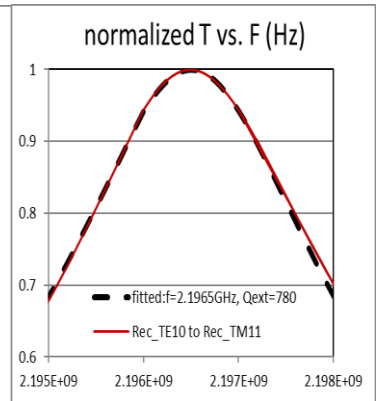
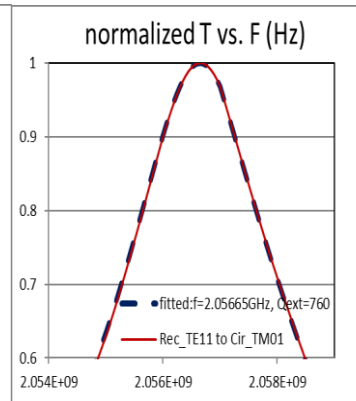
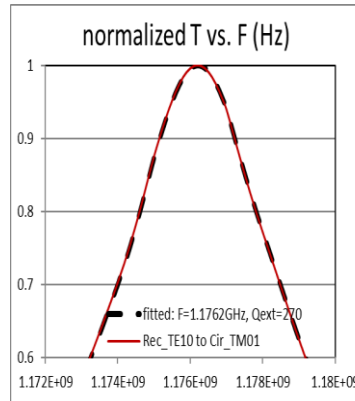
Omega3p-CORK			
Mode	F (GHz)	Qext	E(λ, x)
1	1.1762	271	1.2827e-14
2	2.0567	739	2.1098e-14
3	2.1960	693	4.5892e-15

# Eigenmode Solutions from S3P

- Searching for the eigenmodes below beampipe cutoff (2.2GHz) using S3P;
- Getting eigenmode  $f_0$  and  $Q_{ext}$  through fitting the transmission coefficients.



$$T^2 = \frac{1}{1 + Q_{ext}^2 \left( \frac{f}{f_0} - \frac{f_0}{f} \right)^2}$$



- The cavity eigenmode results from Omega3P and S3P agree well;
- Omega3P provides a direct way to compute the eigenmode parameters;
- Omega3P can solve the eigenmodes with strong damping accurately, which S3P cannot.

	Omega3P		S3P	
Mode	F (GHz)	Qext	F (GHz)	Qext
1	1.1762	271	1.1762	270
2	2.0567	739	2.0566	760
3	2.1960	693	2.1965	780

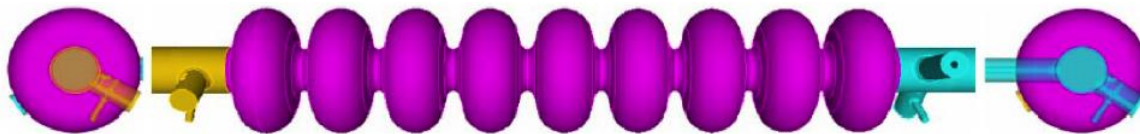
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# 8-Cavity TESLA TDR Cryomodule

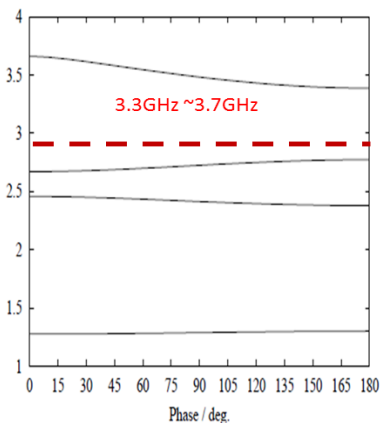
- TESLA cavity has been adopted by XFEL and LCLS-II as well as proposed in ILC;
- The higher-order modes need to be damped properly for reliable operation of the machines;
- The lower monopole and dipole band modes below the beampipe cutoff have been well studied in a single TESLA cavity;
- The HOMs above the beampipe cutoff need to be studied in a whole CM.

TESLA Cavity

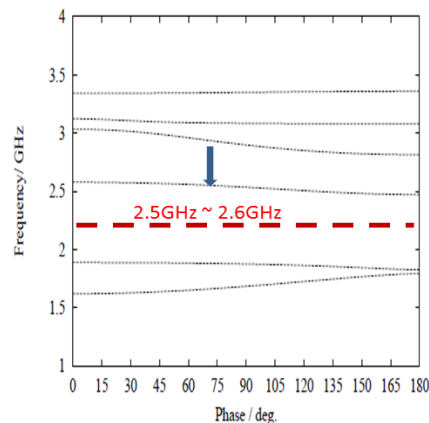


R\_beampipe=39mm, Fc(dipole)=2.2GHz, Fc(mono)=2.9GHz

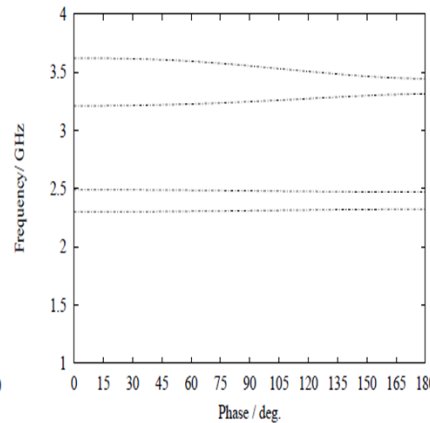
monopole



dipole



Quad.



HOMs under 4GHz

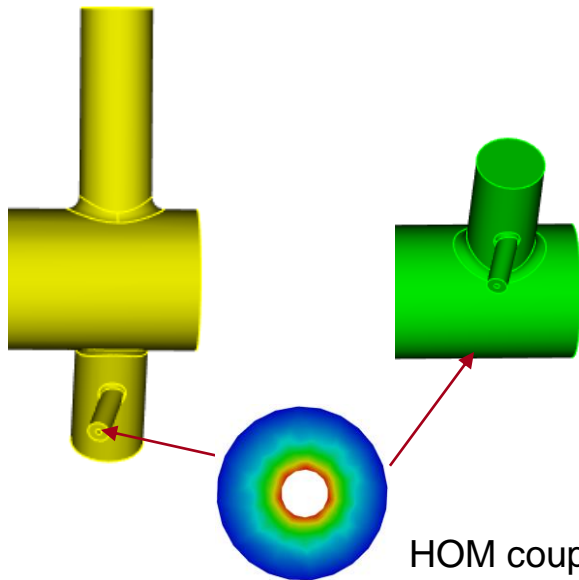


8-Cavity CM for LCLS-II

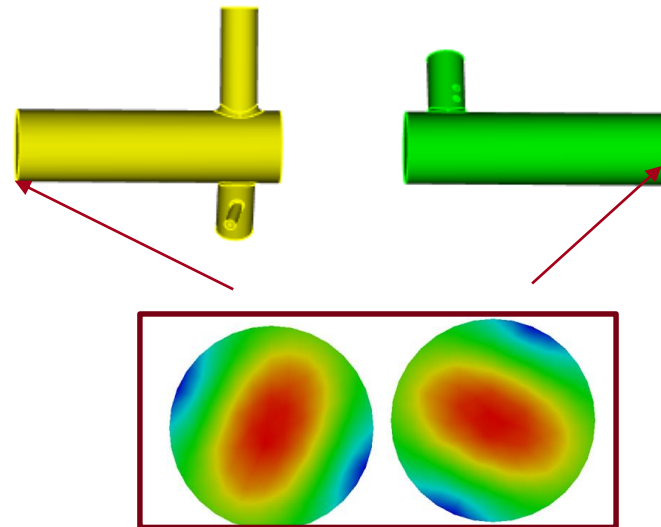
# Nonlinear RF Eigenvalue Problem for 3<sup>rd</sup> Dipole Band Modes

- It was found that there were trapped modes (~2.58GHz) in TDR CM;
- The trapped modes can propagate through the cavities to the beampipe absorber;
- Determining the propagating 3<sup>rd</sup> dipole band mode damping needs to solve a nonlinear eigenvalue problem.

$$F(\lambda) = K - \lambda M + i\lambda W^{TEM} + i\sqrt{\lambda - k_1^2} W_{11}^{TE} + i\sqrt{\lambda - k_2^2} W_{11}^{TE}$$



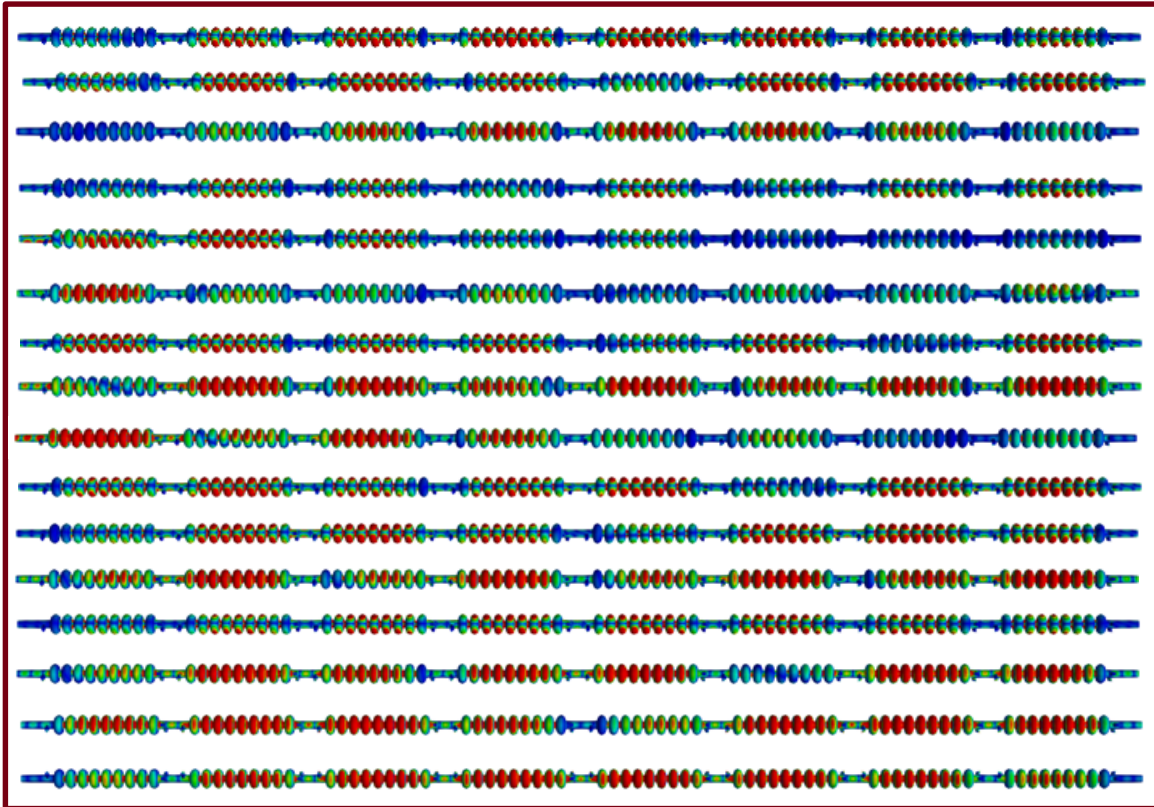
HOM coupler ports:  
TEM,  $F_c=0$



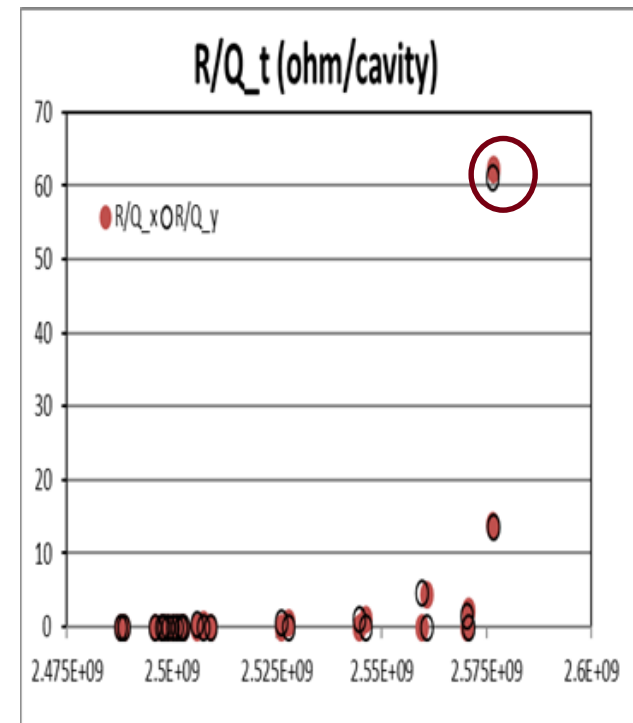
Beampipe ports:  
TE<sub>11</sub>  $F_{c1}=F_{c2}=2.2525\text{GHz}$

# The Last Mode of the 3<sup>rd</sup> Dipole Band in TDR CM

- Searching for the trapped modes in the 3<sup>rd</sup> dipole band in TDR CM;



3<sup>rd</sup> dipole band in a single cavity



## Omega3P simulation parameters –

~ 3 million curved element mesh, ~ 20 million DOFs,  
960 cores, 1.8 TB of memory on NERSC Edison, 1min per mode

# The Trapped Mode Damping Results

- The trapped mode damping factors calculated using Omega3P agree well with measurements at DESY, and the simulated frequencies are shifted to lower values due to cavity imperfection;
- Furthermore, cavity imperfection should be taken into account to study the propagating HOMs damping.

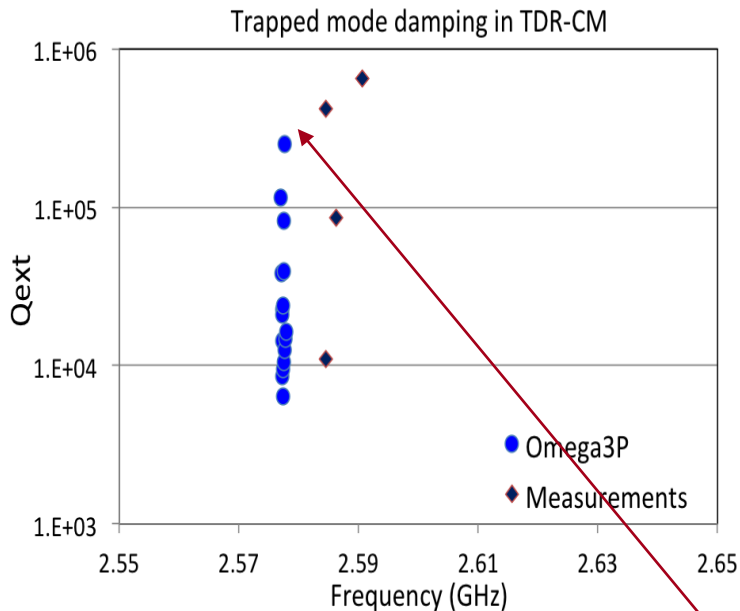


Table 1. Results of HOM investigations for the last mode of the 3<sup>rd</sup> dipole passband ( $R/Q = 15 \Omega/\text{cm}^2$ )

Cavity nr./module	Freq. [GHz]	Q
#3 (S10) / 1	2.5845	$1.1 \cdot 10^4$
#6 (S11) / 1	2.5862	$8.6 \cdot 10^4$
#5 (A15) / 2	2.5845	$4.2 \cdot 10^5$
#7 (S28) / 3	2.5906	$6.5 \cdot 10^5$

Courtesy of Nicoleta-Ionela Baboi



3<sup>rd</sup> dipole band trapped mode in TDR CM

# Summary

- Recently developed CORK algorithm in Omega3P can solve a nonlinear eigenvalue problem in accelerator RF cavities arising from waveguide boundary conditions;
- Omega3P nonlinear eigensolver has been benchmarked against S3P;
- The simulated damping factors of the propagating 3<sup>rd</sup> dipole band trapped modes in TDR 8-cavity CM using Omega3P agree well with measurements at DESY;
- Omega3P parallel computation enables a direct method to calculate propagating modes in large scale problems.

Learn more on modeling capabilities of ACE3P at  
ACE3P CW18 Accelerator Code Workshop, Nov. 5 to 9, 2018

<https://conf.slac.stanford.edu/cw18/>